

# Improving wave-based air-sea momentum flux parameterization in mixed seas

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## Key Points:

- Surface stress at moderate to high winds is dominated by short wind waves.
- COARE3.5 wave based formulation can underestimate surface stress by more than 10 % in mixed sea conditions under moderate to high wind.
- Using the mean wave period or including the directional alignment between wind and wave in COARE3.5 alleviates this issue.

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## Abstract

In winter, the Northwest Tropical Atlantic Ocean can be characterized by various wave age-based interactions among ocean current, surface wind and surface waves, which are critical for accurately describing surface wind stress. In this work, coupled wave-ocean-atmosphere model simulations are conducted using two different wave roughness parameterizations within COARE3.5, including one that relies solely on wind speed and another that uses wave age and wave slope as inputs. Comparisons with the directly measured momentum fluxes during the ATOMIC/EUREC<sup>4</sup>A experiments in winter 2020 show that, for sea states dominated by short wind waves under moderate to strong winds, the wave-based formulation increases the surface roughness length in average by 25% compared to the wind-speed-based approach. For sea states dominated by remotely generated swells under moderate to strong wind intensity, the wave-based formulation predicts significantly lower roughness length and surface stress ( $\approx 15\%$ ), resulting in increased near-surface wind speed above the constant flux layer ( $\approx 5\%$ ). Further investigation of the mixed sea states in the model and data indicates that the impact of swell on wind stress is over-emphasized in the COARE3.5 wave-based formulation, especially under moderate wind regimes. Various approaches are explored to alleviate this deficiency by either introducing directional alignment between wind and waves or using the mean wave period instead of the wave period corresponding to the spectral peak to compute the wave age. The findings of this study are likely to be site-dependent, and mostly concern specific regimes of wind and waves where the original parameterization was deficient.

## Plain Language Summary

Accurately understanding and describing air-sea interactions is critical for weather forecast and regional climate. In this work, we use numerical experiments with and without taking into account the ocean waves to describe air-sea interactions. Most of the momentum exchange between the ocean and the atmosphere is done through locally wind-generated waves, however remotely generated waves, such as swells, can also interfere in these air-sea interactions. Comparisons with observations made during the ATOMIC/EUREC<sup>4</sup>A field campaigns in winter 2020 show in particular that our numerical experiment overestimated the impact of the swell on the atmosphere. Various approaches are explored here to alleviate this deficiency, one of those being the introduction of the effect of the alignment between wind and waves.

## 1 Introduction

Over the ocean, most of the momentum, heat, and mass exchanges with the atmosphere are supported by short wind-waves on spatial scales of  $O(0.1-10\text{m})$ . These wind-waves enhance the surface drag and roughness at the air-sea interface, thereby increasing the wind stress. The wind stress is coupled with the planetary boundary layer (PBL) processes in the atmosphere, modifying the kinematic and thermodynamic profiles in this lowest part of the atmosphere (Janssen, 1989; Moon et al., 2004). In addition to locally generated wind-waves, the sea state is also influenced by the remotely generated swell, especially in the lower latitudes, whose propagation direction is often uncorrelated with local winds. The fast-propagating swell wave that is strongly misaligned with or outruns the local wind can be a conduit for upward momentum and energy transfer from waves to the wind, forming a wave-driven low-level jet (e.g., Harris, 1966; Sullivan et al., 2008; Hanley & Belcher, 2008) and dissipating the swell waves (M. Donelan, 1999; Kahma et al., 2016; Liu et al., 2017).

60 In numerical models, the wind stress over the oceans is parameterized using  
61 bulk flux algorithms, such as the Coupled Ocean-Atmosphere Response Experi-  
62 ment (COARE, Fairall et al., 1996, 2003; Edson et al., 2013). If no coincident wave  
63 fields are available, COARE parameterizes the wave roughness length ( $z_0$ ) using  
64 wind speeds only. In this study, this approach will be referred to as the wind-speed-  
65 dependent formulation (WSDF). Since wind and wind-waves are in near-equilibrium  
66 in many cases over the extratropical open oceans, the COARE’s WSDF tends to  
67 accurately predict the surface roughness and thereby the surface stress (Edson et al.,  
68 2013). However, under trade-wind regimes in the tropics such as our study region  
69 in boreal winter, remotely-generated swell significantly shape the sea state, whose  
70 effect on wind stress cannot be accurately characterized by local wind alone. To  
71 improve estimates of the fluxes under these conditions, “wave-based” formulations  
72 exist in many bulk flux algorithms that model  $z_0$  as a function of wave age or wave  
73 age/slope (e.g., Taylor & Yelland, 2001; Oost et al., 2002; Drennan et al., 2003;  
74 Edson et al., 2013; Sauvage et al., 2020). As there are increasing interests and op-  
75 portunities to incorporate the wave effects on surface fluxes in numerical models,  
76 such wave-based formulations (WBF) in bulk formulas will likely be adopted more  
77 in such models. Since the parameterized surface fluxes serve as lower boundary con-  
78 ditions for turbulent exchanges within the atmospheric and oceanic boundary layers,  
79 the simulation and forecast skills will be influenced by the physics and assump-  
80 tions represented in the bulk formulas. Therefore, it is imperative to understand  
81 the assumptions and deficiencies in current WBFs and offer possible revisions to the  
82 formulations for air-sea fluxes with increased accuracy. The goal of this paper is to  
83 enhance a regime-based understanding of wave-wind interactions via detailed valida-  
84 tion of the parameterized air-sea flux from high-resolution coupled model simulations  
85 against directly measured air-sea fluxes.

86 This study focuses on air-sea momentum flux during the ATOMIC/EUREC<sup>4</sup>A  
87 field campaign. The ATOMIC (Atlantic Tradewind Ocean-Atmosphere Mesoscale  
88 Interaction Campaign) is the U.S. complement to the European field campaign,  
89 EUREC<sup>4</sup>A (Elucidating the Role of Cloud-Circulation Coupling in Climate,  
90 Stevens et al., 2021), both of which took place in the Northwest Tropical Atlantic  
91 Ocean in January-February 2020 (Figure 1). The primary objective of this study is  
92 to determine how well the current WBF in an advanced bulk flux algorithm such as  
93 COARE3.5 reproduces the observed wind stress in the mixed sea conditions com-  
94 pared to the WSDF. By exploiting the fully-coupled ocean-atmosphere-wave model  
95 simulations and extensive analyses of the in situ observational datasets, we will at-  
96 tempt to explain the causes for discrepancies between simulated and measured wind  
97 stresses. Our results indicate that the current COARE3.5 WBF underestimates  $z_0$   
98 and wind stress, particularly over the mixed sea state. We will show that this is due  
99 to either a missing physics of the wave-wind interaction or using an inappropriate  
100 wave input parameter to describe the mixed sea condition.

101 The paper is organized as follows. Section 2.1 describes the technical details  
102 of the latest  $z_0$  formulation in COARE3.5. Sections 2.2 and 2.3 discuss the fully  
103 coupled ocean-atmosphere-wave modeling system used in the investigation, followed  
104 by the details on the experimental design and observational datasets in Section 2.4  
105 and Section 2.5, respectively. The wave impact on  $z_0$ , wind stress, and low-level  
106 winds are discussed in a case study investigation in Section 3. Section 4 provides an  
107 in-depth comparison of the parameterized momentum flux against the direct mea-  
108 surements, identifying the areas and regimes for further improvement. In section  
109 5, possible approaches are proposed and tested to alleviate the biases. Section 6  
110 provides a summary and discussion.

## 2 Air-sea flux parameterization and coupled model

This section provides a brief overview of the wave-mediated momentum flux implemented in the Coupled Ocean-Atmosphere Response Experiment parameterization (COARE3.5, Fairall et al., 1996, 2003; Edson et al., 2013). Hereafter, we will focus on the COARE3.5 version, although a slightly updated version, COARE3.6, has been made publicly available. However, the findings of this study would stay unchanged when using COARE3.6 (not shown).

### 2.1 Roughness length and momentum flux in COARE3.5

The along wind stress in the COARE framework is defined as:

$$\tau = \rho C_D(z, z_0, \psi_m) U_r(z) S_r(z) = \rho u_*^2, \quad (1)$$

where  $\rho_a$  is the air density,  $U_r(z)$  is the magnitude of the along-wind component of the wind vector,  $S_r(z)$  is the scalar wind speed, where the subscript  $r$  denotes relative to the ocean surface; and  $u_*$  the friction velocity.  $C_D$  is the drag coefficient defined as:

$$C_D(z, z_0, \psi_m) = \left[ \frac{\kappa}{\ln(z/z_0) - \psi_m(\zeta)} \right]^2, \quad (2)$$

where  $\kappa$  is the von Kármán constant,  $\psi_m(\zeta)$  is an empirical function of atmospheric stability,  $\zeta$  is the  $z/L$  ratio with  $L$  the Obukhov length and  $z$  the height above the surface (Fairall et al., 1996). The surface roughness length  $z_0$  is parameterized in COARE3.5 as the sum of two terms:

$$z_0 = z_0^{smooth} + z_0^{rough}, \quad (3)$$

where  $z_0^{smooth}$  and  $z_0^{rough}$  represent the smooth and rough flow components of  $z_0$ , respectively (Edson et al., 2013). The smooth flow component is parameterized as

$$z_0^{smooth} = \gamma \frac{\nu}{u_*}, \quad (4)$$

where  $\gamma$  is the roughness Reynolds number for smooth flow, set to be constant at 0.11 based on laboratory experiments, and  $\nu$  is the kinematic viscosity. For smooth flow, the wind stress is mainly supported by viscous stress where  $z_0 \approx z_0^{smooth}$ .

The rough part of the roughness length,  $z_0^{rough}$ , is meant to parameterize the wind-driven gravity waves that support most of the stress above approximately  $5 \text{ ms}^{-1}$  when the sea becomes aerodynamically rough. This component of the roughness is formulated currently in several ways in COARE3.5. The simplest and the most broadly used way is to parameterize it as a function of wind speed only. The so-called wind speed dependent formulation without explicit wave and sea states inputs estimates  $z_0^{rough}$  using the Charnock's relation (Charnock, 1955):

$$z_0^{rough} = \frac{\alpha_{CH} u_*^2}{g}, \quad (5)$$

where  $g$  is the acceleration of gravity and  $\alpha_{CH}$  is the Charnock coefficient that is dependent only on wind speed. COARE3.5 formulates  $\alpha_{CH}$  as

$$\alpha_{ch} = mU_{r10N} + b, \quad (6)$$

where  $U_{r10N}$  is the 10-m wind speed relative to the sea surface under neutral conditions (Edson et al., 2013, Appendix) and coefficients  $m = 0.0017$  and  $b = -0.005$  (?). Hereafter,  $U_{r10N}$  is defined such as:

$$U_{r10N} = \frac{u_*}{\kappa} \ln(10/z_0), \quad (7)$$

The coefficients  $m$ , and  $b$  in Eq. 6, have been determined to fit the average data used in COARE3.5 over wind speeds between 5 and 18  $ms^{-1}$ . If wind speed is below 5  $ms^{-1}$ , the surface roughness is mainly determined by  $z_{smooth}$  in Eq. 4. For wind speeds greater than 18  $ms^{-1}$ , COARE3.5 fixes the value of the Charnock coefficient to its value at 18  $ms^{-1}$ . Note, however, that although  $\alpha_{CH}$  is fixed above 18  $ms^{-1}$ ,  $z_0^{rough}$ ,  $C_D$  and  $\tau$  all continue to increase with the wind speed, just at a lower rate.

