Uppermost mantle velocity beneath the Mid-Atlantic Ridge and transform

faults in the Equatorial Atlantic Ocean

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Abstract

Seismic rays that travel just below the Moho discontinuity provide useful insights into the thermal and compositional properties of the upper-mantle, and can be detected as Pn phases from regional earthquakes. Such phases are routinely identified in the continents. In the oceans, however, detection of Pn phases is limited due to the lack of long-term instrument deployments. Here, we present estimates of upper-mantle velocity in the Equatorial Atlantic Ocean from Pn arrivals beneath, and flanking, the Mid-Atlantic Ridge, and across several transform faults. We analyzed waveforms from 50 regional earthquakes with magnitude $M_W > 3.5$ recorded over 12 months in 2012–2013 by five autonomous hydrophones and a single broadband seismic station located on the St. Peter and St. Paul archipelago. The resulting catalog of 152 ray paths allows us to resolve spatial variations in upper-mantle velocities, which are consistent with estimates from nearby wide-angle seismic experiments. We find relatively fast upper-mantle velocities near to the St. Paul transform system (~8.4 km s⁻¹), compared to slower velocities parallel to the ridge axis (~7.7 km s⁻¹). Hence, this method is able to resolve ridge-transform scale velocity variations. Ray paths in lithosphere younger than 10 Myr show velocities with mean of 7.9 ± 0.5 km s⁻¹, which is slightly slower than those sampling lithosphere older than 20 Myr (8.1 km \pm 0.3 s⁻¹). There is no apparent systematic relationship between velocity and ray azimuth, nor any correlation between uppermantle Pn velocity and shear wave speeds from the global SL2013sv model at depths < 150 km.

- Our results demonstrate that data from long-term deployments of autonomous hydrophones can
- be used to obtain rare and insightful estimates of uppermost mantle velocities over hundreds of
- 26 kilometers, in otherwise inaccessible parts of the deep oceans.

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Introduction

Seismic velocity measurements provide a useful proxy for investigating spatial variations in uppermantle properties, such as temperature and anisotropy, with implications for melt supply and mantle heterogeneity (e.g. Lin and Phipps Morgan, 1992; Dunn et al., 2005). These measurements are relatively straightforward to obtain on the continents (e.g. Chulick and Mooney, 2002; Chulick et al., 2013). However, it remains challenging and expensive to measure upper-mantle seismic velocity in the deep ocean basins, due to their remote location and challenges of deploying longterm instruments on the seafloor. Pn phases are critically refracted rays that are guided through an interface in the lithospheric mantle (Linehan, 1940; Brandsdottir and Menke, 1997), and can be recorded by seafloor-moored hydrophones, which in turn can be used to estimate upper-mantle seismic velocity (Walker et al., 1983; Slack et al., 1999). At the Mid-Atlantic Ridge (MAR) from 10°N to 35°N, Pn arrivals from 48 individual ray paths were identified using hydrophone-recorded data and used to investigate upper-mantle velocities, giving a mean velocity of 8.0 ± 0.1 km s⁻¹ (Dziak et al., 2004). This velocity estimate was higher than that from nearby active source seismic experiments along the ridge axis (7.5–7.9 km s⁻¹; Canales et al., 2000), probably due to the effects of younger and thinner oceanic lithosphere being sampled by the refraction profiles, and the effects of averaging velocities across all rays. Despite these advances, upper-mantle velocities across deep ocean basins remain poorly constrained, and the potential for Pn travel times from hydrophone-recorded data to resolve spatial variations in upper-mantle velocity has not yet been

sufficiently tested.

Here, we use *Pn* arrivals from regional earthquakes to constrain upper-mantle velocity in the remote Equatorial Atlantic Ocean. Arrivals were recorded by a combination of five moored hydrophones and a single seismograph station installed on the St. Peter and St. Paul islets, giving a total of 152 ray paths that sample mantle conditions both on- and off-axis, and across the St Paul transform system. Our study is coincident with several mantle velocity estimates from a wide-angle seismic experiment (Le Pichon *et al.*, 1965), and hence has the opportunity to validate spatial variations in velocity revealed by groups of similar ray paths.

Equatorial Atlantic Ocean

In our study region of the Equatorial Atlantic Ocean (10°N–5°S and 34W-21W), the MAR is offset by some of the longest transform faults on Earth, including the Strakhov, St Paul, and Romanche transforms (Figure 1). The St Paul transform system consists of four transform faults and three intra-transform ridge segments that accommodate an offset of 630 km. The northwest transform fault is currently undergoing transpression, giving rise to the ~200 km-long and ~30 km-wide Atobá ridge (Maia *et al.*, 2016), and also uplift of 1.5mm yr⁻¹ at the St Peter and St Paul islets (Campos *et al.*, 2010; Maia *et al.*, 2016). Other transforms in the system do not host topographic highs or an island related to transpression, and hence presumably are not experiencing uplift. In the three intervening spreading segments, seafloor spreading is slow, at ~20 mm yr⁻¹ average half rate (DeMets *et al.*, 2010). Faulting plays an important role in crustal accretion, and seismicity rates are relatively high, providing a useful tool to investigate the properties of the crust and upper mantle, as well as deformation at long-offset strike-slip systems (e.g. Francis *et al.*, 1978; Abercrombie and Ekstrom, 2001; de Melo and do Nascimento, 2018).

