Air-sea coupling shapes North American hydroclimate response to ice sheets during the Last Glacial Maximum

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22   **Abstract**

23   The Western U.S. is vulnerable to hydrological stress, and insights from past climate periods are

24   helpful for providing historical benchmarks for future climate projections. Myriad evidence from

25   coupled models and paleoclimatic proxies suggests a major reorganization of west coast

26   hydroclimate during the Last Glacial Maximum (LGM, ~17-25 ka), such that the Southwest U.S.

27   was wetter than modern day and the Pacific Northwest was drier. Yet the fundamental

28   mechanisms underlying these hydroclimatic shifts remain unclear. Here, we employ a suite of

29   targeted model simulations to probe the influence of LGM Northern Hemisphere ice sheets on

30   west coast atmospheric dynamics. Whereas previous modeling studies have suggested that the

31   southward shift of LGM west coast precipitation was driven only by the mechanical steering of

32   atmospheric circulation by elevated ice sheet topography, we find this to be an artifact of earlier

33   simulations that neglected realistic air-sea interaction. Instead, our simulations indicate that ice

34   sheet albedo induced a pattern of North Pacific sea surface temperatures, reinforced by ocean-

35   atmosphere feedbacks, that shifted the large-scale atmospheric circulation as well as the

36   latitudinal distribution of west coast precipitation southward during the LGM. Crucially, we find

37   that atmosphere-ocean feedbacks that sustained this ice sheet albedo-induced temperature pattern

38   in the LGM could drive similar hydroclimatic changes today.

39

1 Introduction

The hydroclimate of western North America is heavily influenced by intermittent wintertime storms ("atmospheric rivers") that can produce extreme precipitation, especially when these storms encounter substantial orographic features. For California, in particular, annual water availability hinges on the occurrence of just a handful of such atmospheric river events (Ralph and Dettinger, 2011). This region's acute vulnerability to climate variability and change is exacerbated by its growing human population. Therefore, understanding the fundamental drivers of hydroclimatic change over the west coast of the United States is essential for predicting the future habitability and agricultural viability of this region.

Large amplitude shifts in past climate, beyond the relatively brief instrumental record, provide essential insights into such projections. In this regard, the Last Glacial Maximum (LGM) serves as a useful test case of a relatively recent stable climate state that (a) was substantially different than our present climate and (b) for which we have reasonable constraints on important boundary conditions including orbital geometry, greenhouse gas (GHG) concentrations, sea level, ice volume, etc. To effectively leverage the LGM as a testbed for understanding future change, however, it is necessary to disentangle the climate impacts of forcings that are presently changing (e.g., GHG concentrations) from those that are not (e.g., continental ice sheets).

It has long been suggested that a major non-GHG driver of LGM atmospheric circulation change was the large topographic extent of the ice sheets over North America (Figure 1). Ice sheet forcing has been linked to the dramatic reorganization of west coast hydroclimate during the LGM (Anderson et al., 1988; Broccoli and Manabe, 1987; Cook and Held, 1988; Wong et al., 2016; Yanase and Abe-Ouchi, 2010), which is suggested to have been considerably wetter than present throughout the southwestern U.S. by a preponderance of general circulation model

65 (GCM) simulations and paleoclimate proxy reconstructions (Anderson et al., 1988; Bartlein et

66 al., 2011; Benson et al., 2011; Ibarra et al., 2014; Kirby et al., 2013; Lachniet et al., 2014;

67 Lemons et al., 1996; Lora et al., 2017; Maher et al., 2014; Oster et al., 2015; Owen et al., 2003;

68 Reheis et al., 2012; Seltzer et al., 2019; Thompson et al., 1999). In particular, coupled model

69 simulations nearly universally indicate a north-south dipole in LGM-to-present precipitation

70 change, such that western North America was wetter in the south and drier in the north (Lora,

71 2018; Oster et al., 2015).

72   While models and (to a lesser extent) proxies broadly agree on the existence of this north-

73 south LGM-to-present precipitation dipole response, the physical mechanisms that drove this

74 shift remain an open question. One leading hypothesis is that the North Pacific jet stream was

75 either split or deflected southward, so that extratropical storms and moisture transport occurred

76 further south (Lora et al., 2017). Many studies indicate that the southward deflection of the North

77 Pacific jet may have been the result of mechanical forcing imposed by North American ice

78 sheets, which covered much of modern day Canada and stood ~3 km taller than the present day

79 land surface (Anderson et al., 1988; Bartlein et al., 1998; Bromwich et al., 2004; Cook and Held,

80 1988; Kim et al., 2008; Kutzbach and Wright, 1985; Manabe and Broccoli, 1985). Other studies

81 point to the dominant influence of high latitude cooling induced by the high albedo of ice sheets

82 (Bhattacharya et al., 2017; Manabe and Broccoli, 1985; Morrill et al., 2018; Roberts et al., 2019;

83 Yanase and Abe-Ouchi, 2010).

84   Recently, Roberts et al. (2019) found that wintertime atmospheric stationary wave

85 patterns over the North Pacific during the LGM are primarily driven by ice sheet topography, but

86 are enhanced by ice sheet albedo. In contrast, they showed that summertime stationary wave

87 patterns are mainly driven by diabatic heating from the ice sheet albedo with little contribution

88   from ice sheet topography. In addition, they suggested that ocean dynamics also play a key role

89   in persisting atmospheric circulation anomalies from the summer into the winter. This is

90   consistent with other previous analyses that have broadly implicated ocean-atmosphere coupling

91   as an important factor in modulating the overall response to LGM boundary conditions (DiNezio

92   et al., 2018; Löfverström et al., 2014; Yanase and Abe-Ouchi, 2010).

93   While these recent studies have provided significant insight into how the different

94   characteristics of the Laurentide and Cordilleran ice sheets may have influenced North Pacific

95   atmospheric circulation, several important questions remain. First, how do different ice sheet-

96   driven stationary wave changes (e.g., Roberts et al. 2019) map onto potential shifts in the North

97   Pacific jet stream? Following from this, do different responses of the zonal circulation to ice

98   sheet topography and ice sheet albedo lead to different hydroclimate shifts along western North

99   America? Additionally, what are the specific ocean dynamical processes and air-sea feedbacks

100   that modulate the seasonally varying response of the North Pacific atmosphere to difference ice

101   sheet boundary conditions?

102   Here we address these questions by exploring a suite of model simulations (Table 1)

103   conducted with the Community Earth System Model version 1.2 (CESM1) to uncover the drivers

104   and mechanisms of North American west coast hydroclimate change during the LGM, with

105   targeted experiments to distinguish the mechanical (e.g., orographic) and thermodynamic (e.g.,

106   albedo-driven) impacts of large continental ice sheets on North Pacific atmospheric circulation.

107   In addition, we build upon Roberts et al. (2019) and others by including a comprehensive range

108   of ocean model complexities, which allow us to directly diagnose the specific ocean-atmosphere

109   interactions and ocean circulation changes that are most relevant to the overall climate response.

110

2 Data and Methods

2.1 Climate model simulations

The climate model simulations used in this study were conducted by DiNezio et al. (2018). Here, we will provide a broad overview of the major model components, as well as a brief summary of the LGM boundary conditions applied to our various simulations. However, we direct readers to the Supplementary Materials Sections S3-S5 in DiNezio et al. (2018) for a complete description of the climate model components and boundary conditions. Our simulations were carried out using the CESM1 global climate model, which consists of the Community Atmospheric Model version 5 (CAM5; ~2˚ horizontal resolution with 30 vertical pressure levels). CAM5 is coupled to a full-depth ocean model—the Parallel Ocean Program version 2 (POP2)—featuring ~1˚ horizontal resolution, with increased meridional resolution of ~1/3˚ in the tropics, and 60 vertical levels. Additionally, CESM1 features the Community Land Model version 4 (CLM4) with ~2˚ horizontal resolution.

We compare four sets of equilibrated CESM1 experiments with varying degrees of ocean complexity. They include: (1) A fully forced LGM experiment (LGM$_{Full}$) with all major glacial boundary conditions (GHGs, sea level, orbital, ice sheet topography and albedo, etc.) prescribed following the Paleoclimate Modeling Intercomparison Project Phase III (PMIP3) LGM protocols at 21,000 years ago, (2) A "White Mountain" (WM) simulation that includes all of the LGM ice sheet boundary conditions (e.g., extent, elevation, surface properties), but with the remaining boundary conditions (GHGs, orbit, sea level) set to pre-industrial levels, (3) A "Green Mountain" (GM) simulation that includes identical topography to the WM, but with land surface properties and the remaining boundary conditions set to pre-industrial values, and (4) A pre-industrial control simulation (CTL) with boundary conditions following the PMIP3 pre-industrial

134   protocols. Figure 1 further illustrates the different LGM and CTL land surface properties used in

135   our experiments.