An alternative way to define  $z_0^{rough}$  in COARE3.5 is to use the so-called wave-based formulation (WBF), which requires contemporary information about the wave field and its state of development, such as significant wave height ( $H_s$ ) and phase speed of the waves at the peak of the spectrum ( $c_p$ ). Two WBFs are currently available in COARE3.5, one that uses the wave age only and another that uses both the wave age and wave steepness. In the second form, which is explored in this study in great detail,  $z_0^{rough}$  is expressed as

$$z_0^{rough} = H_s D \left( \frac{u_*}{c_p} \right)^B, \quad (8)$$

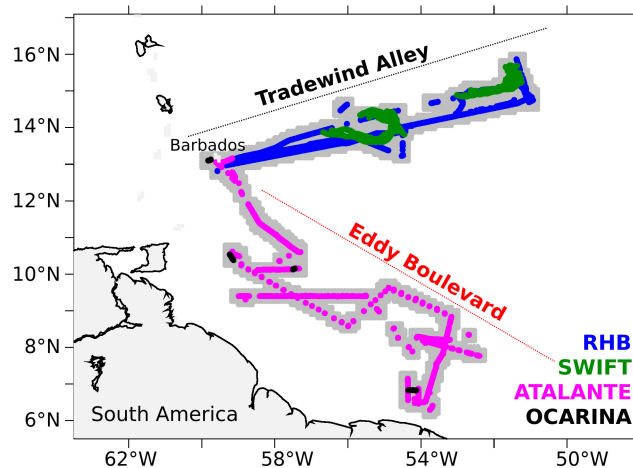
where  $u_*/c_p$  is the inverse wave age based on the friction velocity, and  $D$  and  $B$  are numerical constants given by  $D = 0.09$  and  $B = 2$  in Edson et al. (2013). Hereafter, we will use a definition of wave age based on the ratio of the phase speed of the waves at the spectral peak over the surface wind speed at 10  $m$  defined as

$$\chi = \frac{c_p}{U_{10}}. \quad (9)$$

The wave age is used to describe the state of development of the wave field. For example, a wave age close to 1.2 represents a fully developed sea when the surface waves and stress are largely in equilibrium (e.g., Phillips, 1985), in which the rate that wind does work on the surface waves is balanced by the dissipation rate of breaking waves (microbreakers and whitecaps) and nonlinear wave-wave interactions (e.g., Csanady & Gibson, 2001). Wave ages under 1 are associated with developing seas and young waves, while wave ages well above 1.2 describe decaying seas and swell. It should be noted that in the current COARE3.5,  $c_p$  is defined using the peak period of the waves,  $T_p$ , in deep water such that:

$$c_p = g \frac{T_p}{2\pi}. \quad (10)$$

In Section 3, we will examine the sensitivity of the estimated momentum flux based on the current COARE3.5 algorithm. Guided by comparison to the observations in Section 4, we will then explore the impacts of revised COARE3.5 WBF in Section 5.



**Figure 1.** Tracks of the different platforms measuring surface stress. The gray area denotes where the model outputs are sampled along the tracks of observations. RHB provided data from January 9 to February 13, 2020. SWIFT drifters were deployed from 14 January to 22 January 2020 and from 30 January to 11 February 2020. R/V ATALANTE provided data from January 19 to February 19, 2020 and Ocarina was deployed periodically from January 25 to February 17, 2020.

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## 2.2 SCOAR regional coupled model system

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We use the Scripps Coupled Ocean-Atmosphere Regional (SCOAR) model (Seo et al., 2007, 2021), which couples the Weather Research and Forecast (WRF, Skamarock et al., 2008) Model to the Regional Ocean Modeling System (ROMS, Shchepetkin & McWilliams, 2005) via the COARE3.5 bulk flux algorithm (Fairall et al., 1996, 2003; Edson et al., 2013). In the absence of wave coupling, ROMS is driven by the surface heat flux ( $Q_{NET}$ ), momentum flux ( $\tau$ ), and freshwater flux ( $Q_{FW}$ ) computed from the wind speed-only formulation in COARE3.5 implemented in WRF. In turn, ROMS inputs SST and surface current vectors ( $U_s$ ) to the COARE3.5 to compute the surface fluxes (Figure 2).

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## 2.3 Wave coupling in SCOAR

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This study implemented the coupling of the third-generation spectral wave model WaveWatch-III (WW3 Tolman et al., 2002; The WAVEWATCH III Development Group, 2016) into the SCOAR. Currently, two different ways are implemented to allow coupling waves to the atmosphere. The first option described in Figure 2 is based on the total friction velocity output from WW3 and used to estimate the wind stress and the resulting surface roughness length for computing turbulent heat fluxes. This option won't be used in this study. The second and third options described in Figure 2 are the focus of this manuscript and respectively take advantage of the COARE's WBF from (Edson et al., 2013), and the finding of this study. In this configuration, the centerpiece of the model coupling is the COARE3.5 implemented in the surface layer scheme in WRF to compute the air-sea fluxes. In this study, we use the Mellor-Yamada-Nakanishi-Niino (MYNN) surface layer scheme (Nakanishi & Niino, 2009; Jimnez et al., 2012), which over the ocean grid points computes the surface fluxes using the COARE3.5 WBF. WW3 is forced by the surface wind ( $U_{10}$ ) from WRF and ocean current ( $U_s$ ) from ROMS. WW3 then returns the significant wave height ( $H_s$ ) and the phase speed of the dominant waves ( $c_p$ )

205 determined based on  $T_p$  (Eq. 10) to the MYNN surface layer scheme. In lieu of  $c_p$ ,  
 206 WW3 can alternatively send the mean phase speed ( $c_m$ ) and peak wave direction  
 207 (Section 5). Spatially varying Charnock coefficients ( $\alpha_{CH}$ ) are then updated to pa-  
 208 rameterize the surface roughness length ( $z_0$ ) as a function of dominant wave age  
 209 ( $\chi$ ) and wave steepness (Eq. 8). For this to work in WRF, the MYNN surface layer  
 210 scheme has been modified to allow ingestion of wave age and significant wave height  
 211 ( $H_s$ ) from WW3. The MYNN PBL scheme (Nakanishi & Niino, 2004, 2006) is cou-  
 212 pled to this modified surface layer scheme, allowing for the adjusted  $z_0$ , wind stress  
 213 ( $\tau$ ), and latent ( $Q_{LH}$ ) and sensible ( $Q_{SH}$ ) heat fluxes to influence the kinematic  
 214 and thermodynamics processes in the PBL. The surface layer scheme has also been  
 215 modified to take the ocean surface currents ( $U_s$ ) from ROMS to compute the rela-  
 216 tive wind and thus represent wind-current interaction. This so-called relative wind  
 217 effect is represented in all simulations analyzed here. Wave to ocean coupling is also  
 218 made available and ROMS can be forced by wave fields such as  $H_s$  and wave energy  
 219 ( $FOC$ ) fields. Wave-supported stress ( $\tau^w$ ) and wave dissipation ( $\tau^{ds}$ ) terms can also  
 220 be send to ROMS to compute the ocean-side stress ( $\tau^{oc}$ ). For the purpose of this  
 221 study, wave to ocean coupling is not included and thus on Figure 2 it is assumed  
 222 that  $\tau^{oc} = \tau^a$ , where  $\tau^a$  is the air-side stress.

## 2.4 Experiments

224 In WRF, the deep cumulus convection is represented through the Multi-scale  
 225 Kain-Fritsch scheme (Zheng et al., 2016), the cloud micro-physics by the WRF  
 226 single-moment 6-class scheme (Hong & Lim, 2006). The Goddard radiation scheme  
 227 (Chou & Suarez, 1999) is used for shortwave and longwave radiation. The land  
 228 surface process is treated with the Noah land surface model (F. Chen & Dudhia,  
 229 2001). In ROMS, the KPP (K profile parameterization) scheme (Large et al., 1994)  
 230 determines vertical eddy viscosity and diffusivity. The vertical grid in ROMS is  
 231 stretched to enhance the resolutions near the surface and the bottom, using the so-  
 232 called stretching parameters of  $\theta_s = 7.0$ ,  $\theta_b = 2.0$ , and  $h_{cline} = 300$  m. In WW3, the  
 233 set of parameterizations from Ardhuin et al. (2010) is used, including swell dissipa-  
 234 tion scheme (Ardhuin et al., 2009). Nonlinear wave-wave interactions are computed  
 235 using the discrete interaction approximation (Hasselmann et al., 1985). Reflection  
 236 by shorelines are enabled through Ardhuin and Roland (2012) scheme. The depth-  
 237 induced breaking is based on Battjes and Janssen (1978), and the bottom friction  
 238 formulation follows Ardhuin et al. (2003).

239 The model domain covers the Northwest Tropical Atlantic Ocean (Figure 3).  
 240 The horizontal resolutions in WRF, ROMS, and WW3 are identical 10 km, with  
 241 matching grids and land-sea masks. This horizontal resolution allows us to have  
 242 reasonable description of the mixed sea state influenced by the remotely-generated  
 243 swell and trade winds in the open oceans, which is the focus of this work. However,  
 244 much finer-scale wind-wave and wave-current interactions, as studied in (Ardhuin et  
 245 al., 2017; Bas et al., 2020; Iyer et al., 2022), are not likely captured at this resolu-  
 246 tion, especially in the regions of strong currents and eddy variability. ROMS (WRF)  
 247 is run with a stretched vertical grid with a total of 30 (33) vertical levels, with ap-  
 248 proximately 10 layers in the upper 150 m (below 1300 m). The model coupling is  
 249 activated every 3 hours to account for the diurnal cycle.

250 A set of coupled model simulations presented in Section 4 is run for 6 months  
 251 (November 1, 2019 to May 1, 2020), covering the ATOMIC/EUREC<sup>4</sup>A period, with  
 252 a specific aim to compare with the measurements. In these simulations, the WRF  
 253 model is initialized and driven by 3-hourly ERA5 global reanalysis at  $0.25^\circ$  resolu-  
 254 tion (Hersbach et al., 2018a, 2018b), ROMS by the daily MERCATOR International  
 255 global reanalysis at  $1/12^\circ$  resolution (Lellouche et al., 2018), and WW3 by seven  
 256 spectral points obtained from the global  $1/2^\circ$  resolution WW3 simulations (Rasclé

257 & Ardhuin, 2013). The initial conditions for ROMS and WW3 were obtained from  
 258 the respective ROMS-only and WW3-only spin-up simulations forced by ERA5 at-  
 259 mospheric forcing (starting from January 1, 2019). In ROMS, the tidal forcing is  
 260 obtained using the Oregon State University Tidal Prediction Software (Egbert &  
 261 Erofeeva, 2002) and applied as a 2-D open boundary condition by prescribing the  
 262 tidal period, elevation amplitude, current phase angle, current inclination angle, the  
 263 minimum and maximum tidal current, and ellipse semi-minor axes for 13 major tidal  
 264 constituents. Daily climatology estimates of the Amazon and River and Orinoco  
 265 River discharges are obtained from the Observatory Service SO-HyBAM database  
 266 (<https://hybam.obs-mip.fr/>), which are prescribed as point sources close to the river  
 267 mouths in our grid.

268 The second set of simulations presented in Section 3 is identical to that of the  
 269 6-month-long simulations, except that WRF, ROMS, and WW3 are initialized from  
 270 respectively 3-hourly ERA5 global reanalysis for the atmosphere and ROMS-only  
 271 and WW3-only spin-up simulations for the ocean and waves as described above  
 272 and run on a particular day (January 8, 2020) as a case study investigation. The  
 273 motivation for the short simulations with the identical initial condition is to isolate  
 274 the immediate impacts on  $z_0$  and  $\tau$  before the coupled feedback begins to alter the  
 275 state variables. One could use the identical input state variables to estimate the  
 276 air-sea fluxes offline using different COARE formulations. This yields similar results  
 277 (not shown), indicating that the difference we show in Section 3 is not due to the  
 278 difference in state variables, but due to the formulation difference. One notable ad-  
 279 vantage to use the fully coupled model simulation is that it allows for evaluating the  
 280 wind response beyond the surface layer (e.g., Figure 6c), and potentially large-scale  
 281 feedback effects via the coupling.