Methods

Hydroacoustic Waves

Hydrophones moored in the Sound Fixing and Ranging (SOFAR) channel routinely record T-phases, and also acoustic-converted *P*-waves from regional and teleseismic earthquakes (e.g. Dziak *et al.*, 1997, 2004; Slack *et al.*, 1999; Sohn and Hildebrand, 2001). Seismic waves from regional earthquakes are expected to propagate along the Moho discontinuity as *Pn* phases, and convert to an acoustic phase on the seafloor near to the hydrophone mooring, where they can be detected (e.g. Dziak *et al.*, 2004). Identification of acoustic-converted *Pn* phases on moored autonomous hydrophone records provides a simple and effective method to estimate upper-mantle velocities in remote locations. Compared to long-term broadband ocean bottom seismograph deployments, or wide-angle active source seismic experiments, arrays of these moored instruments are relatively inexpensive to build (e.g. Fox *et al.*, 2001).

Waveform Data

We analyzed *Pn* arrivals in waveform data recorded by a combination of five moored autonomous hydrophones and one land-based seismograph (Figure 2). The five autonomous hydrophone instruments were deployed during two separate experiments: stations EA2 and EA8 were part of the Equatorial Atlantic (EA) array (Smith et al., 2012). Data were recorded at 16-bit resolution and a sampling rate of 250 Hz; for further details on these hydrophone instruments see Fox *et al.* (2001). Hydrophones H2, H4, H5 were deployed during the COLd Mantle Exhumation and Intra-transform Accretion experiment (COLMEIA; Maia *et al.*, 2014, 2016), and recorded data at 24 bit-resolution with a sampling rate of 240 Hz; for further instrument details see D'Eu

et al. (2012). We also used waveform data recorded by a three-component broadband seismograph installed at the Saint Peter and Saint Paul Archipelago Scientific Station on the Belmonte islet (ASPSP; de Melo and do Nascimento, 2018). This station is operated by the Seismological Laboratory of Federal University of Rio Grande do Norte in cooperation with the Brazilian Navy. The sparse distribution and mixed instrument types we used means that data coverage is uneven, as shown in Figure 1b. Waveform data were examined for the time period from July 2012 to July 2013, using a combination of the five hydrophones and the ASPSP seismograph, with recording intervals dictated by technical challenges and vessel schedules.

Prior to manually picking *Pn* arrivals, we applied a 6–20 Hz Butterworth bandpass filter to the hydrophone data in order to suppress unwanted noise. A bandpass filter with range 4–12 Hz was applied prior to picking arrivals from the ASPSP seismograph, to suppress additional microseism noise due to its island location.

Pn analysis

We used regional *P*-wave arrivals recorded by the five autonomous hydrophones and by the ASPSP seismograph to identify *Pn* phases that have travelled within the uppermost mantle in the Equatorial Atlantic Ocean. During the 12 month-long observation period, *P*-wave arrivals from 50 regional earthquakes were recorded. Based upon origin time, events were manually associated with earthquakes in the International Seismological Center Bulletin (ISC; www.isc.ac.uk), yielding hypocenter locations, origin times, and magnitudes ranging from 3.5 to 5.4 M_W. Earthquakes dominantly arise due to strike-slip faulting along the Strakhov, St. Paul, and Romanche transform faults, with additional events due to extension along the intervening spreading ridge segments (Figure 2a). Three example arrivals from these events are shown in

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Figures 3 and 4, highlighting the typical response to strike-slip and normal faulting earthquakes, which range in magnitude from 4.6 to 5.3 M_W. Typical P-arrivals have sufficient signal-to-noise ratio for picking (noted in Figures 3 and 4), and linear move-out, consistent with upper mantle velocity, is evident across the hydrophone array stations due to wave propagation along the crustmantle interface (see common-receiver plots in Supplementary Figures S1–S6). P-arrivals are easily distinguished from T-phase arrivals, which arrive much later than P-arrivals, are emergent in character, and are higher in amplitude than P-arrivals (see hydrophone H5 in Figure 4). The catalog of detected events is given in Table S1. In order to further test whether the detected arrivals were Pn phases, we compared the observed travel times to those predicted by the global IASP91 velocity model (Kennett and Engdahl, 1991). For each source-receiver ray path, we calculated the predicted Pn arrival time using IASP91, with the addition of a station-dependent delay to account for the propagation time from seafloor to hydrophone. This delay (1.2–2.5 s) was estimated using the hydrophone mooring cable length at each station, and the local water sound velocity estimated from the Global Ocean Sound Speed Profile Library (Barlow, 2019), values are given in Table 1. The predicted Pn arrival times differ from the observed Pn arrivals by 5–10 s (Figures 3 and 4), however this is not surprising since the IASP91 model contains a crustal layer that is much thicker (30 km) than that expected in the oceans (~6 km). Hence, the differences in observed and predicted Pn arrival time are probably dominated by this additional crustal layer thickness in the velocity model, plus earthquake location and origin time uncertainties. Although these differences are evident, the waveform character and linear move-out velocity means that we are confident in our identification of these high signal-to-noise phases as *Pn* arrivals.