136   Ice sheet topography and glacier extent in each of our simulations were prescribed

137   following PMIP3 protocols, which combine three different ice sheet reconstructions (Abe-Ouchi

138   et al., 2015). Otto-Bliesner et al. (2006) showed that the simulated LGM climate in earlier

139   versions of CESM could be sensitive to the choice of ice sheet reconstruction. However, our

140   LGM results do not significantly differ when repeating our analysis using an updated ice sheet

141   reconstruction that is independent of the ice sheet blend used in PMIP3.

142   The WM and GM experiments include repeated simulations with an atmospheric model

143   that is: (a) uncoupled from the ocean model (AGCM only; i.e., intrinsic atmospheric variability

144   only), (b) coupled to a fixed, 50 m depth slab ocean model (SOM; i.e., no ocean circulation, but

145   ocean-atmosphere interactions possible), and (c) fully coupled to a dynamical ocean model

146   (DOM). In our AGCM experiments, global sea surface temperature (SST) and sea ice extent

147   were prescribed with climatological values from the CTL. Differences among the SOM/DOM

148   and AGCM configurations highlight the role of interactive ocean processes in shaping the overall

149   climate response in the respective GM/WM experiments, while differences between the DOM

150   and SOM isolate the more subtle roles of ocean circulation changes versus mixed layer ocean-

151   atmosphere interactions. (Hereafter we refer to these different experiments and ocean model

152   complexities using the notation "Experiment$_{model}$", such as "WM$_{DOM}$" for the White Mountain

153   experiment featuring a fully dynamical ocean model). Each of these experiments (summarized in

154   Table 1) was integrated until the climate system (including the deep ocean for DOM simulations)

155   reached close to radiative equilibrium (see DiNezio et al., 2018 Supplenary Materials Section S5

156   for more details). Each was then integrated forward at equilibrium for an additional number of

157   years (Table 1) to calculate climatological fields for comparison.

158   Unless otherwise stated, all climate anomalies reported in this study are at least 95%

159   significant using a Student's t-test and are relative to a CTL experiment run with the

160   corresponding ocean model complexity. When compared to CTL, atmospheric circulation

161   anomalies in WM are due to the combination of ice sheet albedo and topography (i.e., the

162   combination of mechanical and thermodynamic forcing), while circulation anomalies in GM are

163   due to mechanical forcing from the topography alone. Their difference (WM minus GM)

164   estimates the climate response to ice sheet albedo alone (i.e., thermodynamic forcing only).

165   Differencing WM and GM simulations to isolate the direct ice sheet albedo forcing on the

166   climate system implicitly assumes that ice sheet mechanical and thermodynamic forcing are

167   linearly additive. To assess the validity of this assumption, we analyze an additional albedo-

168   forced CESM1 simulation also conducted by DiNezio et al. (2018). This extra simulation

169   features a dynamical ocean model and was forced with the distribution of continental ice set at

170   LGM values, but over the modern topography. All other boundary conditions (GHGs, orbital,

171   topography, etc.) were set to pre-industrial values. This approach isolates the cooling effect of

172   the albedo of the ice sheets relative to modern topography. Following the naming convention of

173   our other experiments, we refer to this fully coupled simulation as White Plains (WP$_{DOM}$; Table 1

174   and Figure 1). By comparing WP$_{DOM}$ anomalies to those from WM$_{DOM}$ we are able to discern

175   whether the elevation of the ice sheet albedo forcing introduces significant nonlinearities in the

176   large-scale climate response.

177   Finally, to test the sensitivity of our LGM$_{Full}$ results to our choice of model, we compare

178   LGM$_{Full}$ to an ensemble of seven models forced with identical LGM boundary conditions taken

179   from the PMIP3 archive. These models are listed in Table S1.

180

181   *2.2 Precipitation and sea surface temperature proxy estimates*

182   We validate our simulated hydroclimate results against several proxy estimates of LGM

183   western North American precipitation change, as in Lora et al. (2017). These proxies consist

184   primarily of pollen data (Bartlein et al., 2011), but also include other quantitative precipitation

185   estimates (Ibarra et al., 2014; Lemons et al., 1996; Maher et al., 2014; Thompson et al., 1999). In

186   addition, we compare simulated North Pacific SST anomalies to the gridded proxy estimates

187   generated by Tierney et al. (2020). These estimates combine a large collection of geochemical

188   proxies of SST with an isotope-enabled version of CESM to produce a spatially coherent

189   reconstruction of LGM temperatures using data assimilation. Specifically, we use version 2.0 of

190   their LGM data assimilation runs, which features a wider array of model priors than the original

191   dataset. See Tierney et al. (2020) for more details.

192

193   **3 Results**

194   Our CESM1 LGM$_{Full}$ experiment produces patterns of North Pacific atmospheric

195   circulation and North American west coast hydroclimate change that compare well with the

196   ensemble mean results of seven PMIP3 models forced with the same LGM boundary conditions

197   (Figure 2). Specifically, both our CESM1 LGM$_{Full}$ simulation and the PMIP3 models produce a

198   clear north-south dipole in boreal winter precipitation anomalies along the North American west

199   coast (Figure 2b,d), such that the Southwestern U.S. is wetter than modern and the Pacific

200 Northwest, from Washington to Alaska, is drier than modern. These similarities are further

201 highlighted by the spatial pattern correlation between our CESM1 LGM$_{Full}$ simulation and the

202 PMIP3 ensemble mean, which is R = 0.92 for the zonal wind anomalies and R = 0.85 for the

203 precipitation change, both significant at 95% confidence.

204   The simulated precipitation anomalies are also highly correlated in space with the proxy

205 reconstructions of LGM-to-present precipitation changes (circles Figure 2b,d and Table S2). The

206 correlation between the proxy estimates and the LGM$_{Full}$ precipitation anomalies at the nearest

207 grid cell is R = 0.67, which is significant at 95% confidence. These precipitation changes are due

208 to an intensified and southward shifted wintertime jet stream, which correspondingly shifts

209 southward the distribution of landfalling North Pacific storms (Figure 2a,c). These model-data

210 comparisons are consistent with many previous studies (Anderson et al., 1988; Bartlein et al.,

211 2011; Ibarra et al., 2014; Lachniet et al., 2014; Lemons et al., 1996; Lora, 2018; Lora et al.,

212 2017; Maher et al., 2014; Oster et al., 2015; Owen et al., 2003; Seltzer et al., 2019; Thompson et

213 al., 1999) and give us confidence in our CESM1 modeling framework.

214

215 *3.1 North Pacific climate response to ice sheet topography*

216   Direct orographic forcing from the ~3 km (Figure 1b) tall North American ice sheets in

217 GM$_{AGCM}$ produces a southward displacement and intensification of the wintertime North Pacific

218 jet, as well as a general tendency for a north-south shift in west coast hydroclimate (Figure 3a-b),

219 which is similar to the coupled response with full LGM forcings (Figure 2a-b). In particular,

220 GM$_{AGCM}$ produces a strong deceleration of the westerlies at ~50°N (Figure 4a) and throughout

221 the atmospheric column that closely aligns with the latitude at which LGM North American ice

222 sheets begin to rise above modern topography (Figure 4a; orange and purple lines), indicative of

223   a direct mechanical slowdown of the westerly flow ahead of the elevated land surface at these

224   latitudes. The weakened atmospheric circulation is then steered southward, leading to an

225   acceleration of the westerlies near 30°N. Overall, these zonal wind anomalies produce an

226   equatorward shift of the North Pacific jet stream, especially in the mid-to-lower troposphere.

227   These circulation changes are apparent in the boreal winter geostrophic streamfunction (Figure

228   S1a), which shows an anomalous low over the North Pacific as part of a generally barotropic

229   stationary wave response throughout the Northern Hemisphere, consistent with results presented

230   by Roberts et al. (2019) using a different climate model.

231   Based on the similarities between GM_{AGCM} (Figure 3a-b) and LGM_{Full} (Figure 2a-b), one

232   might reach the same conclusion as many prior AGCM-based studies that also sought to

233   disentangle the role of ice sheet topography versus ice sheet albedo on LGM climate (Anderson

234   et al., 1988; Bartlein et al., 1998; Bromwich et al., 2004; Cook and Held, 1988; Kim et al., 2008;

235   Kutzbach and Wright, 1985; Manabe and Broccoli, 1985)—that ice sheet topography alone is

236   sufficient to explain the overall changes to the LGM North Pacific jet stream, without the need

237   for ocean-atmosphere interactions. However, when coupling our AGCM to interactive ocean

238   models, we in fact find nearly the opposite result. Specifically, both GM_{SOM} and GM_{DOM} show a

239   broad deceleration of the westerlies without a clear latitudinal shift in the jet stream (Figures 3c,e

240   and 4b-c) as well as drier conditions in the Southwestern U.S. relative to the control (Figure

241   3d,f). This raises a key question: What ocean-atmosphere feedbacks and/or ocean circulation

242   changes lead to such markedly divergent circulation and hydroclimate responses between our

243   uncoupled (GM_{AGCM}) and coupled (GM_{SOM} and GM_{DOM}) simulations?