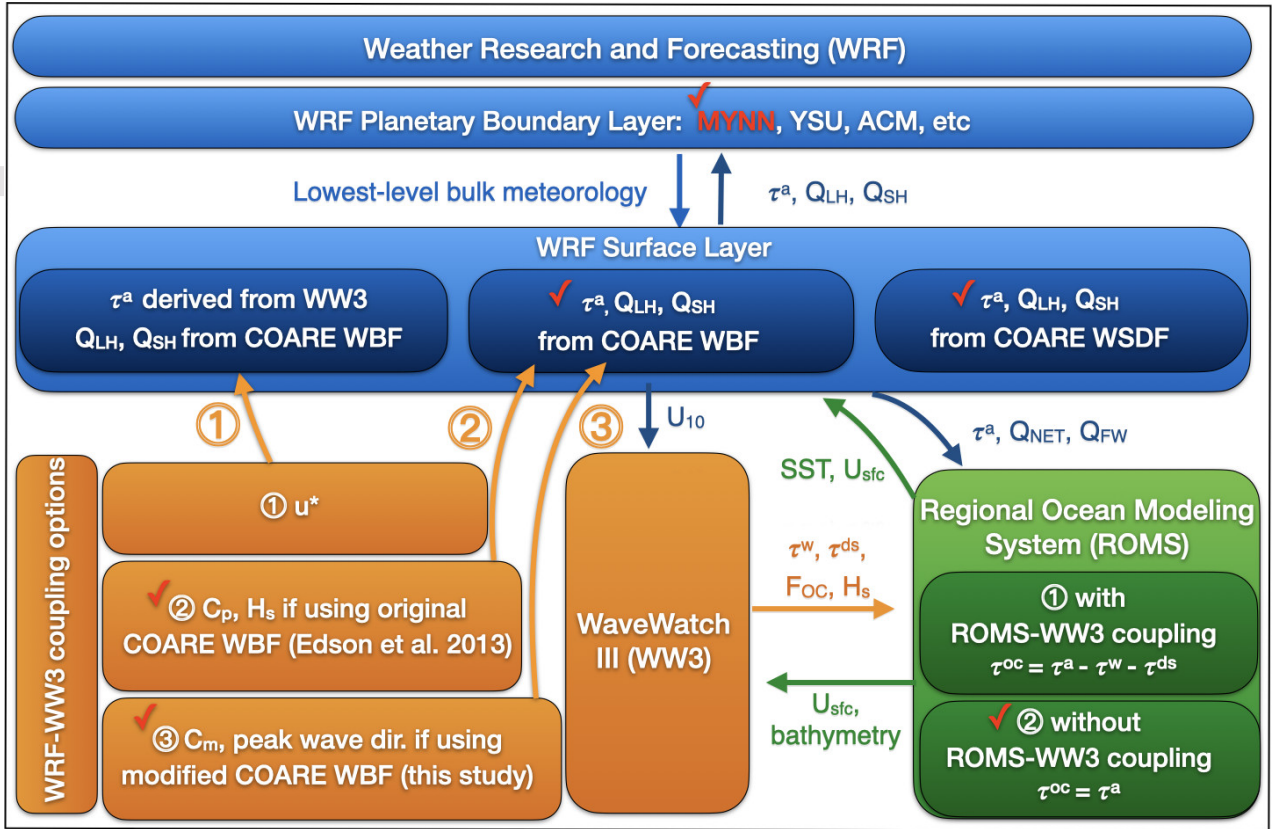
**Table 1.** Summary of the different SCOAR experiments.

Experiments	$z_0$ parameterization	Relative wind	Wave period	misaligned wave
WSDF	wind speed [Eq. 5]	yes	/	/
WBF	wave age + wave steepness [Eq. 8]	yes	$T_p$	no
WBF $_{\theta}$	wave age + wave steepness [Eq. 11]	yes	$T_p$	yes
WBF $_{T_m}$	wave age + wave steepness [Eq. 12]	yes	$T_m$	no

282 Table 1 summarizes 4 experiments conducted in this study, where the only  
 283 difference is in the way  $z_0$  is parameterized in COARE3.5. In the first run (dubbed  
 284 WSDF), the wind speed only formulation is used (hence, only WRF-ROMS cou-  
 285 pling), while in the second run (WBF), the default wave-based formulation is used  
 286 (WRF-ROMS-WW3). These two runs are examined in detail in Sections 3-4. Two  
 287 additional runs, discussed in Section 5, are conducted with a modified wave-based  
 288 formulation. WBF $_{\theta}$  takes into account the directional misalignment between wind  
 289 and wave, while WBF $_{T_m}$  modifies the definition of wave age based on mean wave  
 290 period rather than the peak wave period.

291 All simulations used in this study produce output every 3h. Since this output  
 292 interval is much coarser than the typical sampling intervals used in the observations  
 293 (Section 2e), there is inevitable inconsistency in sampling frequency and the number  
 294 of samples between the model and data. We attempt to increase the model sample  
 295 size and capture more spatio-temporal variability by sampling a slightly broader  
 296 region of the model domain encompassing the particular observational tracks (gray  
 297 areas in Figure 1a). By doing this we assume that the spatial variability sampled in





**Figure 2.** SCOAR WRF-ROMS-WW3 coupling flowchart. See the text for the variable names that are exchanged across the model components. Red ticks denote of the specific schemes and coupling methodology used in this study.

the model would resemble the temporal variability observed, considering that the spatial extent of our model sampling is still relatively close to the different platform tracks.

## 2.5 ATOMIC/EUREC<sup>4</sup>A observations

This study will exploit direct and indirect measurements of momentum fluxes and relevant wave fields (i.e., significant wave height and wave period) from various platforms deployed during the ATOMIC/EUREC<sup>4</sup>A experiment, summarized in Table 2. Figure 1 shows the tracks of the different observational platforms, including the NOAA R/V Ronald H. Brown (RHB, Quinn et al., 2021; Thompson et al., 2021), R/V ATALANTE (Bourras, Geyskens, et al., 2020), SWIFT drifters (Surface Wave Instrument Float with Tracking, Thomson, 2012; Thomson et al., 2019, 2021), and OCARINA (Ocean Coupled to Atmosphere, Research at the Interface with a Novel Autonomous platform, (Bourras, Branger, et al., 2020)) surface naval drone. The RHB provides direct momentum flux measurements every 10 minutes, using the eddy covariance method, in the so-called “Tradewind Alley” region from January 9 to February 13, 2020. The SWIFT drifters were deployed from the RHB, from which the hourly stress can be estimated using the equilibrium frequency range in the wave spectrum. More specifically, the directional wave spectra and bulk wave parameters were estimated from inertial motion observations. Then, the friction velocity at equilibrium  $u_*$  is calculated from the wave spectra, assuming a constant equilibrium

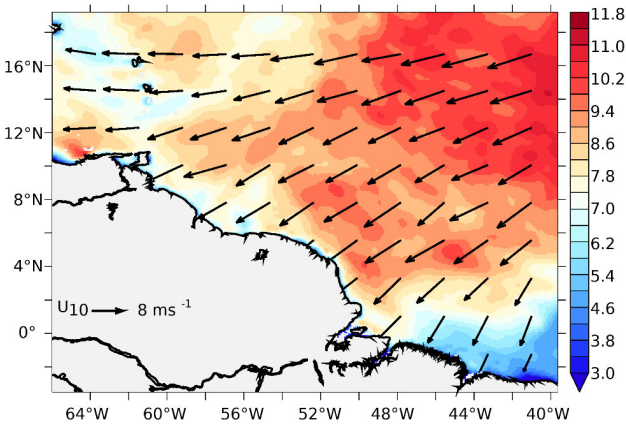
318 frequency range over which the source and sink of wave energy is balanced (Iyer  
 319 et al., 2022). They were deployed from 14 January to 22 January 2020 and from  
 320 30 January to 11 February 2020. The R/V ATALANTE measured the wind stress  
 321 mostly in the “Eddy Boulevard” region based on the inertial dissipation method  
 322 during the period of January 19 to February 19, 2020. OCARINA was deployed  
 323 periodically from the R/V ATALANTE from January 25 to February 17, 2020, pro-  
 324 viding direct wind stress measurements every minute through the eddy covariance  
 325 method.

**Table 2.** Summary of the different ATOMIC/EUREC<sup>4</sup>A observations used in this study.

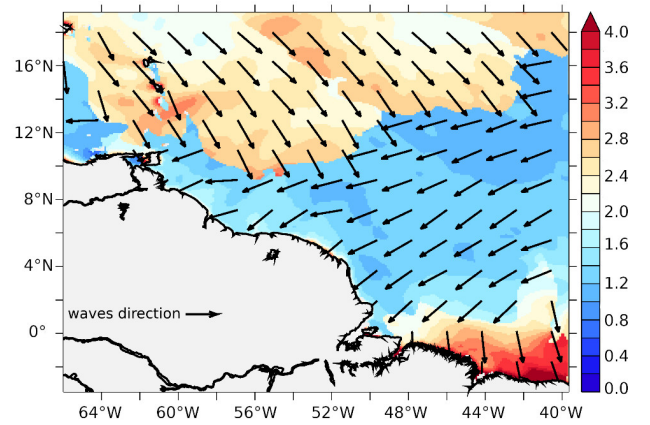
Platforms	R/V Ronald H. Brown	SWIFT	R/V ATALANTE	OCARINA
Observations	wind stress wave periods significant wave height	wind stress wave periods significant wave height	wind stress	wind stress
Methods used in estimating wind stress	eddy covariance	estimated through wave equilibrium subrange	inertial dissipation	eddy covariance
Periods	January 9 to February 13, 2020	14 January to 22 January 2020	January 19 to February 19, 2020	January 25 to February 17, 2020 (periodically)

326 RHB provided data from January 9 to February 13, 2020. SWIFT drifters were  
 327 deployed from 14 January to 22 January 2020 and from 30 January to 11 February  
 328 2020. R/V ATALANTE provided data from January 19 to February 19, 2020 and  
 Ocarina was deployed periodically from January 25 to February 17, 2020

(a) 10-m wind speed and direction



(b) Wave age and wave peak direction



**Figure 3.** Snapshots of (a) 10-m wind speeds (shading,  $ms^{-1}$ ) and direction (arrows) and (b) peak wave age (shading) and wave peak direction (arrows) on January 8, 2020 at 0600 UTC.

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### 3 Impacts of wave and sea state: a case study

To demonstrate the immediate effect of including waves on  $z_0$  and  $\tau$  in the COARE3.5 using a coupled model, we will first compare the simulation results close to the initial condition. By doing so, the input state variables into the bulk formula remain largely identical, and any differences in simulated  $z_0$  and  $\tau$  can be attributed to the difference in the formulations. From this set of experiments, we will compare the results 3 hours after the initial condition.