ISC origin times were subtracted from the Pn arrival times to obtain travel times for each ray path (i.e. each event-station pair), and the Pn velocity estimated by dividing this travel time by the great-circle (i.e. shortest) epicentral distance from event to station. We account for oceanic crust by subtracting ray path distances and travel times for the portion of the path that travels through the crust, assuming that all events occurred at 10 km depth (ISC catalog), and that crustal thickness is uniformly 6 km (Christeson $et\ al.$, 2019). For each station, we then calculate the distance and travel time for the portion of the ray path that extends from an earthquake in the crust to the Moho, and back from the Moho through the entire oceanic crustal thickness. Details of these corrections for each station are given in Table 1.

Our approach yielded a total of 152 ray paths and hence 152 separate *Pn* velocity estimates from the catalog of 50 regional earthquakes (Figure 5). Although epicentral distance ranges greatly from 32 km to ~1095 km, all 50 events were detected at all available stations, implying that the detection threshold of the combined hydrophones and ASPSP station is at least as low as M_W 3.5. Since most stations were located either near to, or to the north of, the St Paul fracture zone, our ray path coverage is more comprehensive in the northern part of the study area. Ray paths sampling upper-mantle velocities to the south of the St Paul fracture zone are restricted to events detected by hydrophone EA8, and those originating from four earthquakes located at the eastern end of the Romanche transform fault (Figure 5).

Pn velocity uncertainty

The two most significant potential sources of error are in the hypocenter locations of events in the ISC Catalog, and in the Pn arrival picks. These two sources of error directly influence ray path length and travel time, respectively. We estimated hypocenter location (and hence epicentral

distance) error to be \pm 10 km, based upon ISC catalogue location uncertainty, and typical error in global earthquake location (Lohman and Simons, 2005; Weston *et al.*, 2012). Arrival pick (and hence also travel time) errors were estimated based upon qualitative station-dependent pick uncertainties, defined as \pm 0.5 s for EA2 and EA8; \pm 1.0 s for H2, H4 and H5; and \pm 0.3 s for ASPSP. The total uncertainty in our velocity estimate, δv , was estimated by formally propagating the errors in epicentral distance, d, and travel time, t, as follows

$$\delta v = v \sqrt{\left(\frac{\delta d}{d}\right)^2 + \left(\frac{\delta t}{t}\right)^2},$$

where the d and δd are epicentral distance and its respective error, and t and δt are travel time and its respective error (e.g. Taylor, 1997).

Although receiver location uncertainty is negligible for the land station ASPSP (located via the Global Positioning System), there is potential location uncertainty for the moored hydrophones in our network. Moored hydrophone locations were obtained by acoustic triangulation between the mooring acoustic release and the deployment vessel soon after the moorings settled on the seafloor, within error of several meters. In order to account for the possibility of abnormally strong current motion, each instrument is fitted with a pressure and temperature logger below the floatation package, so that any significant hydrophone depth changes would be recorded (e.g. Fox et al., 2001). Significant depth changes were not detected during our experiment, and thus we assume that the hydrophone location was constant during data collection, and hence uncertainty is relatively minor.

Results

Pn velocities

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The catalog of 152 Pn ray paths (Figure 5b) and travel times (Figure 6) indicates uppermantle velocities that vary considerably across the study area, with estimates ranging between 7.2 and 11.1 km s⁻¹, with uncertainties ranging from 0.1 to 1.9 km s⁻¹ (Table S2). At the center of the study area there appears to be a longitudinal variation in Pn velocity, with events originating near the St Paul transform system, and sampling adjacent lithosphere, having higher velocities than those from the adjacent spreading centers (Figure 5a). The best constrained estimate for sub-axis, ridge-parallel mantle velocity comes from ray paths that sample the portion of the spreading axis between the Strakhov fracture zone and stations near the St Paul fracture zone (H2, H5 and ASPSP). Here, Pn travel times consistently imply relatively slow velocities, with a mean of 7.7 km s⁻¹. Slightly faster velocities ranging between 7.8 and 8.2 km s⁻¹ are indicated by ray paths between hydrophone EA2 and the Strakhov fracture zone, oriented roughly parallel to a plate spreading flowline. Ray paths oriented southwest-northeast (azimuth ~060°), i.e. oblique to the spreading direction, between events on the St Paul fracture zone and detected at hydrophone EA2, have some of the fastest mantle velocities (between 7.6 and 8.5 km s⁻¹) compared to other rays sampling areas unaffected by fracture zones. Velocity estimates in the vicinity of the St Paul fracture zone itself (from transform faulting events detected by hydrophones H2, H4 and H5, and ASPSP) show considerable variation, ranging from 8.0 to 9.1 km s⁻¹ and a mean of 8.4 km s⁻¹, and little apparent spatial consistency. Among these events, we encountered one of the fastest Pn velocities (9.0 km \pm 0.2 s⁻¹) in this study, for a ray path oriented roughly parallel to the St Paul transform fault (ray azimuth ~105°) between an event near the St Paul islets and detected by hydrophone H4.

South of the St Paul fracture zone, ray paths from events detected by hydrophone EA8 showed considerable variation in upper-mantle velocity, which range from 7.2 to 9.0 km s⁻¹. Ray paths originating from the spreading axis north of the St Paul transform fault and trending $\sim 170^{\circ}$

towards EA8, have velocities of 7.3–8.1 km s⁻¹, while ray paths from the St Paul transform fault trending \sim 185° towards EA8 have consistently faster velocities of 7.6–9.1 km s⁻¹.