244   Of central importance to addressing this question is the SST warming observed

245   throughout much of the subpolar North Pacific in GM_{DOM} (Figure 5d) and GM_{SOM} (Figure S2a).

246   These warm anomalies extend upward throughout the atmospheric column (Figure S3b-c),

247   decreasing the meridional temperature gradient between the tropics and the poles, and weakening

248   the North Pacific jet stream by thermal wind balance. An ocean mixed layer heat budget

249   (Supplementary Materials Section S1; Figure S4) indicates that this anomalous subpolar

250   warming can be attributed primarily to three important factors: (1) increased surface shortwave

251   radiation due to decreased low clouds, (2) decreased evaporative cooling driven by the

252   mechanically weakened surface westerlies, and (3) a dynamical ocean response that primarily

253   reflects warm SST advection via Ekman transport (Figure S6a-b). In contrast, without the added

254   atmospheric response to anomalous ocean temperatures, GM_{AGCM} produces cool anomalies in the

255   mid-troposphere near ~40°N (Figure S3a), consistent with cold air advection from high-latitudes

256   by the anomalous low (Figure S1a). This cooling enhances the local meridional temperature

257   gradient and reinforces the direct mechanical shift of the jet stream, opposite to the behavior

258   observed in GM_{SOM} and GM_{DOM}.

259   The subpolar warming seen in GM_{DOM} is a robust response to ice sheet mechanical

260   forcing, even when considering the additional impacts of lower LGM GHGs (see Supplementary

261   Materials Section S2). This is because reducing GHGs does not significantly alter the anomalous

262   meridional temperature gradient set up by the direct mechanical ice sheet forcing in GM_{DOM}.

263   However, adding in the full range of altered LGM boundary conditions (including ice sheet

264   albedo) does lead to a much different SST response than in GM_{DOM}. In particular, LGM_{Full}

265   exhibits colder-than-modern conditions throughout the North Pacific (Figure 5c), with the

266   coldest SST anomalies found near the Kuroshio Extension region in a pattern that broadly

267   resembles a positive phase of the Pacific Decadal Oscillation (PDO; Mantua et al., 1997).

268   This PDO-like cooling is supported by various proxy reconstructions of North Pacific

269   SSTs (Figure 5a) (Gray et al., 2020; Rae et al., 2020; Tao et al., 2013; Tierney et al., 2020), and

270   is also seen in the PMIP3 models analyzed here (Figure 5b and Figure S8) as well as more recent

271   PMIP4-era models (Kageyama et al., 2021). We do note that, while both our LGM$_{Full}$ simulation

272   and the PMIP3 ensemble mean show weak cooling in the subpolar North Pacific, proxy

273   estimates indicate that the subpolar LGM may have been slightly warmer than modern (Rae et

274   al., 2020). Although, these warm anomalies may have been largely confined to a relatively small

275   area north of 50°N (Figure 5a). Regardless, the broad-scale differences between the ice sheet-

276   driven SST anomalies in our GM$_{DOM}$ and LGM$_{Full}$ experiments further suggest that ocean

277   adjustment to direct mechanical ice sheet forcing (i.e., extensive subpolar North Pacific

278   warming) is not a foundational component of the overall response of the coupled climate system

279   to full LGM boundary conditions.

280

281   *3.2 North Pacific climate response to ice sheet albedo*

282   We now consider the combined contributions of elevated ice sheet topography (as in

283   GM) and high ice sheet surface albedo to the overall LGM North Pacific climate response. In

284   WM$_{AGCM}$, the zonal wind and precipitation anomalies are very similar to GM$_{AGCM}$ (Figures 3a-b

285   and 3g-h). Differencing the WM$_{AGCM}$ and GM$_{AGCM}$ simulations, we see little response in the

286   zonal wind and precipitation, thus ice sheet topography is almost entirely responsible for the

287   large-scale atmospheric response in both of these uncoupled simulations (Figures 3m-n). This

288   result is consistent with Roberts et al. (2019), who showed that the wintertime North Pacific

289   stationary wave response to ice sheet topography is similar to the response to ice sheet

290   topography and ice sheet albedo combined (see their Figure 2). Indeed, our WM$_{AGCM}$ and

291   GM$_{AGCM}$ simulations produce nearly identical wintertime stationary wave fields to theirs

292   (Figures S1a and S1g). When coupled to an ocean model in WM$_{SOM}$ and WM$_{DOM}$, the additional

293   thermodynamic forcing from high ice sheet albedo is responsible for an intensification and

294   southward displacement of the boreal winter North Pacific jet, particularly in the central/eastern

295   North Pacific, leading to the characteristic north-south dipole in western North American

296   precipitation (Figure 3i-l). This can be seen more clearly when differencing the respective WM

297   and GM simulations (Figure 3o-r).

298   To what extent can the differences between our WM and GM simulations be explained

299   by ice sheet albedo versus potential nonlinearities in the climate system? Our WP$_{DOM}$ simulation

300   isolates the climate response to ice sheet albedo relative to modern day topography, and provides

301   a useful tool to address this question. In WP$_{DOM}$, the boreal winter North Pacific jet experiences

302   an intensification and southward shift that is qualitatively similar to that seen in WM$_{DOM}$ minus

303   GM$_{DOM}$ (Figure S9a,c). This leads to a north-south shift in precipitation in WP$_{DOM}$ that is also

304   comparable to WM$_{DOM}$ minus GM$_{DOM}$ (Figure S9b,d). Similar results are seen when comparing

305   anomalies from WM$_{DOM}$ to the linear combination of anomalies from WP$_{DOM}$ and GM$_{DOM}$

306   (Figure S9e-h). While the spatial patterns of anomalies between the different ice sheet albedo

307   forcing estimates are qualitatively similar, the zonal wind anomalies in WP$_{DOM}$ tend to be weaker

308   than in WM$_{DOM}$ and the moistening signal in the southwest U.S. is overall less clear in WP$_{DOM}$.

309   These results suggest that some nonlinearities likely contribute to the differences between our

310   WM and GM simulations and that WM minus GM should not be quantitatively interpreted.

311   Instead, the albedo forcing estimates provided by taking WM minus GM (i.e., Figures 3m-r and

312   4g-i) should be seen as a qualitative indicator of how elevated ice sheet albedo forcing imprints

313   on large-scale climate.

314   The WP$_{DOM}$ experiment discussed above is an idealized representation of LGM

315   conditions (i.e., it is unlikely that there would ever be such extensive ice sheet albedo forcing

316   across relatively low topography in the real world). However, our WM$_{DOM}$ simulation provides a

317   much more realistic representation of total ice sheet forcing with which to compare to our full

318   LGM simulation. For example, the WM$_{SOM}$ and WM$_{DOM}$ zonal wind and precipitation anomalies

319   are nearly identical to those seen in LGM$_{Full}$ (comparing Figures 2a-b and 3i-l), suggesting that

320   thermodynamic forcing from the bright ice sheets is key in setting the overall LGM North Pacific

321   atmospheric response. However, these anomalies are also heavily reliant upon coupled ocean-

322   atmosphere interactions, as indicated by the difference of WM minus GM zonal wind fields

323   across the different ocean configurations (Figure 3m-r and Figure 4g-i). In particular, ocean-

324   atmosphere coupling enables an albedo-forced acceleration of the westerlies throughout the

325   atmospheric column between ~30-40°N in WM$_{SOM}$ and WM$_{DOM}$. This response is only slightly

326   larger in WM$_{DOM}$, indicating that dynamical ocean circulation changes are of secondary

327   importance to mixed layer air-sea interactions in driving these anomalies. Overall, our

328   experiments suggest that the inclusion of ocean-atmosphere coupling in WM$_{SOM}$ and WM$_{DOM}$

329   acts to reinforce the mechanically induced southward jet shift implied by the uncoupled

330   WM$_{AGCM}$ (Figure 4g).

331   How do ocean-atmosphere feedbacks and/or ocean circulation changes combine with the

332   direct atmospheric response to thermodynamic ice sheet forcing to form the wintertime WM$_{SOM}$

333   and WM$_{DOM}$ North Pacific climate anomalies? Both WM$_{DOM}$ (Figure 5e) and WM$_{SOM}$ (Figure

334   S2b) produce PDO-like SST anomaly patterns that are similar to that seen in proxy estimates, our

335   CESM1 LGM$_{Full}$, and the PMIP3 ensemble mean (Figure 5a-c). However, the peak anomalies in

15

336   WM$_{DOM}$—exceeding -4˚C in magnitude—are only about half as large as in LGM$_{Full}$ since the

337   latter experiences additional cooling from changes in orbital configuration and lower GHGs.

338   Using a mixed layer heat budget analysis, we find that the WM$_{DOM}$ cooling primarily

339   results from: (1) decreases in downward surface shortwave radiation south of 40˚N, (2) ocean

340   circulation changes poleward of 40˚N, and (3) increased evaporation from strengthened surface

341   winds and an increased air-sea temperature contrast (Figure S5). The cooling generated by ocean

342   circulation changes has contributions from anomalous Ekman heat advection (Figure S6c-d), but

343   is also consistent with a southward expansion of the North Pacific subpolar gyre circulation and

344   a southward shift of the subarctic SST front in response to strengthened polar easterlies (Gray et

345   al., 2020). The anomalous temperatures around 40˚N extend throughout much of the atmospheric

346   column to ~300 mb, increasing the meridional temperature gradient over the North Pacific and

347   strengthening the jet stream in WM$_{SOM}$ and WM$_{DOM}$ by thermal wind balance (Figure S3e-f).