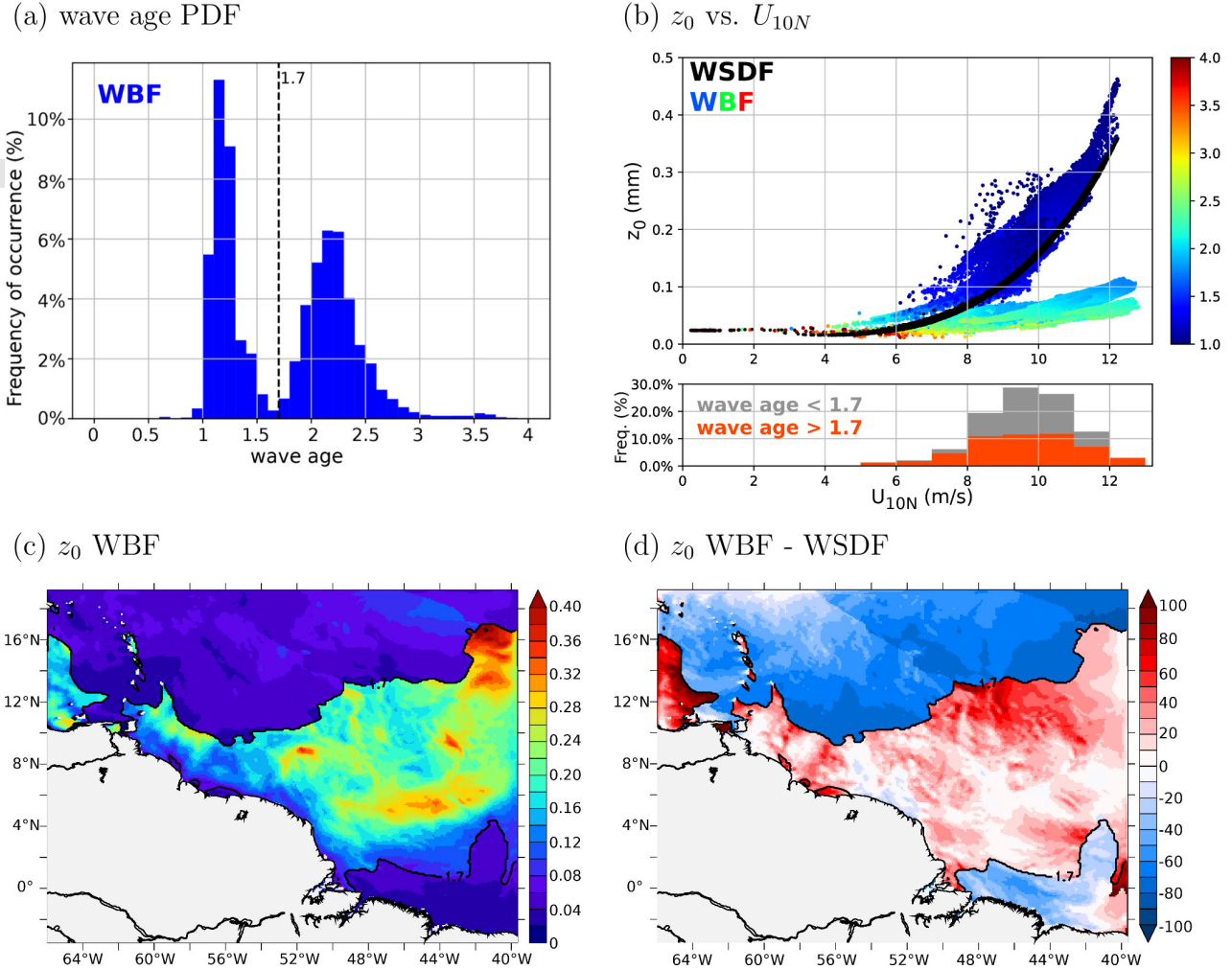
The sea state and wind fields on January 8, 2020 at 0600 UTC, shown in Figure 3a, illustrate the archetypal synoptic condition observed in this region during the boreal winter. Much of the domain was under the influence of northeasterly trade winds with wind speeds of 7-13  $ms^{-1}$ , while the northern and southeastern parts of the domain experienced much weaker ( $<7 ms^{-1}$ ) easterly and northerly winds, respectively. Figure 3b shows the corresponding wave age and peak wave direction. In the Tradewind Alley region, surface waves were predominantly downwind with relatively small wave age, indicating the developing seas with young waves. Away from the trade winds, especially in the northern part of the domain, the wave vectors are generally misaligned with the local wind vectors, and the wave age is high, indicative of the swell-dominated sea state.

To illustrate sea state distribution differently, Figure 4a shows the probability density function (PDF) of wave age for the same period. Two distinct peaks of wave age stand out clearly. The first peak resides on wave age between 0.8 and 1.7, corresponding to developing (young) waves to fully developed (mature) seas. The secondary peak is found over a wide range of wave age greater than 1.7, reaching up to 4-5, the latter representing swell. Indeed, the fact that there is a gap at 1.7 strongly suggests that the older waves are swell, as opposed to the continuum of longer/older wind waves. Thus, in this case, we choose to use 1.7 as a threshold for fully developed seas and not the usual value of 1.2 which is what you might expect for wind waves dominated region. As a matter of fact, this swell-dominated sea state is frequently observed in the ATOMIC region in the boreal winter (e.g., Semedo et al., 2011; Jiang & Chen, 2013). Indeed, if considering the entire month of January 2020 in our simulations, we find that wave ages greater than 2 occur more than 60% of the time in this domain.

Figure 4b compares the  $z_0$  against wind speed from the WSDF (black) and WBF (color) runs for this period.  $z_0$  from WBF is color-coded to denote the corresponding wave age. The bottom panel shows stacked PDFs of 10-m wind speeds from WBF, with the red (gray) parts representing the proportion of wind associated with wave age over (under) 1.7. The WSDF in COARE3.5 assumes young seas under moderate to high winds, and hence the parameterized  $z_0$  (black) obeys the well-known quadratic dependence on wind speed. The surface roughness  $z_0$  from WSDF shows less scatter because it is based solely on wind speed.

In contrast, WBF captures the two wave age-dependent regimes of  $z_0$  that appear distinct from WSDF. The first is the cluster of  $z_0$ , which increases more rapidly with wind speed than WSDF  $z_0$  and occurs over 4-12  $ms^{-1}$ . The wave age of this cluster (shading) is typically less than 1.7, corresponding to the first wave age peak in Figure 4a of small-scale young waves. Thus, the developing and equilibrium waves under these wind speeds and wave age conditions increase  $z_0$  in WBF compared to WSDF.

The second cluster indicates significantly decreased  $z_0$  in WBF with wind speed up to 12  $ms^{-1}$ . This cluster can be further split into two different wind speed groups, under and above 8  $ms^{-1}$ , color-coded by the PDF of winds (Figure 4b). Below 8  $ms^{-1}$  (red, weak winds), the wave age mainly constitutes the tail of the



**Figure 4.** (a) PDF of wave age from the entire model domain on January 8, 2020 at 0600 UTC. The dotted vertical line denotes the wave age of 1.7, below (above) which the sea state is characterized as developing, equilibrium and slightly old waves (mature waves and swell). The upper panel of (b) is a scatter plot of  $z_0$  (mm) vs.  $U_{10N}$  ( $ms^{-1}$ ).  $z_0$  from WSDF is shown in black, while  $z_0$  from WBF is color-coded to denote the corresponding wave age. The stacked PDFs of  $U_{10N}$  in the lower panel of (b) are constructed when wave age is above 1.7 (red) and below 1.7 (gray). (c) A map of  $z_0$  from WBF, superposed with a contour of wave age = 1.7. (d) A map of percentage difference of  $z_0$  between WBF and WSDF

381 PDF distribution shown in Figure 4a with an average wave age of 2.7. It is where re-  
 382 motely generated swell appears to dominate the sea state. However, the wind speeds  
 383 under  $8 ms^{-1}$  account for less than 10% of the total wind speed data, and thereby it  
 384 has a relatively small impact on the space/time-averaged  $z_0$ . Indeed, when averaged  
 385 for wind speed below  $8 ms^{-1}$ , the percentage difference in  $z_0$  between WSDF and  
 386 WBF, defined as  $(WBF-WSDF/WSDF)*100$ , is only -1.7%.

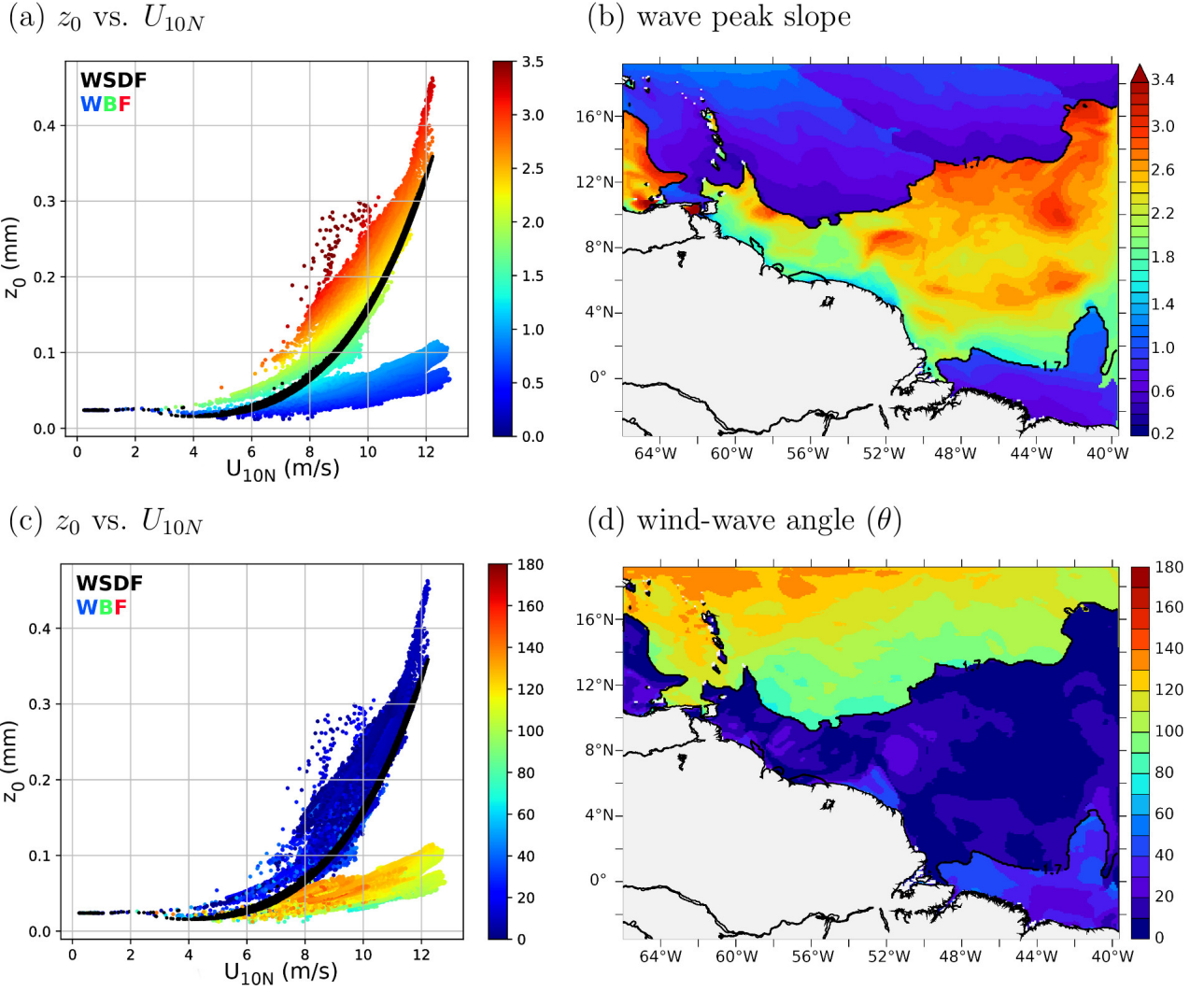
387 During this day, most of the wind speed is above  $8 ms^{-1}$ . In addition to the  
 388 proportion of low wave age expected under this moderately high wind speed, we also  
 389 find an increased occurrence of large wave age, accounting for 44% of the data (Fig-  
 390 ure 4b). The co-existence of high wind and swell indicates a mixed sea condition. In

391 this case, when averaged over wind speed above  $8 \text{ ms}^{-1}$ , the swell impact appears  
 392 much more significant, with  $z_0$  in WBF being 15.7% lower than that in WSDF. The  
 393 working hypothesis is that the use of the phase speed at the spectral peak causes the  
 394 WBF to assume that the swell is supporting most of the stress even under moderate  
 395 winds. This strong impact of swell on  $z_0$  at such moderately strong winds is ques-  
 396 tionable, in the sense that the majority of air-sea momentum exchanges should still  
 397 be supported by short-scale coupled wind waves despite the co-existence with the  
 398 long-wave swell.

399 The spatial distribution of  $z_0$  from WBF is shown in Figure 4c. The  $z_0$  dif-  
 400 ference between WBF and WSDF is shown in Figure 4d. As in Figure 4a,b, two  
 401 distinct regimes of  $z_0$  are readily apparent on the map, delineated sharply by the  
 402 contour of wave age 1.7 (black). The horizontal discontinuities in the wave and  $z_0$   
 403 fields (Figure 4c,d) appear only with the use of the peak period, while the use of  
 404 average wave period produces much smoother fields (not shown). The location of the  
 405 front is only because this is a snapshot of the sea state on 8 January at 0600 UTC.  
 406 Snapshots 3h before/after would show the swell front displaced to another location  
 407 as the swell is moving/dissipating. In the first regime of increased  $z_0$  in WBF under  
 408 moderate to strong trade winds, the WBF predicts an increased  $z_0$  by on average  
 409 25% compared to WSDF. This increased  $z_0$  is expected as the WBF  $z_0$  formulation  
 410 (Eq. 8) takes into account the effect of wave slope on the aerodynamic roughness  
 411 of the sea surface. That is, Figure 5a,b show that wave slope under young waves is  
 412 higher, where the choppy sea surface increases  $z_0$ . Figure 5c,d shows the angle ( $\theta$ )  
 413 between the wind direction and peak wave direction. If  $\theta = 0^\circ$ , wind and waves are  
 414 perfectly aligned, whereas  $\theta = 180^\circ$  means wind and waves are opposed. Collocated  
 415 with the regime of increased  $z_0$ , the peak wave direction is largely downwind, since  
 416  $\theta$  is generally less than  $50^\circ$ . This corroborates that these waves are young waves  
 417 driven by local winds. In the present study only the peak wave direction is used to  
 418 defined alignment/misalignment with the local wind. However, at times, the wave  
 419 field can yield significant directional spreading, this aspect is discussed later on in  
 420 Section 5.2.