Only 12 ray paths sampling the upper-mantle parallel and adjacent to the spreading axis between the southern extent of the St Paul transform fault and the Romanche transform fault are available. This relatively poor coverage in ray paths in this area hinders our interpretation, where velocities range from 7.2 to 8.3 km s⁻¹.

Discussion

Upper-mantle velocity structure

In general, rays originating from the St Paul transform system are faster than those originating from active spreading centers to the east and west (Figure 5a), probably due to cooler conditions at the Moho along the transform, giving higher velocities. Our estimates of uppermantle *Pn* velocities broadly agree (within error) with *Pn* velocities from radially stratified velocity models such as PREM (Dziewonski and Anderson, 1981) and IASP91 (Figure 6; Kennett and Engdahl, 1991). *Pn* velocity estimates from this study are also consistent with mantle velocity estimates from a series of reversed wide-angle refraction seismic profiles (i.e. with multiple shot points giving overlapping coverage) collected in the equatorial Atlantic during R/V *Atlantis* cruise A180 (Figure 5b; Le Pichon *et al.*, 1965). The modal difference in velocity between refraction profiles from Le Pichon *et al.* (1965) and all intersecting ray paths is 0.2 km s⁻¹ (see histogram in Figure 5c), although our *Pn* velocity estimates are typically slower than those reported by Le Pichon *et al.* (1965), with a maximum disagreement of 1.2 km s⁻¹. A mantle velocity of 8.30 km s⁻¹ was reported along profile A180-48, which is 283 km-long, and crosses the eastern side of the St Paul transform fault (near ~26.3°W), trending northeast-southwest (Figure 5b). This velocity

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is consistent with that inferred from Pn ray paths with a similar orientation, originating from earthquakes on the St Paul transform fault that were detected by hydrophone EA8. Ray paths that intersect profile A180-48 (at angles either perpendicular or oblique to the trend of the refraction profile) typically indicate slower upper-mantle velocities, ranging from 7.3 to 8.1 km s⁻¹, with the exception of one anomalous ray path oriented parallel with the St Paul transform fault with a velocity of 9.0 km s⁻¹. Refraction profiles A180-40 and -42 are oriented roughly east-west, are located ~100 km north of the Romanche transform fault, and have velocities of 8.03 and 8.49 km s⁻¹, respectively. Although there are only four Pn ray paths near to these profiles, with nearperpendicular orientation, they indicate velocities ranging from 7.6 to 8.2 km s⁻¹, and hence are in Our results are also in close agreement with a velocity estimate of 8 km s⁻¹ broad agreement. from an active source experiment near 18°W roughly perpendicular to the St Paul fracture zone, which at this longitude separates 40 Myr old crust in the south from 70 Myr old crust in the north (Growe et al., 2019). The generally good agreement between upper-mantle velocities from the refraction profiles and our Pn arrivals validates our results, and implies that spatial trends observed in the study area are likely to be real. Elsewhere along the MAR, between 10° to 40°N, a mean upper-mantle velocity of 8.0 km

Elsewhere along the MAR, between 10° to 40° N, a mean upper-mantle velocity of 8.0 km \pm 0.1 km s⁻¹ was estimated using a similar method to this study with Pn arrivals detected by an array of autonomous hydrophones (Dziak *et al.*, 2004). In common with this study, ray paths used by Dziak *et al.*, (2004) often crossed the ridge axis, spanned a series of fracture zones, and extended onto older crust, which may explain the close agreement in results. This result suggests that off-axis and on-axis upper mantle characteristics are similar in the north and equatorial Atlantic Ocean.

Near the Oceanographer transform fault (~35°N), two-dimensional tomographic inversion of wide-angle seismic refraction data located along the ridge axis suggest velocities of 7.4–7.8 km

s⁻¹ (Canales *et al.*, 2000; Hooft *et al.*, 2000). These results agree closely with our estimates of *Pn* velocity from rays sampling on-axis upper-mantle to the north of the St Paul transform fault (Figure 5b).

Upper-mantle velocity and plate age

Due to upwelling of hot material, seismic velocities in the upper-mantle near to the ridge axis are expected to be lower than in off-axis areas (e.g. Turcotte and Schubert, 2002). Following the removal of minor gridding artifacts associated with fracture zone traces, we used the global crustal age model (Müller et al., 2008) to assign a mean age along each ray path, for comparison with Pn velocity (Figure 7a).

Ray paths along lithosphere younger than 10 Ma show a wide range of velocities, with a mean of 7.9 km s⁻¹ and standard deviation of 0.5 km s⁻¹. Twenty of the ray paths yield velocities less than 7.5 km s⁻¹. *Pn* velocities for ray paths along lithosphere older than 20 Ma are typically slightly faster, with a mean of 8.1 km s⁻¹ and standard deviation of 0.3 km s⁻¹, while only two estimates give velocities slower than 7.5 km s⁻¹ (Figure 7a). Most rays cover a wide range of crustal ages, so this geometry, and our averaging approach, may smear the possible effects of lithospheric ageing. The lack of rays travelling exclusively via older lithosphere may also obscure any progressive trend between upper-mantle velocity and crustal age. However, the tendency toward the inclusion of slower velocities in younger crust (Figure 7) reflects the expected variation with respect to the zone of axial upwelling.