348   To confirm that the underlying SST field is indeed driving the boreal winter upper-level

349   zonal wind response, we conducted an additional set of AGCM simulations forced with pre-

350   industrial land surfaces (i.e., no LGM ice sheets) and the long-term monthly mean SST output

351   from WM$_{DOM}$ (referred to as WM$_{SST}$; see Supplementary Materials Section S3). The consistency

352   between the WM$_{SST}$ global experiment and WM$_{DOM}$ is striking (Figures 4f and 6b), and suggests

353   that nearly the entire LGM North Pacific atmospheric circulation—and thus North American

354   hydroclimatic—response during boreal winter is attributable to the ice sheet-driven SST field

355   and not to the direct atmospheric response to the combined effects of ice sheet orography and

356   albedo (as estimated by WM$_{AGCM}$). Regional SST forced experiments show that the

357   aforementioned North Pacific SST pattern accounts for the majority of the global SST forced

358   response, particularly the accelerated westerlies between ~30˚-40˚N (Figure 6c-e). Therefore, the

359   direct atmospheric response to continental ice sheets is only important to the extent that it

360   contributes to the formation of this PDO-like SST pattern (Figure 5e).

361

362   *3.3 Physical mechanisms linking ice sheet albedo to North Pacific SST changes*

363   The North Pacific ocean and atmosphere are highly coupled in boreal winter (Alexander,

364   1992; Alexander and Scott, 1997). Therefore, it is important to diagnose the extent to which the

365   SST cooling described above is a source or a symptom of the atmospheric circulation anomalies

366   in the WM experiments. In doing so, we can better understand the physical pathways by which

367   the direct large-scale atmospheric response to a tall/bright ice sheet (as in WM_{AGCM}) first

368   imprints onto the North Pacific ocean before coupling back to the atmosphere (as in

369   WM_{SOM}/WM_{DOM}). The direct thermodynamic impact of North American ice sheets on the large-

370   scale atmosphere should be strongest in boreal summer when the zonal winds are weakest and

371   most sensitive to diabatic heating (Roberts et al., 2019; Ting, 1994). Additionally, anomalous

372   cooling over land caused by high ice sheet albedo will be largest in summer when there is more

373   insolation to reflect.

374   The seasonally varying influence of ice sheet albedo on large-scale atmospheric

375   circulation was discussed in detail by Roberts et al. (2019), and is reproduced here as the

376   difference between WM and GM stationary wave fields in our AGCM simulations (Figure S1m-

377   r). Importantly, the summertime atmospheric circulation anomalies are more similar across the

378   AGCM/SOM/DOM configurations than those in winter for both WM and GM simulations. The

379   insensitivity of the boreal summer climate response to the presence of an ocean in our

380   simulations suggests that the ocean and the large-scale atmosphere are largely decoupled during

381   this season. This is confirmed in our additional WM_{SST} experiment, which does not reproduce

382   the WM$_{DOM}$ zonal wind anomalies in boreal summer (not shown). Therefore, by relating the WM

383   annual mean SST cooling (Figure 5e) to summer North Pacific climate anomalies, we can

384   reliably diagnose the direct atmospheric response to the tall/bright ice sheet, even in our coupled

385   simulations.

386   As discussed previously, decreases in surface shortwave radiation play an integral role in

387   cooling the North Pacific in WM$_{DOM}$, with cloud changes accounting for nearly the entire surface

388   shortwave contribution to the SST anomalies (Figure 7a-c). Our simulation shows an increase in

389   total cloud fraction in this region; however, the shortwave cloud radiative forcing (Figure 7b) is

390   primarily due to an increase in low clouds, such as marine stratocumulus (Figure 8a). Low-lying

391   marine stratocumulus clouds are an important amplifying factor for SST anomalies in the North

392   Pacific (Norris et al., 1998; Ronca and Battisti, 1997), often producing a positive feedback with

393   increasing (decreasing) SSTs leading to decreasing (increasing) low cloud amounts and

394   consequently more (less) downward shortwave radiation at the surface (Amaya et al., 2020).

395   Direct ice sheet-driven large-scale atmospheric circulation changes can be connected to

396   local SST anomalies independent of these ocean-atmosphere feedbacks in WM$_{AGCM}$. This

397   uncoupled simulation produces a strikingly similar pattern of surface shortwave-induced SST

398   anomalies, which are also dominated by changes in shortwave cloud radiative forcing (Figure

399   7d-f). Aside from the similar spatial pattern, the anomalies are generally weaker in WM$_{AGCM}$

400   than in WM$_{DOM}$, suggesting that local low cloud-SST feedback amplifies both the low cloud and

401   SST response in WM$_{DOM}$. Averaging the fractional change in low cloud amount in the central

402   North Pacific, we see that the increase in low clouds in WM$_{AGCM}$ and WM$_{DOM}$ is largest in boreal

403   summer (Figure 7g), with the difference between the two further highlighting the role of ocean-

404   atmosphere feedbacks in amplifying these changes in WM$_{DOM}$ throughout the year.

18

405   The increase in summertime low cloud coverage in both WM$_{AGCM}$ and WM$_{DOM}$ is due

406   primarily to a stronger temperature inversion above the marine boundary layer west of ~160˚W

407   and increased moisture east of ~160˚W (Figure 8). These changes are driven by an anomalous

408   cyclonic circulation in the lower atmosphere seen in each of our WM simulations during boreal

409   summer (Figure S1j-l). This low pressure anomaly advects cold air from the high latitudes in the

410   western North Pacific (increasing stability at low levels), while advecting additional moisture

411   from the tropics in the eastern North Pacific. See Supplementary Materials Section S1.4 for more

412   details on the cloud parameters analyzed here.

413   Overall, the consistent change in cloudiness in the WM$_{AGCM}$ simulation provides

414   evidence that the direct large-scale atmospheric response to North American ice sheets first

415   influences North Pacific SSTs during boreal summer through changes in cloud fraction and

416   surface shortwave radiation. These SST anomalies then persist into winter through a combination

417   of positive feedbacks including low cloud-SST and wind-evaporation-SST interactions, as well

418   as ocean circulation adjustments, where they then strongly influence the position of the North

419   Pacific jet and, by extension, North American west coast hydroclimate.

420

421   **4 Summary**

422   In this study, we investigated the influence of continental ice sheets on North Pacific

423   atmospheric circulation and North American west coast hydroclimate during the LGM. Our

424   model results showed that, in response to full LGM boundary conditions (e.g., GHGs, ice sheets,

425   sea level, orbital configuration, etc.), the North Pacific jet is intensified and shifted southward

426   during boreal winter. This results in a wetter-than-modern-day Southwest U.S. and a drier-than-

427   modern-day Pacific Northwest, consistent with previous model simulations and paleoproxy

428   reconstructions (Anderson et al., 1988; Bartlein et al., 2011; Benson et al., 2011; Ibarra et al.,

429   2014; Kirby et al., 2013; Lachniet et al., 2014; Lemons et al., 1996; Lora et al., 2017; Maher et

430   al., 2014; Oster et al., 2015; Owen et al., 2003; Reheis et al., 2012; Seltzer et al., 2019;

431   Thompson et al., 1999). Decomposing this climate response into contributions from ice sheet

432   topography alone (i.e., Green Mountain) and ice sheet topography plus albedo (i.e., White

433   Mountain), we found that these circulation changes result from a complex combination of

434   mechanical and thermodynamic ice sheet effects that are critically dependent on the presence of

435   ocean-atmosphere interactions.

436   The step-by-step response of the North Pacific atmosphere to realistic continental ice

437   sheets (i.e., a White Mountain) is summarized in the schematic Figure 9. In short, the large-scale

438   stationary wave response to tall/bright ice sheets includes a low-level cyclonic circulation over

439   the North Pacific during boreal summer, supported also by previous studies (Roberts et al.,

440   2019). This circulation anomaly increases moisture transport to the eastern North Pacific, while

441   simultaneously advecting cool air from high-latitudes over the western North Pacific. The

442   increase in moisture combined with the increased lower atmosphere stability leads to an increase

443   in low clouds, which cool North Pacific SSTs between 30°N-40°N by reducing downward

444   surface shortwave radiation. A positive low cloud-SST feedback and subpolar gyre dynamical

445   adjustments (Gray et al., 2020) then combine to amplify both the cloud changes and the North

446   Pacific SST anomalies, allowing them to persist into boreal winter. During the winter, the ocean

447   couples back to the large-scale atmosphere and reinforces the tendency for a topographically-

448   forced southward shift of the North Pacific jet stream by increasing the North Pacific meridional

449   temperature gradient and accelerating the westerlies on the equatorward flank of the jet.