421 Figure 4d also shows the second regime of decreased  $z_0$  with the inclusion of  
 422 waves, especially in the northern part of the domain. In this region, the remotely  
 423 generated swell propagates into the domain through the northern boundary and  
 424 forms a sea state with the aerodynamically smooth sea surfaces (Figure 5a,b) and  
 425 with waves whose direction is strongly misaligned ( $\theta = 60-160^\circ$ ) with the local wind  
 426 (Figure 5c,d). In particular, the reduced  $z_0$  over swell persists under wind speed of  
 427 up to  $12 \text{ ms}^{-1}$  (Figure 3a), despite the expectation that under such a high wind, the  
 428 wind-waves would still strongly increase the aerodynamic roughness and stress.

429 Figure 6a,b compare the parameterized wind stress in WBF and WSDF. One  
 430 can see from these plots a consistent difference in wind stress due to the inclusion of  
 431 waves. Wind stress decreases sharply in wind speeds of  $8-12 \text{ ms}^{-1}$  over the northerly  
 432 swell, where wave age  $>1.7$ . At the highest wind speed during the event, the per-  
 433 centage difference in wind stress magnitude exceeds 10%. Conversely, wind stress  
 434 is increased in WBF by  $\approx 4\%$  over fully developed seas (wave age  $<1.7$ ) and high  
 435 winds, consistent with the increase in  $z_0$  there (Figure 4c). By comparing to the  
 436 direct momentum flux observations, we will determine in Section 4 if such reduced  
 437  $z_0$  and  $\tau$  over swell conditions at moderate to high wind speeds are consistent with  
 438 the observations. As COARE3.5 does not consider the misaligned waves with winds,  
 439 these conditions may constitute a source of uncertainty in the parameterized  $z_0$  and  
 440  $\tau$  via COARE3.5 WBF. As for the large wave age in the southeastern corner of the  
 441 domain, it is concurrent with weaker winds (Figure 3a), and hence the assumptions  
 442 about the swell under weaker wind seem valid in this region. This leads to a small  
 443 difference in  $z_0$  between WBF and WSDF.

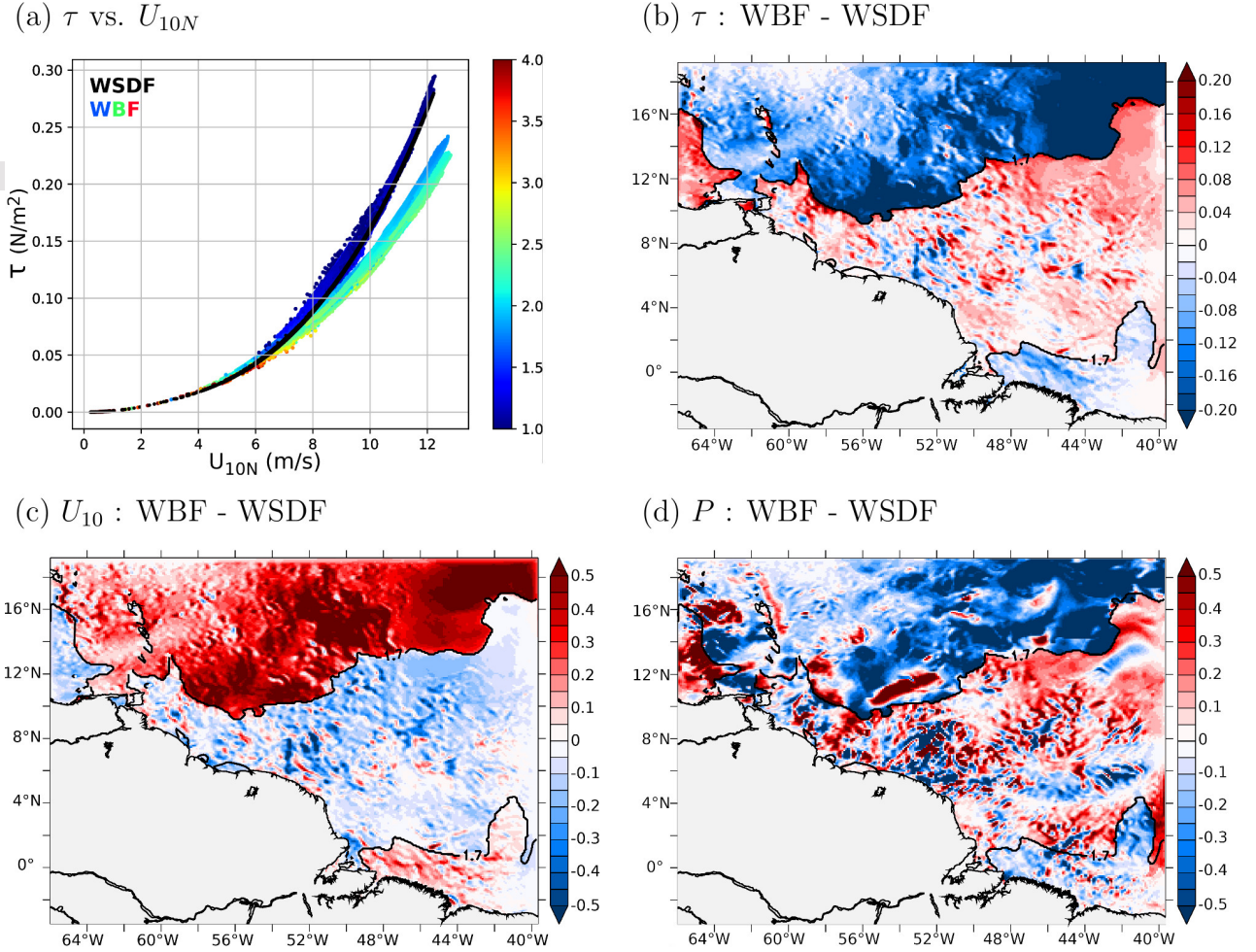


**Figure 5.** (a) Scatter plot of  $z_0$  (mm) vs.  $U_{10N}$  ( $ms^{-1}$ ) from WSDF in black and WBF color-coded to denote the corresponding wave peak slope ( $10^{-2}$ ) defined as  $H_s/L_p$  where  $L_p$  is the peak wavelength. (b) A map of wave slope peak ( $10^{-2}$ ), superposed with a contour of wave age = 1.7 on January 8, 2020 at 0600 UTC. (c,d) As in (a-b) except that colored scatters and shading denote the angle between the wind and wave directions ( $^{\circ}$ ).

444 The altered stress directly influences the low-level winds via the surface drag.  
 445 Here, we estimate the response in low-level winds at the lowest WRF model layer,  
 446 at about 27 m above the sea surface. Figure 6c shows that the low-level wind is  
 447 increased over the aerodynamically smooth sea surface due to swell by  $>0.5 ms^{-1}$ ,  
 448 accounting for 5-20% of the wind speed in WBF. In contrast, where young waves  
 449 dominate in WBF, the wind stress is increased by 5% and the wind speed is de-  
 450 creased.

451 One relevant physical process that represents the air-sea momentum transfer  
 452 affecting the winds and surface currents, is the wind work ( $P$ ),

$$P = \frac{1}{\rho_o} (\overline{u_s \tau_x} + \overline{v_s \tau_y}), \quad (11)$$

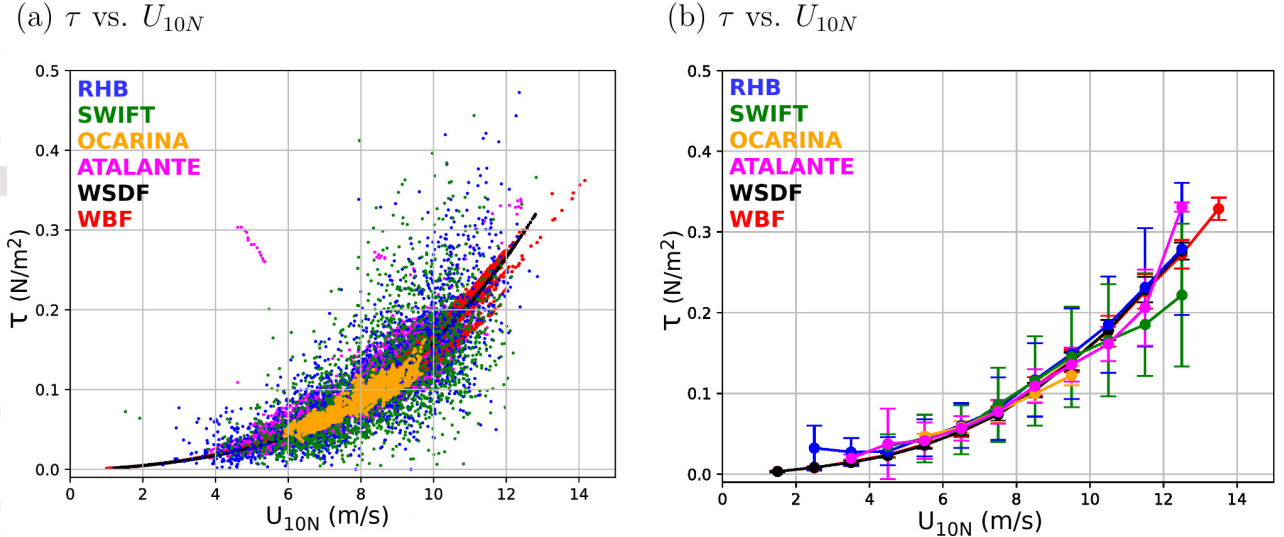


**Figure 6.** (a) Scatter plot of  $\tau$  ( $\text{Nm}^{-2}$ ) vs.  $U_{10N}$  ( $\text{ms}^{-1}$ ) from WSDF in black and WBF color-coded to denote the corresponding wave age. (b,c,d) Difference maps between WBF and WSDF of (b)  $\tau$  ( $10^{-1} \text{Nm}^{-2}$ ), (c)  $U_{10}$  ( $\text{ms}^{-1}$ ), and (d) wind work ( $P$ ,  $10^{-5} \text{m}^3 \text{s}^{-3}$ ) on January 8, 2020 at 0600 UTC, superposed with a contour of wave age = 1.7.

453 where  $(u_s, v_s)$  are the surface current vectors,  $(\tau_x, \tau_y)$  are the wind stress vectors, and the overbar denotes the time-average. When  $P$  is positive, the mechanical  
 454 work is done by the wind stress on the ocean surface currents, increasing the ocean  
 455 kinetic energy (e.g., Wunsch, 1998). When negative, it represents the diversion of  
 456 the ocean energy by the current to the wind, accelerating the low-level winds at  
 457 the expense of weakened surface currents (e.g., Renault et al., 2016, 2017; Seo et  
 458 al., 2019, 2021). Figure 6d shows the difference in  $P$  between WBF and WSDF for  
 459 this snapshot. The region of reduced  $\tau$  and increased low-level wind in the swell-  
 460 dominated region is congruent with the region of the robust decrease in  $P$ , while the  
 461 opposite is true in the Tradewind Alley region. The difference in  $P$  mainly reflects  
 462 the changes in wind stress due to waves (Figure 6b).  
 463

#### 464 4 Modeled and observed momentum fluxes during ATOMIC

465 Determining whether or not the parameterized  $z_0$  and  $\tau$  with WBF represents  
 466 an improvement over WSDF requires a detailed comparison to direct covariance



**Figure 7.** (a) Scatter plot comparing the two parameterized  $\tau$  ( $Nm^{-2}$ ) using COARE3.5 WSDF (black) and WBF (red) against the various types of measurements of  $\tau$  (see Section 2e for a description of the various methodologies). (b) As in (a) except that measurements are bin-averaged with a wind speed bin-size of  $U_{10N} = 1 \text{ ms}^{-1}$ . The error bars represent  $\pm 1$  standard deviation. Only bins with more than 5 points are plotted.