Azimuthal Seismic Anisotropy

Laboratory experiments have shown that the mantle can experience significant shear strain during corner flow at the ridge axis, leaving an anisotropic fabric in the lithospheric mantle as minerals (e.g. olivine) are aligned into a lattice preferred orientation (LPO; e.g. Zhang and Karato, 1995; Nicolas and Christensen, 2011). Anisotropy consistent with a LPO formed by two-dimensional mantle flow has been measured at some locations in the oceanic upper mantle (e.g. Raitt *et al.*, 1969; Lin *et al.*, 2016), however the strength of anisotropy varies widely, and debate remains about its origins (e.g. Mark *et al.*, 2019). Since isochrons in this region are fairly uniform (Figure 7), V_{Pn} anisotropy could be expected parallel to paleo-relative plate motion, although this assumption has been shown to not apply everywhere (VanderBeek and Toomey, 2017).

We investigated the dependence of mantle velocity with azimuth, and use epicentral distance as a proxy for depth of mantle penetration to group rays (Figure 7b). No discernable pattern is evident in rays grouped by epicentral distance, including those expected to sample deepest in the mantle with epicentral distances > 700 km (blue lines in Figure 7c). Removing rays with V_{Pn} error > 0.4 km s⁻¹ also does not resolve any azimuthal dependence (Figure 7d), nor does separating rays by mean crustal age (Figures 7e and 7f).

The apparent lack of such azimuthal dependence could be due to several reasons. First, azimuthal dependence may be too subtle to be resolved by our V_{Pn} estimates, given the uncertainties in hypocenter location and crustal thickness discussed above. Second, the slow spreading rate of the MAR 22 mm/yr total rate) may result in a thickened lithosphere that is dominantly cooled by conduction, thus inhibiting corner flow (e.g. Sleep, 1975). As a result, deformation could be accommodated by faulting at depths of 5–10 km beneath the Moho, reducing the viscous strain in the mantle at these depths, and suppressing the anisotropy recorded in the mantle (e.g. Ribe, 1989). Observations of weaker or anomalous anisotropy elsewhere in the

Atlantic Ocean are consistent with our findings (e.g. Gaherty et al., 2004; Dunn et al., 2005). Third, complex, three-dimensional upwelling patterns near the ridge axis could result in anisotropy on relatively short wavelengths (Lin and Phipps Morgan, 1992), which would be smeared along our relatively long ray paths, and hence not be resolved.

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Pn and surface wave velocity

To explore the relationship between V_{Pn} and the thermal structure of the asthenospheric upper-mantle, we compared our velocity estimates with a global, vertically polarized shear speed model SL2013sv (Schaeffer and Lebedev, 2013). Our objective is to consider our observations of uppermost mantle properties in the context of deeper mantle anomalies. We do not aim to validate our V_{Pn} estimates via this comparison. This model was chosen because it is particularly sensitive to anomalies within the upper-mantle, and hence provides a window into the upper mantle structure directly beneath our Pn ray paths. We extracted values of vertically polarized tomographic shear velocity anomaly (%dVs) at 100 km intervals along each ray path from slices through the SL2013sv model at depths of 25, 50, 75 and 150 km. We then calculated the mean %dVs along each ray path, at each depth interval (Figure 8). At 25 and 50 km depths, the effects of the ridge axis are evident, with faster velocities associated with ray paths travelling offaxis (detected by EA2 and EA8), and hence not sampling the relatively slow axial region (Figures 8a and 8b). This effect is less pronounced at 75 km depth (Figure 8c), and is not apparent at 150 km depth, which presumably reflects sub-plate velocities. The lack of correlation between SL2013sv and Pn velocities at 150 km suggests that our V_{Pn} estimates, sensitive to the velocity structure directly beneath the Moho, do not record deeper, larger-scale sub-plate (i.e.

asthenospheric) processes and anomalies. Hence our observed V_{Pn} variability may instead arise due to local variations in melt supply, lithospheric thickness, or faulting.

Conclusions

We used a network of five autonomous hydrophones combined with a single island seismograph to detect Pn arrivals from regional earthquakes in the Equatorial Atlantic Ocean over a period of 12 months between 2012 and 2013. We estimated upper-mantle velocities using the travel times of 152 Pn arrivals, and find that our estimates are in good agreement with independent estimates of upper-mantle velocities from nearby seismic refraction experiments.

We find that the upper-mantle near the St Paul transform system has consistently fast velocities (>8 km s⁻¹), compared to relatively slow velocities (\sim 7.5 km s⁻¹) in the adjacent Mid-Atlantic Ridge spreading segments northwest of the transform. This spatial pattern is consistent with the notion that Pn ray paths sample slower velocities in areas of crustal accretion near the ridge axis, and faster velocities in areas of transforms which are in general, cooler, despite the presence of intra-transform spreading segments. We do not resolve any dependence between V_{Pn} and azimuth, which could either be due to observational uncertainty, or, due to the combined effects of thickened lithosphere and more complex mantle upwelling patterns. We also do not find any correlation between V_{Pn} velocities and vertically polarized shear speed from the global SL2013sv model, indicating that our method is not sensitive to properties of the asthenosphere. The good agreement between our results and those from seismic refraction experiments demonstrates that the relatively simple method of using sparse arrays of autonomous hydrophones to detect Pn arrivals can be used to obtain accurate estimates of upper-mantle velocities. Hence, this method

provides a useful complement to deployments of other seafloor instruments such as ocean bottom seismographs, in remote areas where direct observations are typically elusive.