450

5 Discussion

451

452   Many of the above results are consistent with previous studies (e.g., Roberts et al., 2019;

453 Yanase and Abe-Ouchi, 2010), who used similar model simulations to investigate the importance

454 of different ice sheet characteristics to LGM atmospheric circulation as well as the influence of

455 ocean coupling in shaping those responses. However, our detailed analysis of North Pacific SST

456 patterns helps to fill a notable gap in our understanding of the specific physical pathways (e.g.,

457 low cloud feedbacks) by which the coupled climate system adjusts to LGM ice sheet forcing, and

458 helps to further highlight the relevance of LGM climate for modern applications. In particular,

459 our finding that ice sheet-driven SST patterns—rather than the direct atmospheric response to

460 mechanical forcing induced by the elevated topography of the LGM ice sheets—drove much of

461 the North Pacific atmospheric circulation response during the LGM has key implications for

462 present day climate variability and future change.

463   We suggest that the primary role of ice sheets on LGM hydroclimate was to provide a

464 suitably large perturbation to the climate system that led to the formation of an SST pattern that

465 bears strong resemblance to a well-known climatic mode (e.g., the PDO). In this sense, while the

466 presence (or removal) of a ~3 km tall ice sheet over North America bears no direct relevance to

467 modern climate, the possible emergence of North Pacific SST anomalies that are similar in

468 pattern and magnitude as those observed in our LGM simulations is distinctly plausible. While

469 such patterns may result from natural climate variations, our findings call attention to ongoing

470 anthropogenic factors—such as changes in Southeast Asian aerosol emissions—that have been

471 shown to promote a PDO-like SST response in coupled models (Smith et al., 2016) and may

472 exert a strong influence on western North American hydroclimate through air-sea interactions.

473 These regional SST patterns may further serve to alter the expected response of the mid-latitude

21

474 jet streams to future global warming (Lu et al., 2008), due to the modification of the local

475 meridional temperature gradient at the surface.

476   Finally, we acknowledge several important caveats and possible sensitives regarding our

477 analysis. For example, we note that our analysis is primarily based on simulations from a single

478 model—CESM1.2. However, portions of our results (particularly the large-scale atmospheric

479 circulation changes associated with different ice sheet characteristics) are nearly identical to

480 Roberts et al. (2019), who conducted similar White Mountain and Green Mountain experiments

481 using a different climate model. Indeed, the North Pacific stationary wave patterns in our

482 simulations (Figure S1) compare very well to those seen in their simulations (see Roberts et al.

483 2019, Figures 2 and 3). Therefore, we have confidence that the ice sheet-driven atmospheric

484 circulation changes discussed here are qualitatively insensitive to model choice.

485   Additionally, although we have identified a series of physically consistent pathways by

486 which the coupled climate system adjusts to ice sheet forcing (Figure 9), we note that these

487 mechanisms depend on an accurate representation of marine stratocumulus clouds and their

488 feedbacks. While the simulation of marine stratocumulus clouds has improved in recent model

489 generations (Engström et al., 2014), there is still large uncertainty among models in the expected

490 response of clouds to future greenhouse gas increase (Qu et al., 2014; Webb et al., 2013).

491 Nevertheless, Xiao et al., (2014) showed that CESM1 credibly reproduces the climate conditions

492 necessary to accurately simulate marine stratocumulus clouds in the North Pacific, including the

493 depth of the planetary boundary layer, which gives us confidence in our results. Regardless,

494 future work is needed to further explore the physical mechanisms identified here across a range

495 of other model configurations

496

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684

685 **Table 1** List of the CESM1 experiments used in this study. Crosses indicate the boundary

686 condition has been modified following PMIP3 protocols. All boundary conditions have been

687 applied globally. The availability of dynamical ocean model (DOM), slab ocean model (SOM),

688 or AGCM data is noted in parentheses. The equilibration period for each DOM run is shown in

689 the far right column. The column second from right indicates the additional integration period

690 used for calculating climatologies once each DOM run had reached close to radiative

691 equilibrium. Each SOM and AGCM simulation was spun-up for 10 years, and integrated for an

692 additional 100 years for calculating climatologies. See DiNezio et al. (2018) for complete details

693 regarding the initial conditions and the boundary conditions.

	GHGs	Orbit	Land mask & bathymetry	Topography	Albedo	Duration (yrs)	Equilibration (yrs)
LGM$_{Full}$ (DOM)	X	X	X	X	X	600	500
White Mountain (DOM, SOM, AGCM)				X	X	400	700
Green Mountain (DOM, SOM, AGCM)				X		200	100
White Plains (DOM)					X	400	500

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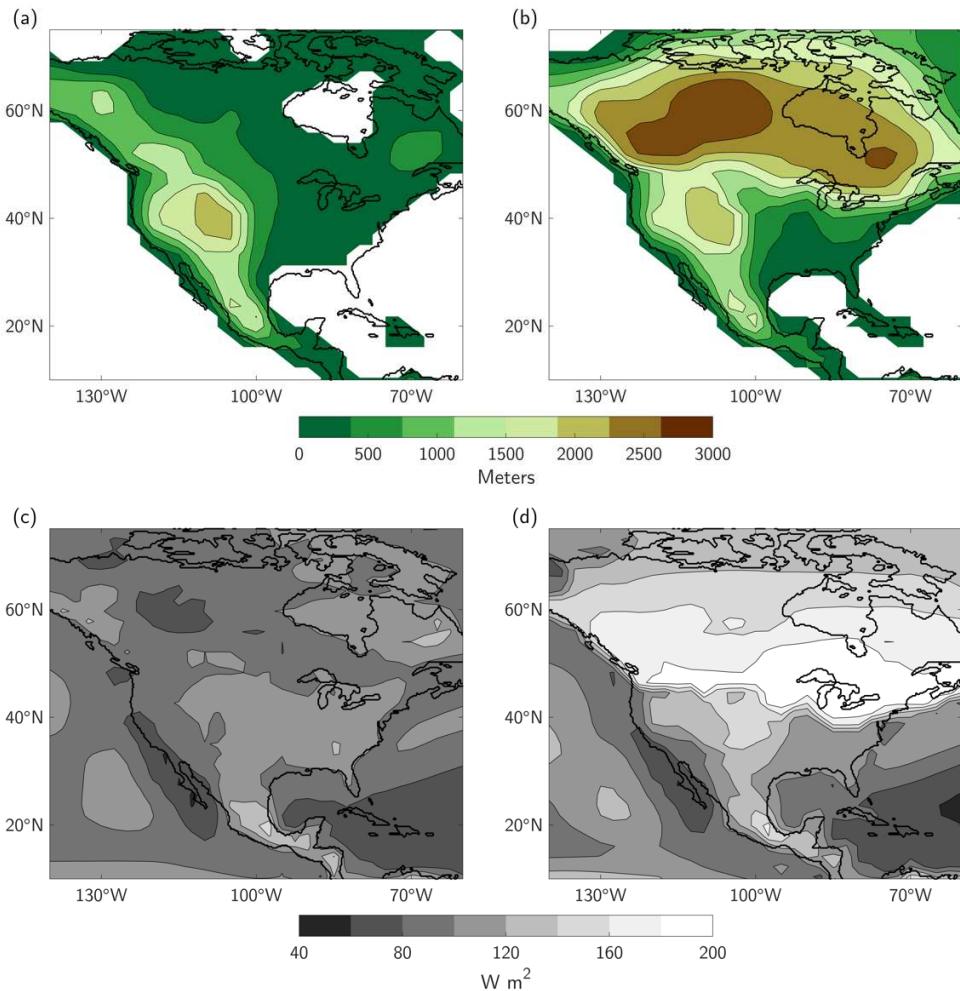
Figure 1 Ice sheet characteristics in our CESM1 simulations. (a)-(b) Surface geopotential divided by 9.8 m/s^2 (meters) and (c)-(d) annual mean upward shortwave radiation at the top of the atmosphere (W m^{-2}) in our CTL (left column) and LGM (right column) simulations. LGM$_{Full}$ and WM simulations have the ice sheet surface properties of (b) and (d) (i.e., tall and bright). GM simulations have the ice sheet surface properties of (b) and (c) (i.e., tall and dark). WP simulations have the ice sheet surface properties of (a) and (d) (i.e., flat and bright).