467 stress measurements. In this section, we will compare the model simulation with  
 468 the observations during the EUREC<sup>4</sup>A/ATOMIC experiments to evaluate the ac-  
 469 curacy of the wave-based parameterized  $\tau$  and identify the regimes where further  
 470 improvements might be needed.

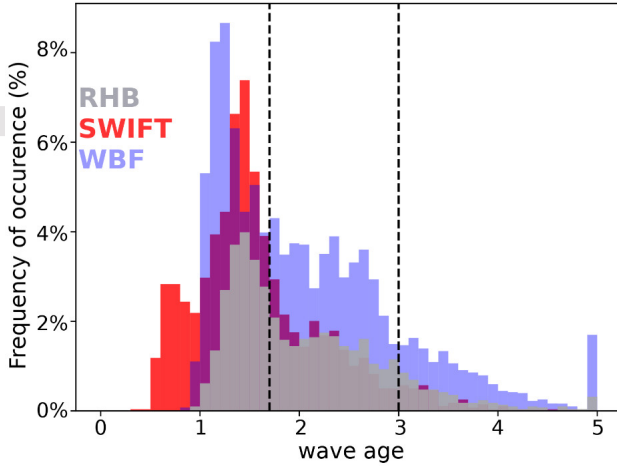
471 Figure 7a compares the two modeled stresses to the observations. All observa-  
 472 tions and the two model simulations display the quadratic relationship of wind stress  
 473 with wind speed. RHB and SWIFT, sampling the stress mainly in the Tradewind  
 474 Alley region, produce greater scatter compared to ATALANTE and OCARINA,  
 475 which were deployed further south in the Eddy Boulevard region (1a). The signifi-  
 476 cant departure from this curve in the Tradewind Alley region may reflect the greater  
 477 uncertainties in determining  $\tau$  from these measurements. Between the model simu-  
 478 lations, WBF produces a larger spread than WSDF, yet their averages at given wind  
 479 speed are similar (Figure 7b). Overall, parameterized stresses by WSDF and WBF  
 480 both agree well with the observations to within the observational errors during the  
 481 campaign.

482 Figure 8a compares the histograms of the wave age from the WBF run to those  
 483 from the SWIFT drifters and the RHB. It should be noted that in both the model  
 484 and measurements, the wave age is estimated using the peak period ( $T_p$ ). The ob-  
 485 servations and model simulation show the bi-modal distribution of wave age as was  
 486 seen from the snapshot case in Section 3 (Figure 4a), with the first peak near wave  
 487 age 1.7 and the secondary, much broader, peak between 2.5-3. The SWIFT obser-  
 488 vations (in red) capture a higher occurrence of young waves than the RHB obser-  
 489 vations or the WBF simulation. WBF also features a fatter tail of the distribution  
 490 toward larger wave ages, indicating that the model overemphasizes the occurrences  
 491 of swell and decaying waves compared to these observed estimates.

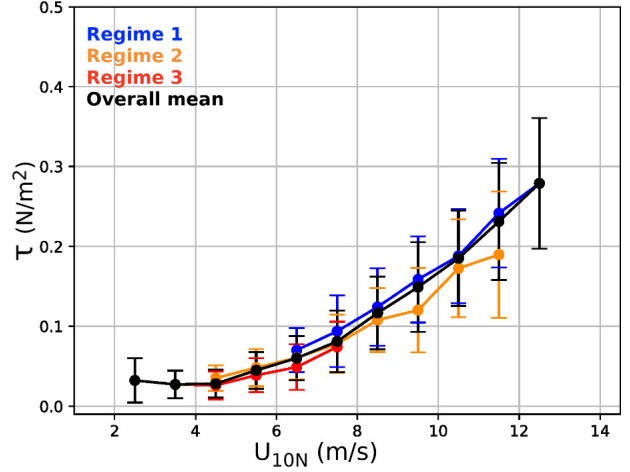
492 Given the wave age distributions, we then divide the distribution into 3 dif-  
 493 ferent “Regimes” to better understand the wave age-dependent  $z_0$ -wind speed and



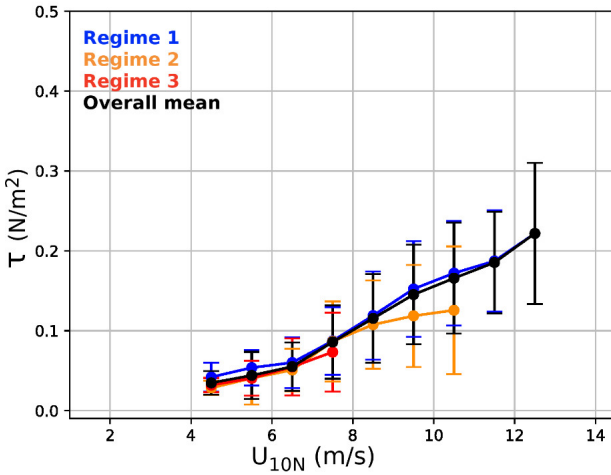
(a) wave age PDFs



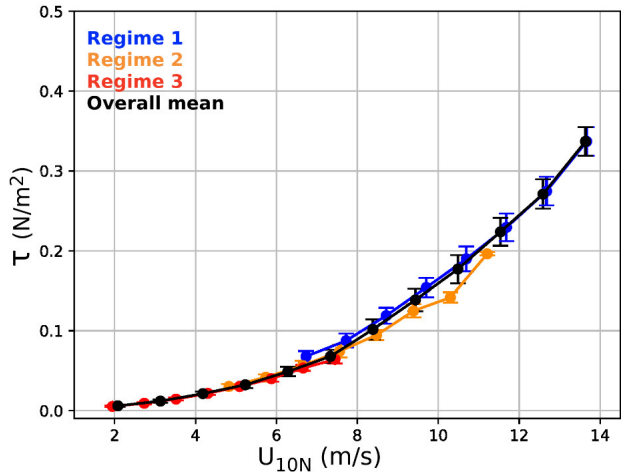
(b) RHB:  $\tau$  vs.  $U_{10N}$



(c) SWIFT:  $\tau$  vs.  $U_{10N}$



(d) WBF:  $\tau$  vs.  $U_{10N}$



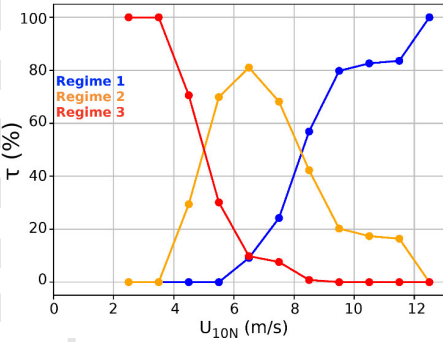
**Figure 8.** (a) Peak wave age distribution estimated from SWIFT (red), RHB (gray), and WBF (blue). Here, wave age is capped at 5. Three wave age regimes are defined: Regime 1 (blue) when wave age  $< 1.7$  denotes the young sea to fully developed sea, Regime 2 (orange) when wave age is between 1.7 and 3 indicates the mature to old sea, and Regime 3 (red) when wave age  $> 3$  represents the old sea and non-locally generated swell. (b-c) Binned scatter plots of  $\tau$  ( $\text{Nm}^{-2}$ ) vs.  $U_{10N}$  ( $\text{ms}^{-1}$ ), color-coded to show the three different wave age Regimes, with the bin-average of  $1 \text{ ms}^{-1}$ . The error bars represent  $\pm 1$  standard deviation. Only bins with more than 5 points are plotted. The mean of all wave ages is shown in black. (d) As in (b) and (c) except from the WBF run. Here WBF is sampled along-track of the RHB and SWIFT.

494  $\tau$ -wind speed relationships. Regime 1 refers to young to fully developed seas, defined  
 495 as when wave age  $< 1.7$ , while Regime 2 indicates the mature to old sea, including  
 496 mixed sea state, which is diagnosed as wave ages between 1.7 and 3. Finally, the old  
 497 sea and non-locally generated swell characterizes Regime 3 estimated as when wave  
 498 age  $> 3$ . When using the peak period, and to stay consistent throughout the paper,  
 499 thresholds are kept the same. However, these thresholds are not necessarily universal  
 500 but can vary in different times or regions under consideration.

501 The colored lines in Figures 8b and c show the bin-averaged surface stress  
 502 from the RHB and the SWIFT from the 3 Regimes. The black lines denote the  
 503 bin-averaged surface stress across all wave age regimes. Despite the significant error  
 504 bars, which represent  $\pm 1$  standard deviation, one can observe the consistent relation-  
 505 ship between the measured stress and the wind speed across different wave age.  
 506 For example, the measured stress over Regime 1 (blue) is higher than the overall  
 507 average (black) as the short-wind waves support the bulk of momentum exchanges.  
 508 In contrast, the stress over Regime 2 (orange) and Regime 3 (red) is lower than the  
 509 overall average, as the sea state is characterized by mixed and older seas. This sea  
 510 state dependence of wind stress is also somewhat evident in the WBF simulation  
 511 (Figure 8d) despite the smaller error bars likely due to smaller number of samples in  
 512 the model, as discussed in Section 2d.

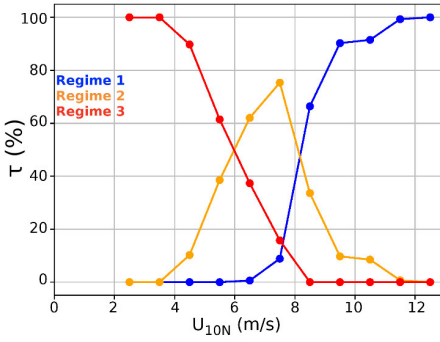
(a) RHB

$\tau$  vs.  $U_{10N}$



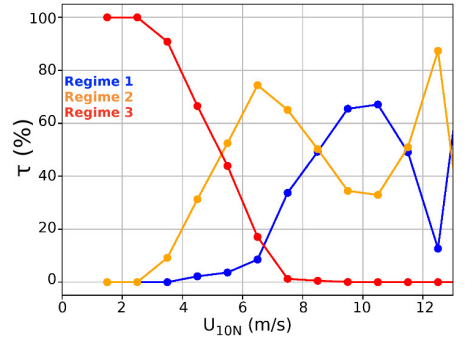
(b) WBF (along the RHB track)

$\tau$  vs.  $U_{10N}$



(c) WBF (Jan 8, 06h)