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Figures

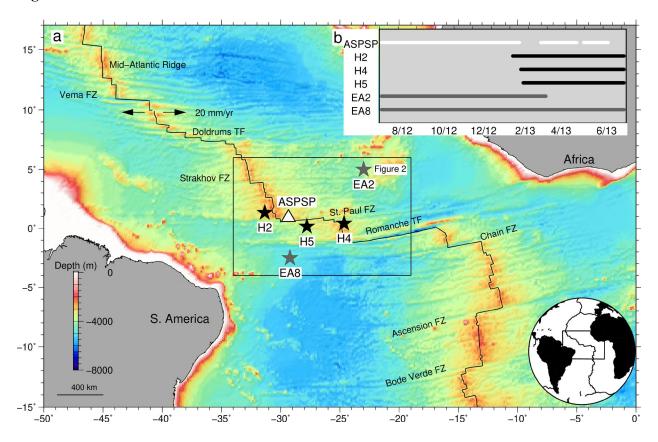


Figure 1. a) Regional bathymetric map of equatorial Atlantic ocean. White triangle shows ASPSP seismograph station, located on St. Peter and St. Paul islets; black/gray stars are COLMEIA / EA hydrophone networks, respectively (Smith *et al.*, 2012; Maia *et al.*, 2014); black line is Mid-Atlantic Ridge, with selected transforms and half-spreading rate noted (arrows). Black box shows location of Figure 2. b) Bars show operation intervals of instruments for each network: ASPSP (white); COLMEIA (black) and EA (gray).

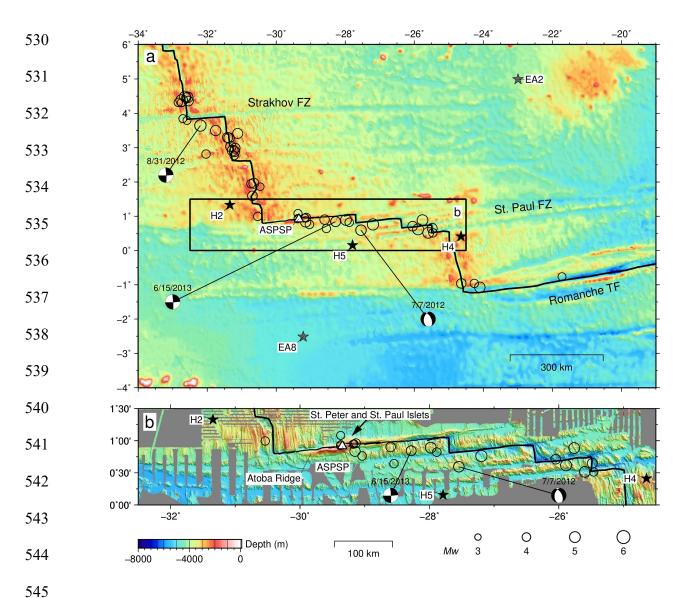


Figure 2. a) Bathymetric map of equatorial Atlantic ocean. Black box shows location of (b). Circles are earthquakes used in *Pn* analysis, scaled by M_W; triangle shows ASPSP station; black/gray stars are COLMEIA / EA hydrophone networks, respectively (Maia *et al.*, 2013; Smith *et al.*, 2012); black line is Mid-Atlantic Ridge, with selected transforms labeled; beach-balls are centroid moment tensors for three exemplar earthquakes (Ekström *et al.*, 2012), waveforms shown in Figures 4 and 5. b) Bathymetric map showing details of St Peter and St Paul fracture zone (from Udintsev *et al.*, 1996; Gasperini *et al.*, 1997; Maia *et al.*, 2016).

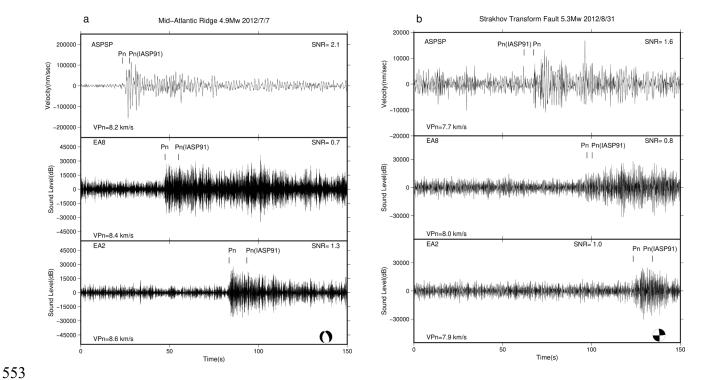


Figure 3. Example waveforms recorded by the ASPSP seismograph and EA array hydrophones, with 4–12 Hz and 6–20 Hz Butterworth filter applied, respectively. **a)** M_w 4.9 normal faulting event on 7^{th} July 2012, located on the St Paul transform fault at 27.5°W. Picked Pn arrivals, and Pn arrivals predicted by IASP91 model are marked; beach-balls are centroid moment tensors (Ekström et al., 2012); V_{Pn} and signal to noise ratio (SNR) noted for each station (this study), SNR calculated in a window 1.5 s either side of Pn arrival. **b)** M_W 5.3 strike-slip event on 31st August 2012, located on Strakhov transform fault near 32.5°W.