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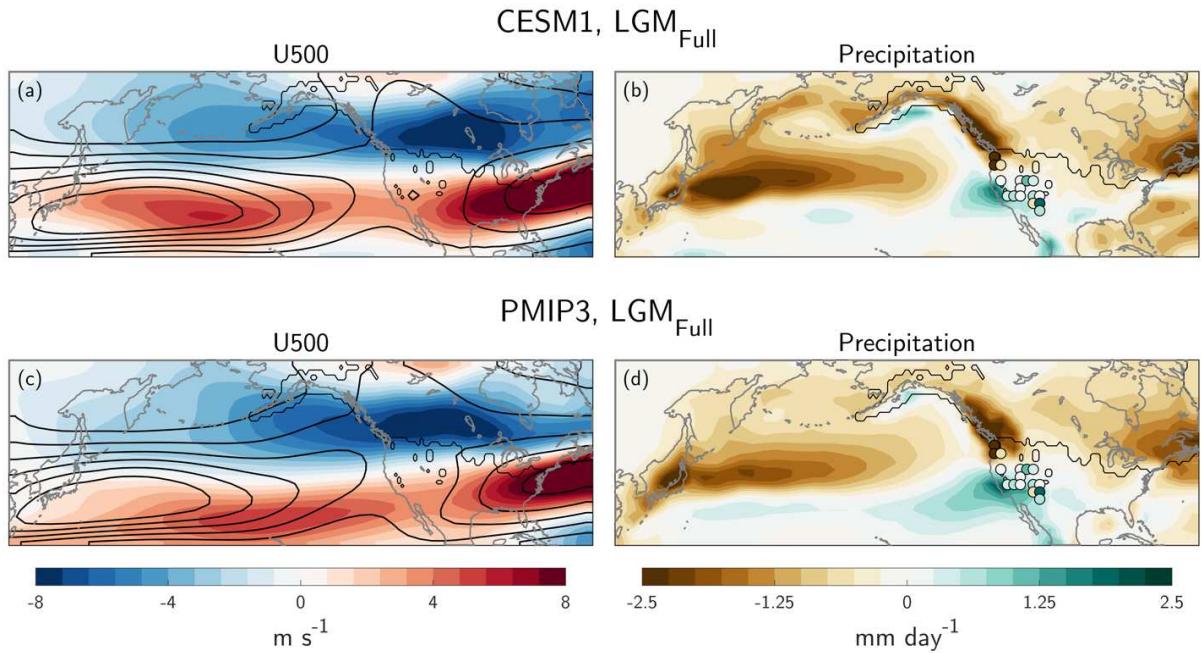
Figure 2 North Pacific climate at the Last Glacial Maximum. Difference maps of LGM$_{Full}$ minus CTL (a) 500 mb zonal wind (U500; m s^{-1}) and (b) precipitation (mm day^{-1}) averaged during boreal winter (December-February). Thick black contours in (a) show wintertime U500 climatology in the pre-industrial control (contour interval is 5 m s^{-1} with max value of 30 m s^{-1}). (c)-(d) As in (a)-(b), but for the ensemble mean of seven PMIP3 models forced with full LGM boundary conditions minus their respective pre-industrial control simulations. Thin black contour in each panel marks approximate ice sheet edge at the LGM. Filled circles in (b) and (d) show precipitation changes from proxy estimates (Table S2).
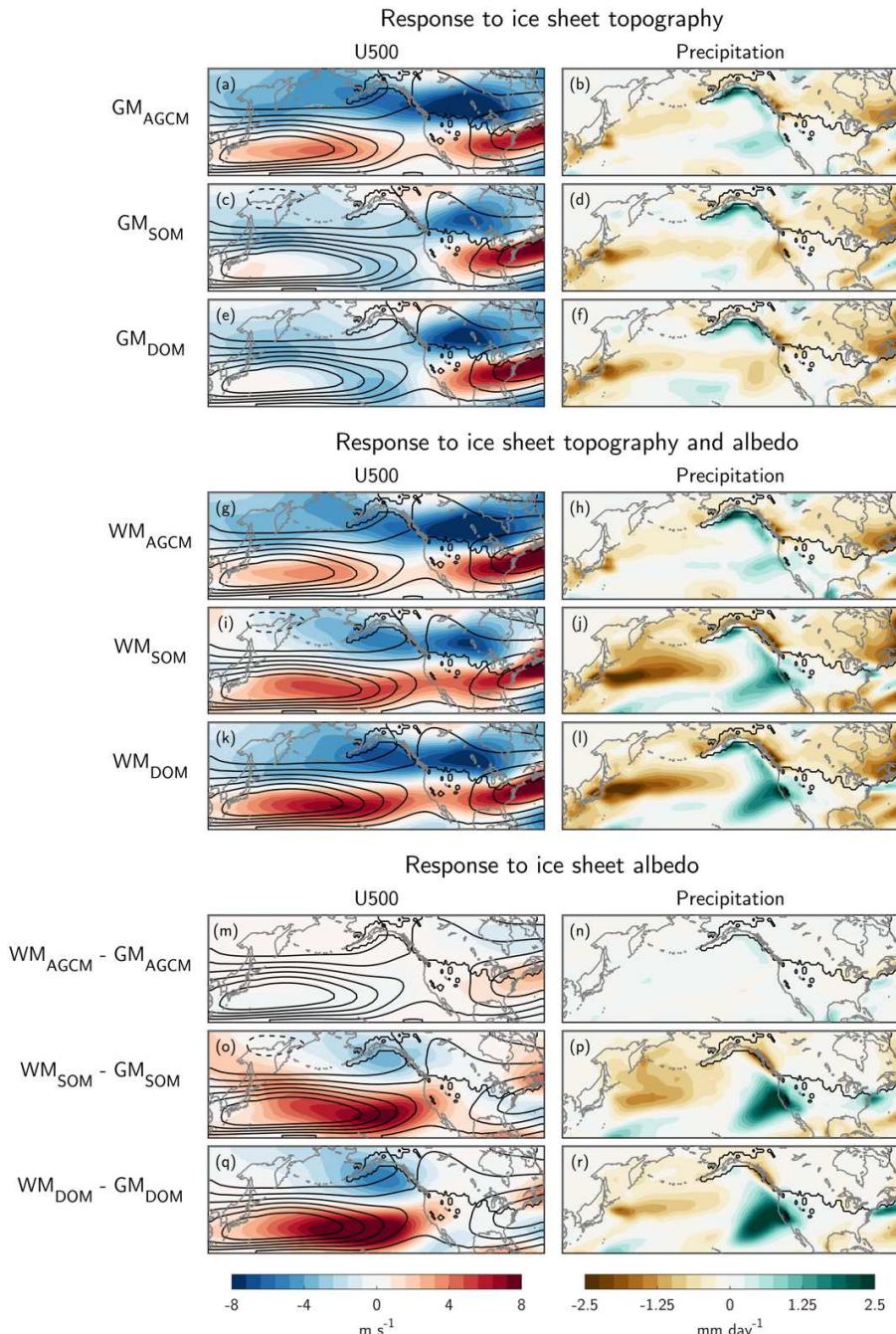
Figure 3 North Pacific boreal winter response of 500 mb zonal wind (U500) and precipitation in single forcing experiments: (a)–(f) GM−CTL, (g)-(i) WM−CTL, and (m)–(r) WM−GM in their respective ocean configurations. Thick black contours in show wintertime U500 climatology in the pre-industrial control (contour interval is 5 m s^{-1} with max value of 30 m s^{-1}).
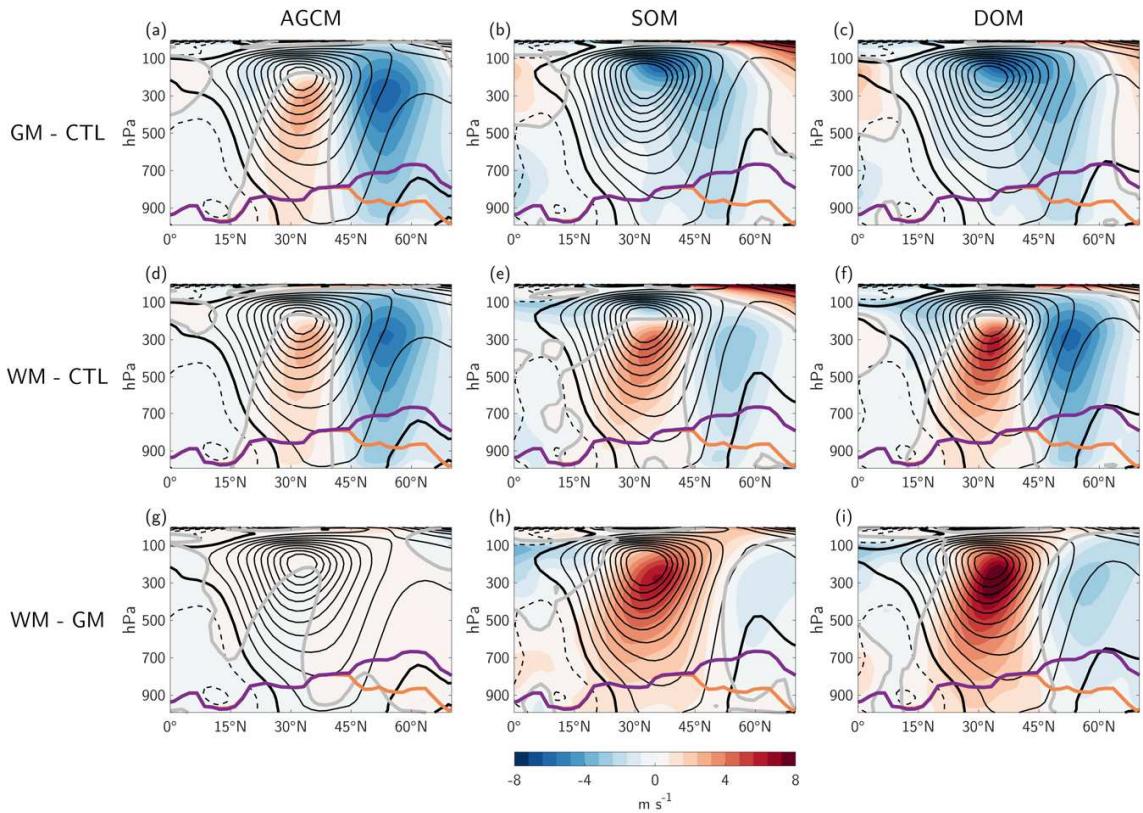
Figure 4 North Pacific zonal wind response. Latitude-height cross-sections of boreal winter (December-February) zonal wind anomalies (shading; m s^{-1}) zonally averaged over the Pacific basin (120˚E-110˚W). Anomalies are for (a)-(c) GM minus CTL, (d)-(f) WM minus CTL, and (g)-(i) WM minus GM in their respective (left column) AGCM, (middle column) SOM, and (right column) DOM ocean configurations. Purple and orange contours roughly outline the profile of North American topography at each latitude (i.e., as if looking westward from the Atlantic Ocean) for LGM and CTL, respectively. Black contours denote the winter CTL zonal wind climatology for this region. Solid (dashed) black contours denote positive (negative) values. Black contour interval is 5 m s^{-1} starting at 0 m s^{-1} (thick black contour). Thick gray contour marks the 0 m s^{-1} zonal wind anomaly contour.