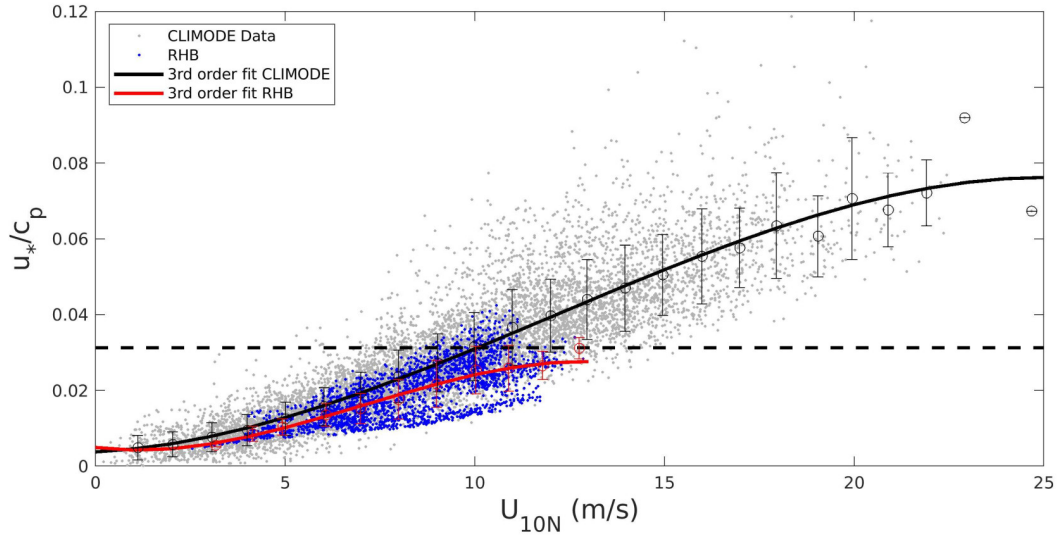
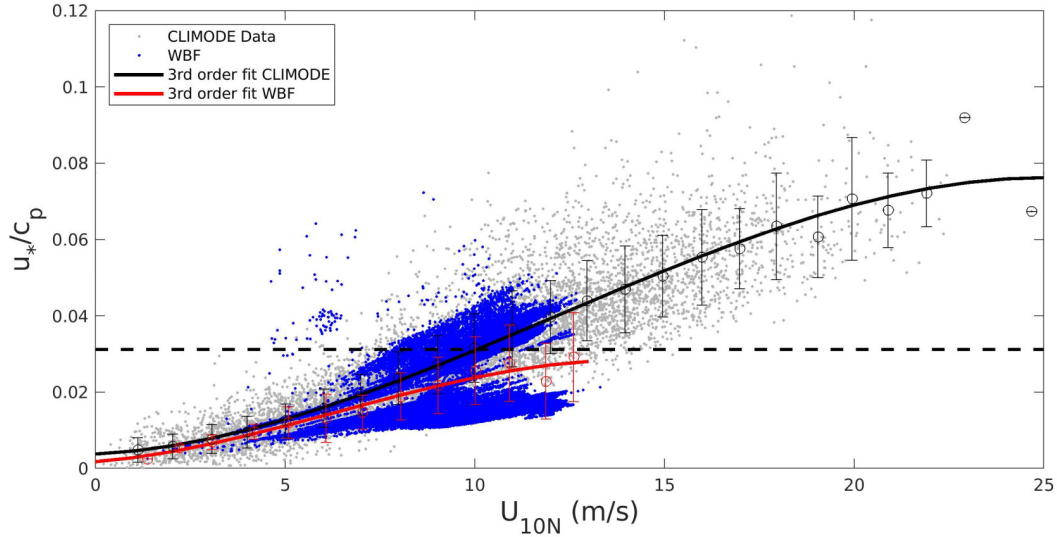
$\tau$  vs.  $U_{10N}$



**Figure 9.** Percentage contribution of  $\tau$  (%) by the three different wave age Regime at a given wind speed (bin averaged every  $1 \text{ ms}^{-1}$ ) from (a) RHB, (b) WBF sampled along the RHB track between January 9 and February 13, 2020 and (c) WBF sampled over the whole model domain on January 8, 2020 at 0600 UTC. The different colors denote the different wave age categories described in Figure 8.

513 To further quantify this relationship, Figure 9a shows the percentage of stress  
 514 supported by the different wave-age Regimes from the RHB observations, binned  
 515 over  $1 \text{ ms}^{-1}$  intervals. Under  $4 \text{ ms}^{-1}$  wind speeds, the surface stress is mainly  
 516 supported by Regime 3 (red), whereas above  $8 \text{ ms}^{-1}$ , Regime 1 (blue) dominates  
 517 the contribution to the stress. Regime 2, which represents mixed sea conditions  
 518 (orange), mainly supports the surface stress at low to moderate wind speeds ( $4\text{-}8$   
 519  $\text{ms}^{-1}$ ) and contributes to less than 20% of the stress above  $10 \text{ ms}^{-1}$ . Figure 9b  
 520 shows the same diagnostics, but for the WBF run sampled along the track of RHB.  
 521 It shows that the WBF overall exhibits a similar fractional contribution to stress.

522 When the model is compared to the observations at this particular track, WBF  
 523 appears to accurately characterize the observed stress relationship with wave age  
 524 (See also Figure 8). However, if sampled over a broader region of the same mixed  
 525 sea conditions from the model, a different result is obtained. Figure 9c shows the  
 526 same results as Figure 9b, except that the entire model domain is sampled under  
 527 the same synoptic condition examined in Section 3. It shows that the parameterized  
 528 stress under  $8\text{-}12 \text{ ms}^{-1}$  wind speeds supported by Regime 2 (orange) is comparable  
 529 to the stress supported by Regime 1 (blue) as also seen in Figure 6. In reality, short  
 530 wind waves under such wind speeds should still support the increased stress despite

(a)  $u_*/c_p$  vs.  $U_{10N}$ (b)  $u_*/c_p$  vs.  $U_{10N}$ 

**Figure 10.** (a) Scatter plot of inverse peak wave age ( $u_*/c_p$ ) vs.  $U_{10N}$  ( $m s^{-1}$ ) for CLIMODE data (gray) and RHB data (a, blue). Bin-averages with the 1 standard deviation error bars are overlaid, at  $1 m s^{-1}$  interval, along with the 3rd order fit (line) for CLIMODE (black) and RHB (red). The horizontal dashed line is  $u_*/c_p = 0.03$ , denoting the threshold for fully developed seas (equivalent to  $c_p/U_{10N} = 1.2$ ). (b) As in (a) but RHB data is replaced with WBF, for the whole domain

on January 8, 2020 at 0600 UTC.

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the higher wave age, we believe this is a form of deficiency in COARE3.5 WBF in representing the wind stress over mixed swell-dominated seas.

532

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In fact, the COARE3.5's WBF was developed and tuned primarily by using the wave data collected from the extratropics, where sea state tends to be dominated by growing and fully-developed waves under high winds (see Figure 2 in Edson et

534

535

al., 2013). Figure 10 compares the sea state used to tune COARE3.5, taken during the CLIMODE campaign (CLIVAR Mode Water Dynamic Experiment, Marshall et al., 2009), with the sea state observed by RHB during January-February 2020 and modeled in WBF on January 8, 2020 at 0600 UTC in the ATOMIC region. It shows the relationship between the inverse wave age and  $U_{10N}$ . Here, a low inverse wave age is indicative of decaying seas and swells. An inverse wave age of 0.03 (dashed line) is roughly equivalent to an equilibrium wave age of 1.2. As expected, the sea state captured in the ATOMIC region is very different and much older than the one used in COARE3.5. Therefore, the wind stress under moderate winds and swell dominated conditions observed here, and possibly in other tropical oceans, may not be currently well parameterized in the COARE3.5 WBF. The specific deficiency identified from this analysis is that, for mixed seas (Regime 2) where high wave age and moderately strong wind co-occur, the current COARE3.5 WBF overemphasizes the swell impact on wind stress, leading to the low-stress bias despite the moderately strong winds.

## 5 The revised wave-based formulation in COARE3.5

In the following, we present two experimental revisions to the  $z_0$  formulation in the current COARE3.5 WBF for swell conditions coincident with moderate to high winds, the condition that is frequently observed in the northern ATOMIC region in the boreal winter. One method is to replace the peak wave period ( $T_p$ ) with the mean wave period ( $T_m$ ) in the definition of the phase speed and thus wave age, and another is to incorporate the effect of misaligned waves with local wind on aerodynamic roughness in the  $z_0$  parameterization. In essence, these two observationally-guided approaches desensitize the impact of swell on  $z_0$  and  $\tau$  estimates at moderate winds and alleviate the low biases in the current COARE3.5 WBF. For this, we now return to the case study on January 8, 2020 as in Section 3.

### 5.1 The mean wave period

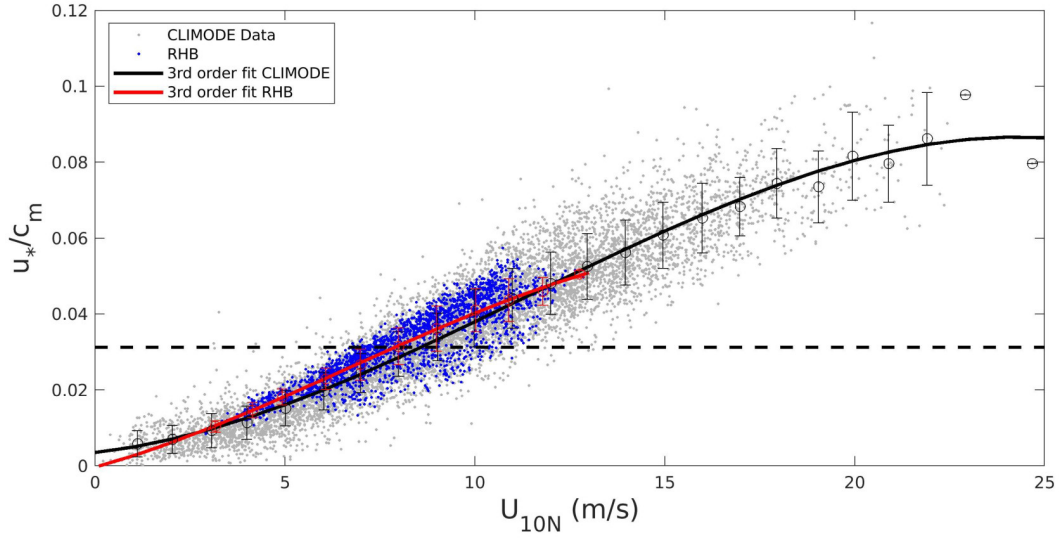
One possible approach to mitigate the overestimation of the swell impact on  $z_0$  and  $\tau$  under moderate to high winds is to use the wave's mean period,  $T_m$ , to calculate the average phase speed,  $c_m$ , in the wave age definition. This change is motivated by the finding that  $T_p$  does not accurately describe a mixed-sea state where swell and wind-sea co-exist, as shown in Figure 10.  $T_p$  can be also sensitive to the spectral shape of the wave energy and the chosen filter, while  $T_m$  can be reliably estimated from observations and WW3 as either an energy-weighted average period or zero-crossing period. A similar argument has been made recently by (Colosi et al., 2021) as they chose to use a wave age dependent computed with the mean period to construct the seasonal probability of swell over global oceans.

We carried out an additional coupled simulation, dubbed WBF- $T_m$ , where  $T_p$  is replaced with  $T_m$  to get the mean phase speed of the waves  $c_m$  in Eq. 12:

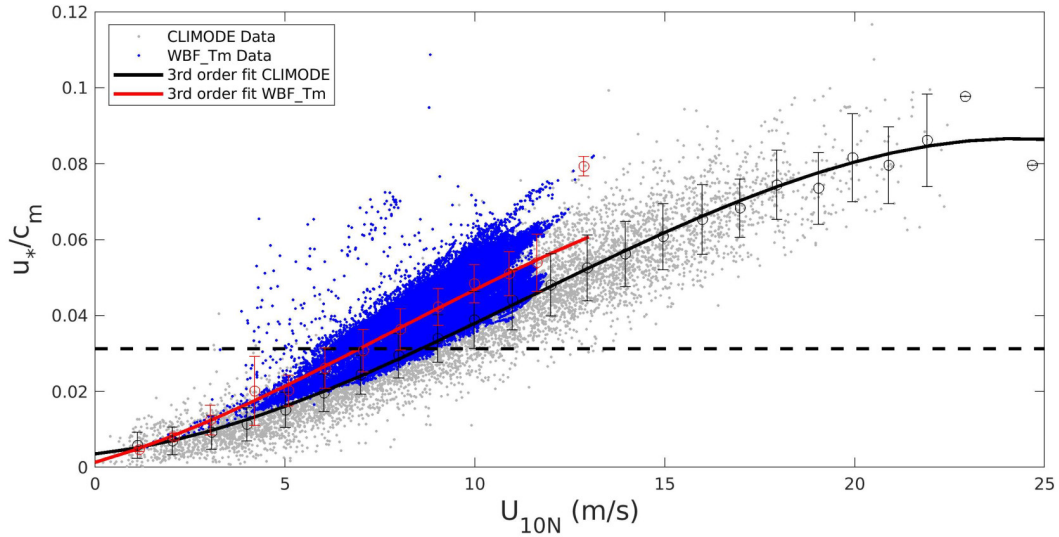
$$z_{rough} = H_s D \left( \frac{u_*}{c_m} \right)^B, \quad (12)$$

where  $D=0.39$  and  $B=2.6$ , which have been tuned using the COARE3.5 set of observations. We will estimate  $T_m$  based on the zero-crossing period, as it is the one used to describe  $T_m$  in the observation. Figure 11 shows the same diagnostics as in Figure 10 but this time using  $c_m$  to calculate the inverse wave age in both the observations, CLIMODE and RHB, and the WBF- $T_m$  run. The general trend of both sets of observations are now in good agreement (Fig. 11a). In WBF- $T_m$ , the use of  $c_m$  in eq. 12 alleviates the bias over the mixed sea (Regime 2) (Figure 10b vs. Fig. 11b) and shows a better agreement of the general trends from the observations.