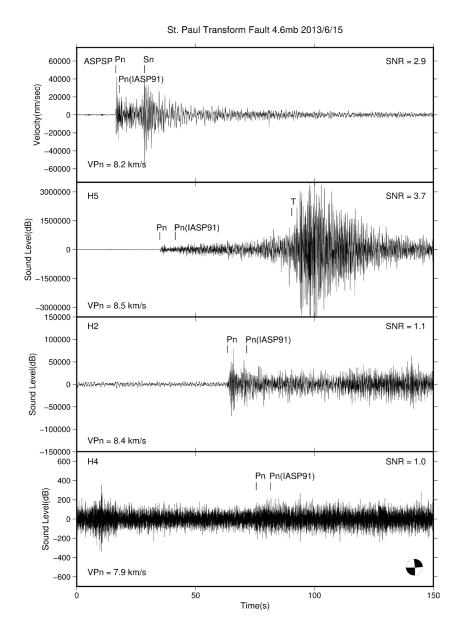
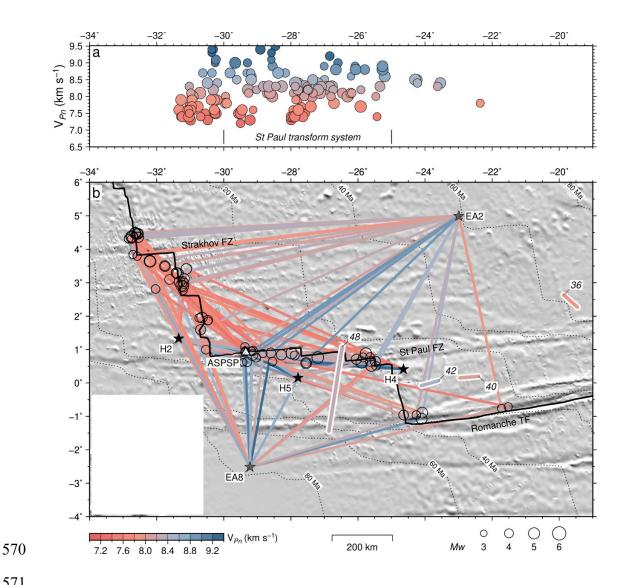


Figure 4. Example of waveforms recorded by the ASPSP seismograph and COLMEIA hydrophones, with 4–12 Hz and 6–20 Hz Butterworth filter applied, respectively, for mb 4.6 strikeslip event on 15th June 2013, located near St Paul transform fault at 29.5°W. Picked Pn arrivals, and Pn arrivals predicted by IASP91 model are marked; beach-balls are centroid moment tensors (Ekström et al., 2012); V_{Pn} and SNR noted for each station (this study).



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Figure 5. a) V_{Pn} plotted against mean longitude of ray path. Circle radius scaled by magnitude of source event; colored by V_{Pn} ; St Paul transform system marked by vertical bars. **b)** Shaded relief map showing stations, earthquakes, and ray paths. Circles are earthquakes used in Pn analysis, scaled by Mw; thin colored lines are ray paths shaded with Pn velocity; white triangle is ASPSP station; black/gray stars are COLMEIA / EA hydrophone networks, respectively (Maia et al., 2013; Smith et al., 2012); thick lines numbered 48, 42, 40, and 36 are seismic refraction profiles from cruise AT40-180 (Le Pichon et al., 1965), shaded by velocity; dotted lines are isochrons

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, modified from Müller *et al.* (2008) to remove artifacts associated with fracture zone traces. **c**)

Histogram of difference between velocity estimates from refraction experiment (Le Pichon *et al.*, 1965), and intersecting ray paths from this study; positive values indicate faster velocities estimated by refraction experiment; dark/light gray bars are velocities from profiles 48 and 42, respectively.

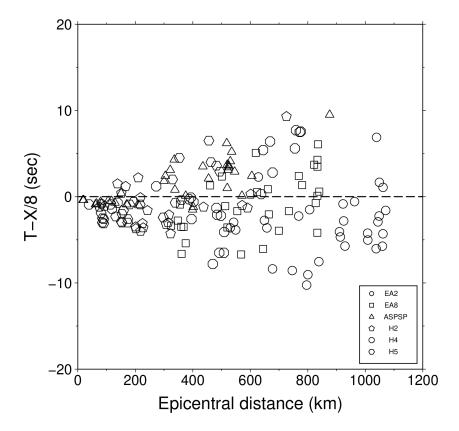


Figure 6. Reduced travel time versus epicentral distance for the 152 ray paths, plotted with a reduction velocity of 8 km s⁻¹, approximately corresponding to velocity immediately Moho from PREM and IASP91 models (dotted line; Dziewonski and Anderson, 1981; Kennett and Engdahl, 1991); key shows recording station symbols.