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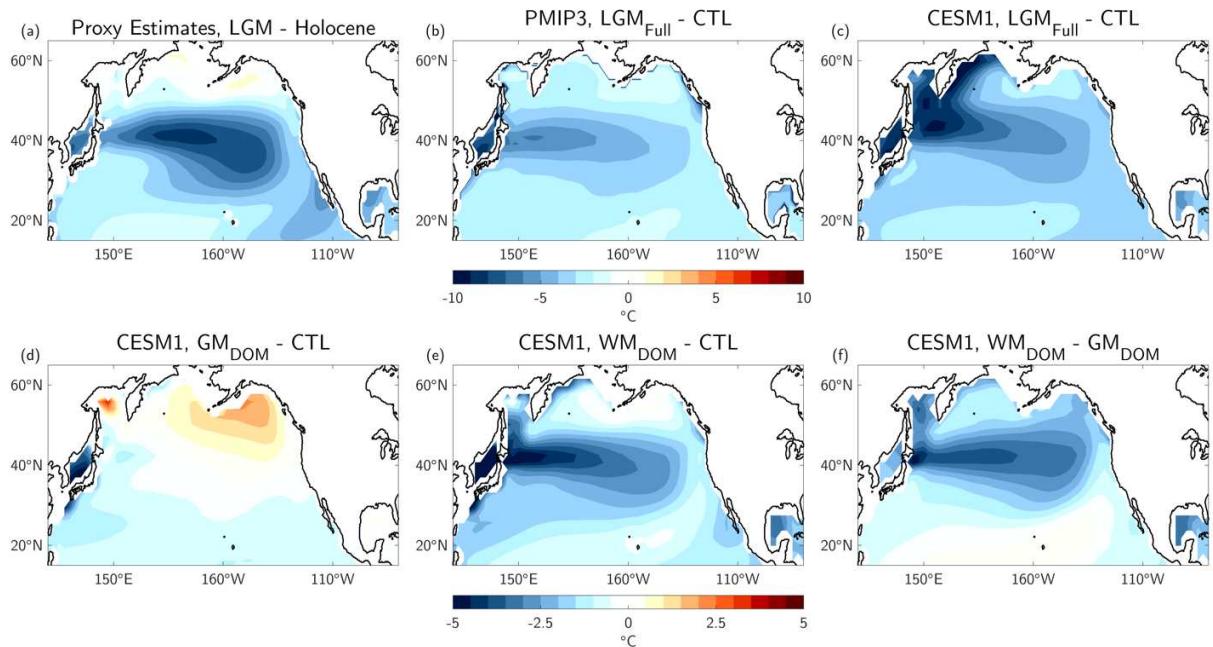
Figure 5 Simulated ocean temperature response. Difference maps of annual mean SSTs (˚C) in (a) LGM minus late Holocene data assimilated proxy estimates from Tierney et al. (2020), (b) the ensemble mean of seven PMIP3 models forced with full LGM boundary conditions minus their respective pre-industrial control simulations, (c) CESM1 LGM_Full minus CTL simulations, (d) CESM1 GM_DOM minus CTL, (e) CESM1 WM_DOM minus CTL, and (f) CESM1 WM_DOM minus GM_DOM.