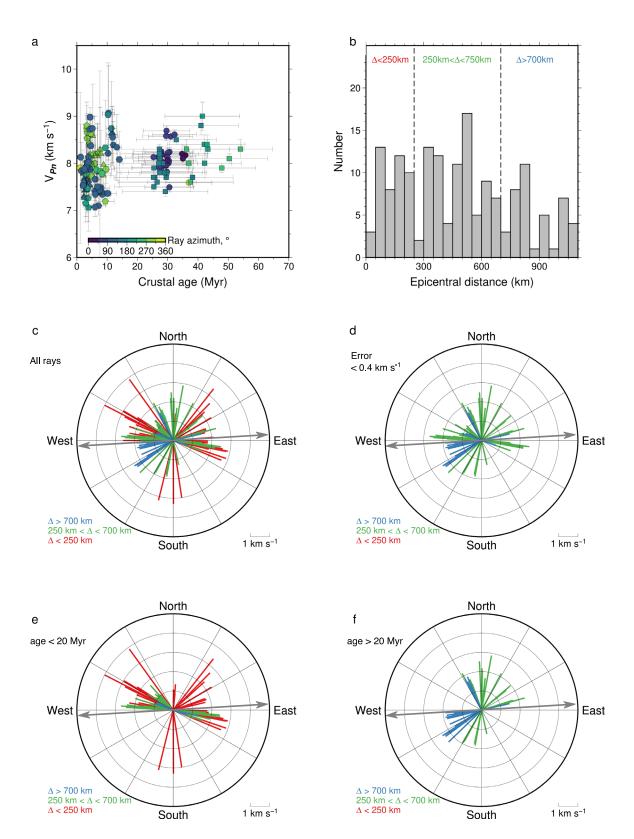


Figure 7. a) V_{Pn} plotted against oceanic crustal age at epicental location (crustal ages assigned from model of Müller *et al.* (2008), colored by ray azimuth. Triangles are estimates from ASPSP station; black/gray symbols are from COLMEIA / EA hydrophone stations, respectively (see key in Figure 6); horizontal error bars are 2σ crustal age along ray path, vertical error bars from uncertainty described in text. b) Histogram of epicentral distances; dotted lines show cut-offs used to define categories in anisotropy analysis. c) Rose diagram showing V_{Pn} vs. azimuth for all rays; length of lines scaled by V_{Pn} , colored by epicentral distance; gray arrows show plate spreading azimuth. d) V_{Pn} vs. azimuth for rays with uncertainty < 0.4 km s⁻¹, colored by epicentral distance. e) V_{Pn} vs. azimuth for rays sampling crust < 20 Myr in age, colored by epicentral distance. f) V_{Pn} vs. azimuth for rays sampling crust > 20 Myr in age, colored by epicentral distance.

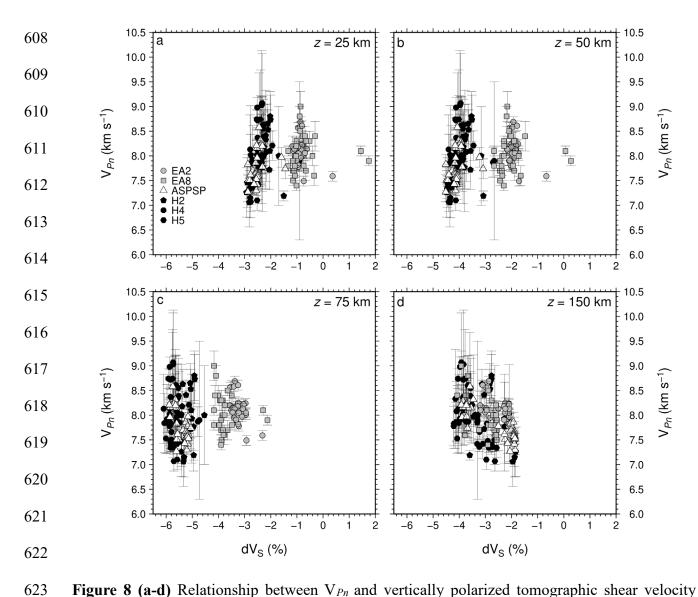


Figure 8 (a-d) Relationship between V_{Pn} and vertically polarized tomographic shear velocity anomaly at depths of 25, 50, 75 and 150 km, respectively from global model SL2013sv (Schaeffer and Lebedev, 2013). Triangles are estimates from ASPSP station; black/gray symbols are from COLMEIA / EA hydrophone stations, respectively (see key).

Station	Sensor				Cable	Water	Crust path	Crust travel time
name	type	Lat, °N	Lon, °E	Depth bsl, m	length, m	delay, s	correction, s	correction, s
ASPSP	S	0.9169	-29.3459	-16	-	-	12.5	1.9
EA2	Н	4.9907	-22.9931	800	3912	2.10	23.8	7.2
EA8	Н	-2.5159	-29.2181	800	3242	2.54	23.0	6.5
H2	Н	1.3297	-31.3445	700	2260	1.57	21.8	5.5
H4	Н	0.4123	-24.6437	700	1860	1.24	21.3	5
H5	Н	0.1552	-27.7875	700	3060	2.04	22.8	6.3

Table 1. Details of seismograph (S) and hydrophone (H) sensors used for Pn analysis in this study; including depth below sea level (bsl), cable length, and water/crust corrections.