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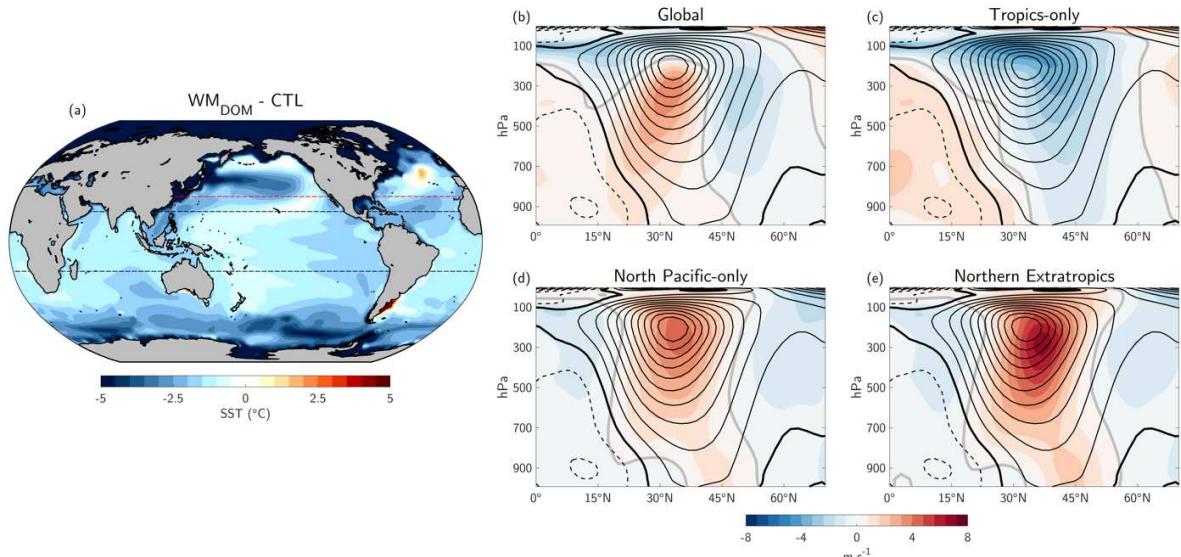
Figure 6 North Pacific zonal wind response in WM$_{SST}$. (a) Global SST anomalies (˚C) for WM$_{DOM}$ minus CTL averaged during December-February. (b)-(e) Latitude-height cross-section of boreal winter (December-February) zonal wind anomalies (shading; m s^{-1}) zonally averaged over the Pacific basin (120˚E-110˚W). Anomalies are from the WM$_{SST}$ CAM5 simulations forced with WM$_{DOM}$ SSTs (b) globally, (c) in the tropics-only (20˚S-20˚N; equatorward of dashed black lines), (d) in the North Pacific-only (>30˚N; poleward of black dashed line in North Pacific), and (e) in the Northern Hemisphere extratropics (>30˚N; poleward of the black dashed line at all longitudes). Black contours in (b)-(e) denote the winter WM$_{SST}$ control zonal wind climatology for this region. Solid (dashed) black contours denote positive (negative) values. Black contour interval is 5 m s^{-1}, with the thick black contour marking 0 m s^{-1}. Thick gray contour marks the 0 m s^{-1} zonal wind anomaly contour.
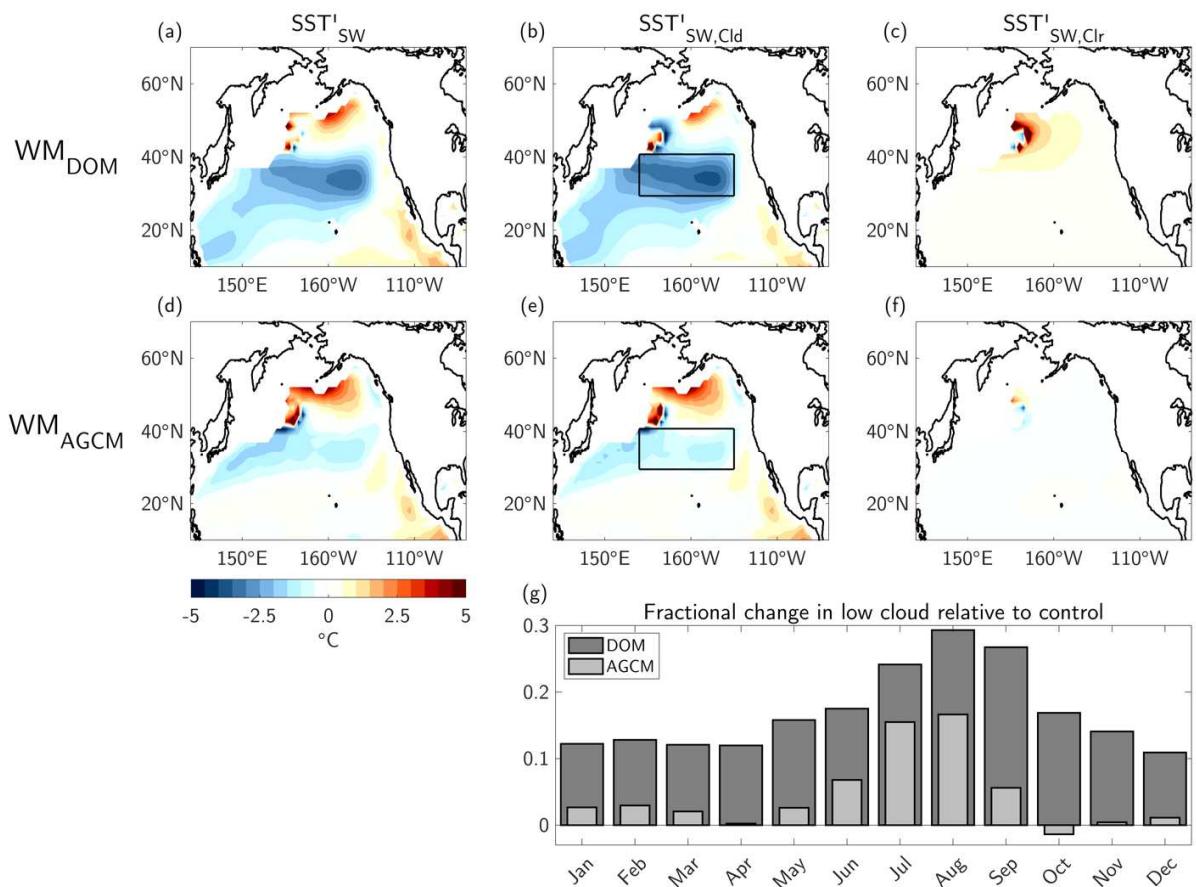
Figure 7 Impact of clouds on ocean temperatures. SST anomalies implied by changes in annual mean (a) net surface shortwave, (b) shortwave cloud radiative forcing, and (c) clear-sky surface shortwave in WM$_{DOM}$. (d)-(f) Same as (a)-(c), but for WM$_{AGCM}$. (g) Monthly fractional change (relative to CTL) in low cloud amount, averaged in black boxes in (b) and (e) for WM$_{DOM}$ (dark gray) and WM$_{AGCM}$ (light gray).
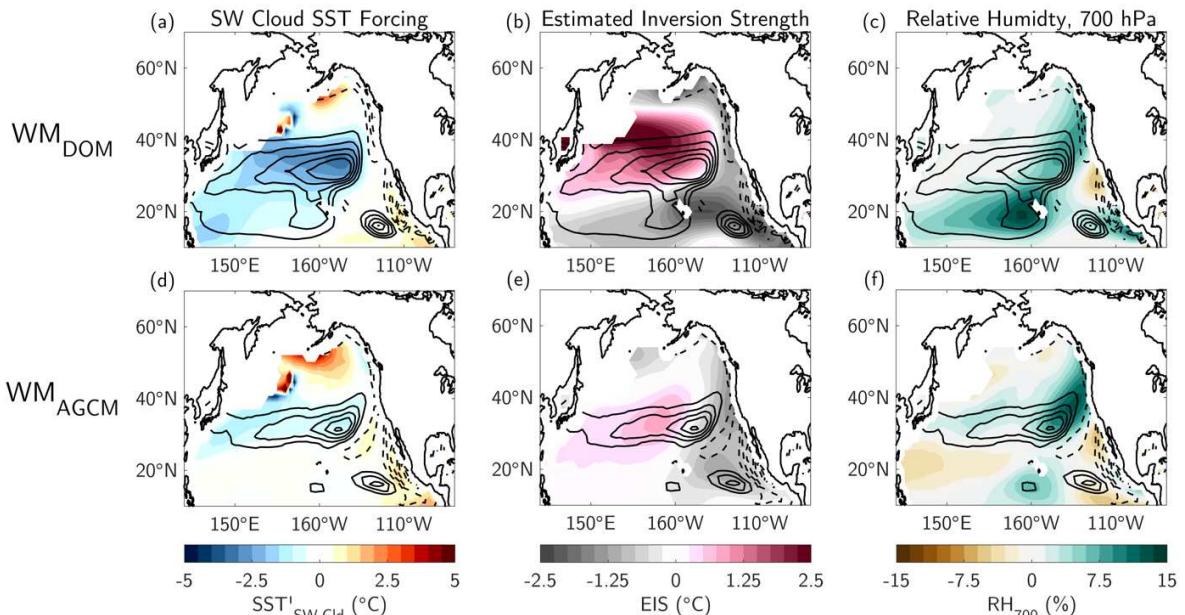
Figure 8 Low cloud controlling factors. (left column) SST anomalies implied by changes in annual mean shortwave cloud radiative forcing (shading; ˚C), (middle column) Estimated Inversion Strength (EIS; shading; ˚C) anomalies averaged during boreal summer (June-August), and (right column) relative humidity anomalies (shading; %) at 700 mb averaged during boreal summer in (a)-(c) WM$_{DOM}$ minus CTL and (d)-(f) WM$_{AGCM}$ minus CTL. Black contours in each panel depict the corresponding fractional change in summertime averaged low cloud fraction (e.g., WM$_{DOM}$ minus CTL over CTL). Solid (dashed) black contours denote positive (negative) values. Contour interval is 0.1 starting at 0.1. See Supplementary Materials Section S1.4 for more details on the cloud parameters analyzed here.
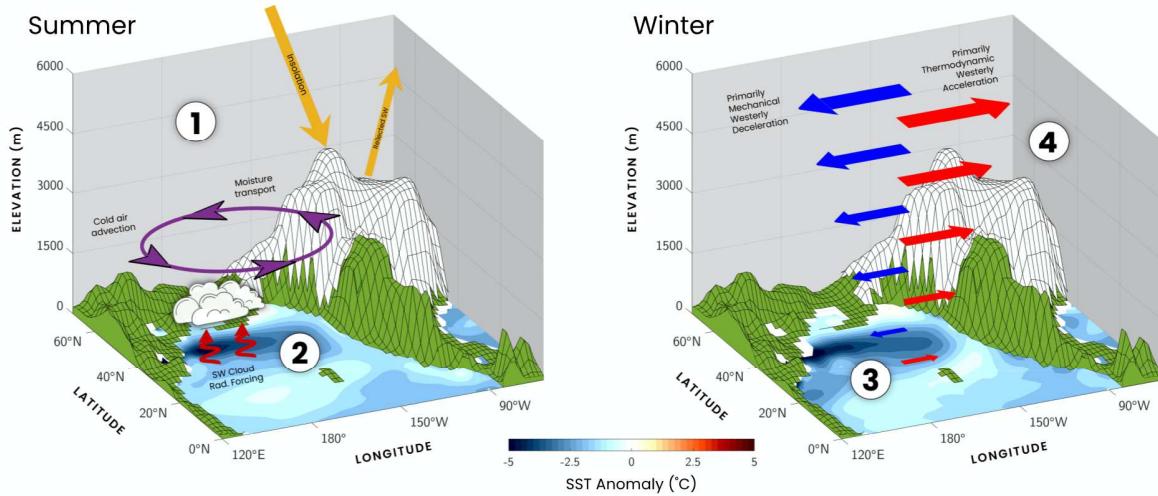
Figure 9 Schematic summary of the North Pacific coupled climate response to tall/bright North American ice sheets during the LGM. (1) A direct stationary wave response to the tall/bright ice sheets during boreal summer produces a low-level cyclonic circulation over the North Pacific. These circulation anomalies drive cold air advection in the western North Pacific and increased water vapor transport in the eastern North Pacific. (2) Enhanced stability coupled with increased atmospheric moisture leads to more marine stratocumulus clouds, which reflect sunlight and cool the North Pacific Ocean near ~40°N. (3) Low cloud-SST feedback and ocean dynamical adjustments amplify the SST and cloud changes, persisting them into boreal winter. (4) The enhanced North Pacific meridional temperature gradient accelerates the westerlies on the equatorward flank of the jet stream, while mechanical interactions with ice sheet topography decelerate the flow on the poleward flank. Overall, this leads to a southward shift of both the North Pacific jet stream and North American west coast hydroclimate during the LGM.