(a)  $u_*/c_m$  vs.  $U_{10N}$



(b)  $u_*/c_m$  vs.  $U_{10N}$

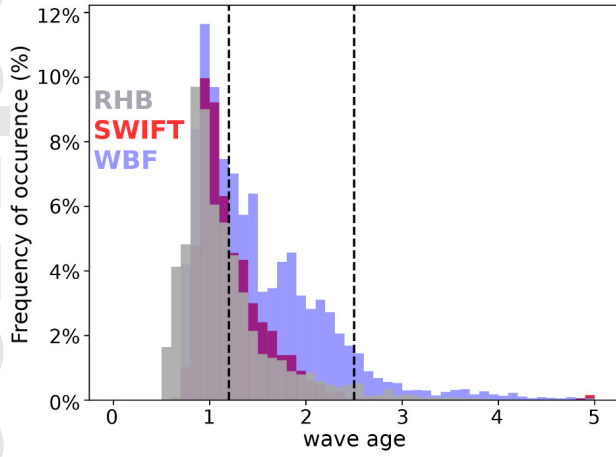


**Figure 11.** (a) As in Figure 10a, but with inverse mean wave age ( $u_*/c_m$ ). The dashed line is  $u_*/c_m = 0.03$ , denoting the threshold for fully developed seas (equivalent to  $c_m/U_{10N} = 1.2$ ). (b) As in Figure 10b except for showing the result from WBF  $T_m$

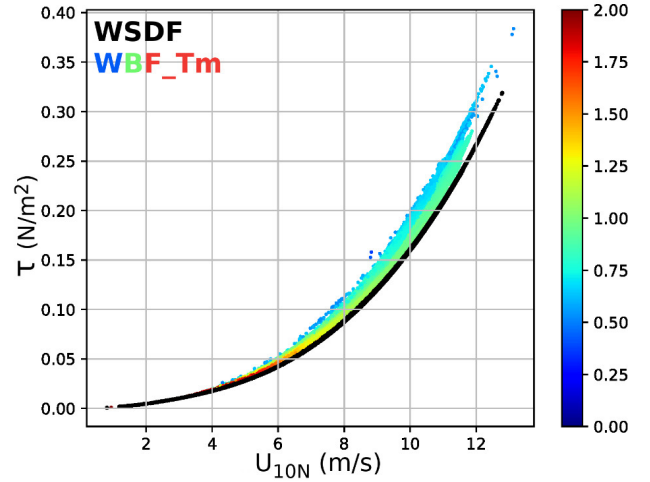
583 Further refinement of coefficients in eq. 12 will be addressed in more detail in the  
 584 future release of the COARE4.0 algorithm.

585 Figure 12a shows the PDF of wave age for RHB (gray), SWIFT (red), and  
 586 WBF  $T_m$  (blue) computed using  $T_m$ . This figure should be compared to Figure 8a  
 587 where RHB, SWIFT and WBF wave age PDFs were computed using  $T_p$ . Similar to  
 588 Figure 8a, wave age is capped at 5 to show the tail of the distribution. In contrast  
 589 to the bi-modal distribution of wave age with the pronounced secondary peak of  
 590 wave age estimate with  $T_p$ , the use of  $T_m$  effectively removes this secondary peak  
 591 in both the model and observations, yielding a markedly different distribution with  
 592 an overall prevalence of younger sea state. We adjusted the different categories of

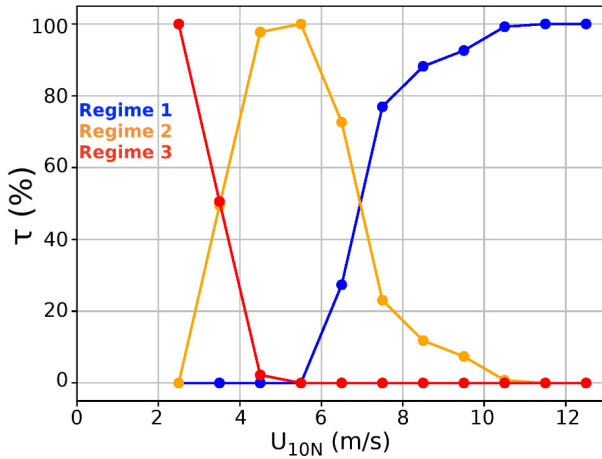
(a) wave age PDF



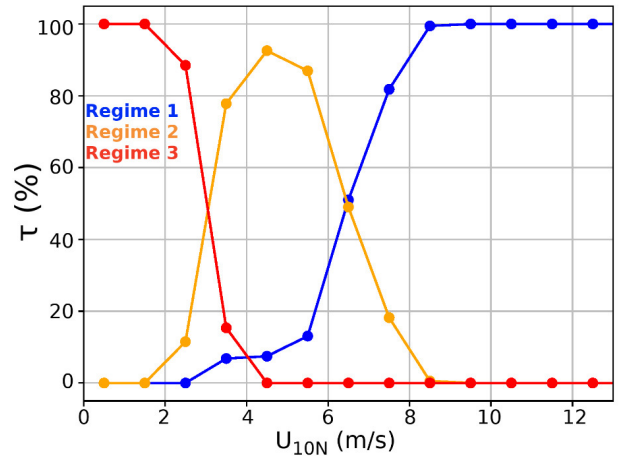
(b)  $\tau$  vs.  $U_{10N}$



(c) RHB:  $\tau$  vs.  $U_{10N}$



(d) WBF- $T_m$  (Jan. 8, 06h):  $\tau$  vs.  $U_{10N}$



**Figure 12.** (a) Mean wave age distributions estimated from RHB (gray), SWIFT (red), and WBF- $T_m$  (blue). WBF- $T_m$  is sampled along-track of the RHB and SWIFT. (b) Scatter plot of  $\tau$  ( $\text{Nm}^{-2}$ ) vs.  $U_{10N}$  ( $\text{ms}^{-1}$ ) from WSDF in black and WBF- $T_m$  color-coded to denote the corresponding wave age on January 8, 2020 at 0600 UTC. (c,d) As in Figure 9a,c, except that the wave age is defined with  $T_m$  for (c) RHB and (d) WBF- $T_m$ .

593 wave age defined previously to fit the new wave age distribution based on  $T_m$ . Figure  
 594 12b shows  $\tau$  on January 8, 2020 at 0600 UTC from WBF- $T_m$ , with wave age  
 595 color-coded. The cluster of low  $z_0$  with high wave age seen in Figure 4b is elimi-  
 596 nated in WBF- $T_m$ , because of the elevated  $z_0$  and  $\tau$  under moderate to high wind  
 597 speeds. Finally, Figure 12c,d, to be compared to Figure 9a,c shows the percentage of  
 598  $\tau$  supported by each category of wave age for RHB and for WBF- $T_m$ , respectively.  
 599 With the use of  $T_m$ , WBF- $T_m$  agrees well with RHB concerning the fractional con-  
 600 tribution from each sea state to the surface stress. Particularly over  $7 \text{ m s}^{-1}$ , most of  
 601 the contribution to  $\tau$  now comes from the wind sea (blue), whereas the contribution  
 602 of mature seas and swell subsides rapidly with the increased wind speeds. This is  
 603 a clear improvement from  $\tau$  parameterized using  $T_p$  (Figure 9c) and is much more  
 604 consistent with the observations (Figures 9a, 12c).

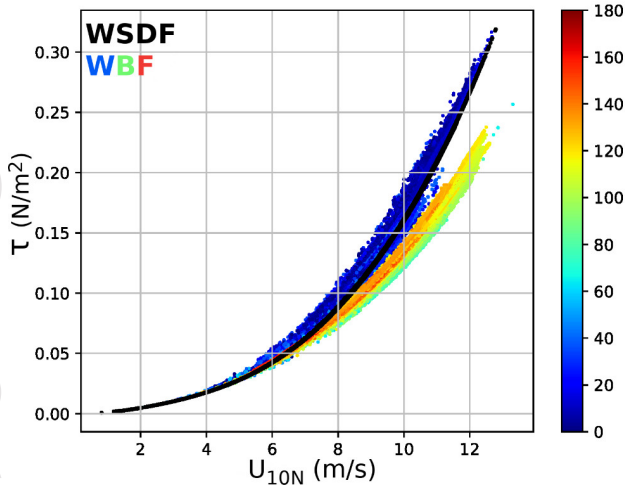
## 5.2 Including the (mis)aligned wind-wave directions

As discussed in Section 2, the COARE3.5 assumes the wave stress as a scalar roughness parameter, and hence the direction of wave-stress vectors is aligned with the mean wind vectors. However, wave stress and mean wind vectors can be misaligned under various conditions, including under rapidly translating storms (e.g., S. S. Chen et al., 2013), near strong vorticity and divergence gradients and density fronts (e.g., Villas Bas & Young, 2020), or over mixed seas where wind waves and swells co-exist under high winds. Such nonequilibrium wave motions can influence wave slope, roughness length, and wind stress (Janssen, 1991; Rieder et al., 1994; Zou et al., 2019; Patton et al., 2019; Porchetta et al., 2021; Deskos et al., 2021). Here, we attempt to incorporate the directionality of the wind and waves following Patton et al. (2019) and Porchetta et al. (2019), such that

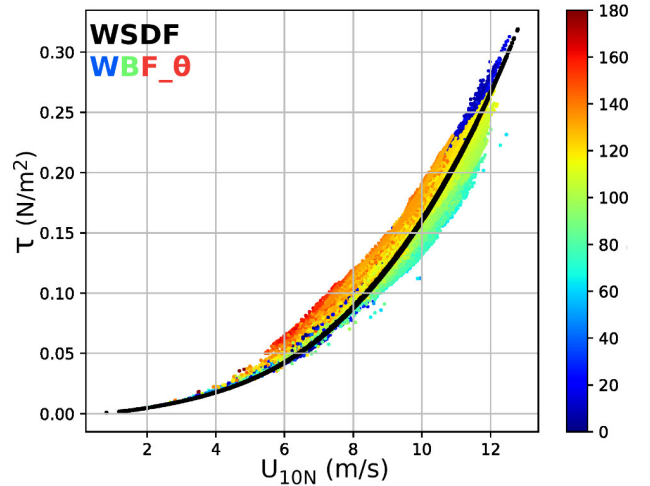
$$z_{rough} = H_s D \cos(a\theta) \left(\frac{u_*}{c_p}\right)^{B \cos(b\theta)}. \quad (13)$$

$D$  and  $B$  are the coefficients taken from COARE3.5 (See Eq. 8), while the coefficients  $a = 0.4$  and  $b = 0.32$  are adopted from (Porchetta et al., 2019). In principle, all these coefficients require site-specific tuning. For example, (Porchetta et al., 2019) used the high wind conditions observed from the FINO platform in the North Sea and the Air-Sea Interaction Tower (ASIT) in the New England Shelf, which represents different wind speed and wave age conditions from the trade-wind and swell-dominated tropical oceans as in the ATOMIC domain. Additional tuning exploiting direct momentum flux measurements would be needed to develop a refined set of coefficients for the tropical oceans. This is beyond the scope of the study. Using this new formulation, we conducted an additional coupled experiment, dubbed  $WBF_\theta$ , which is to be compared to the default wave-based formulation in COARE3.5, where  $\theta = 0$ .

(a)  $\tau$  vs.  $U_{10N}$



(b)  $\tau$  vs.  $U_{10N}$

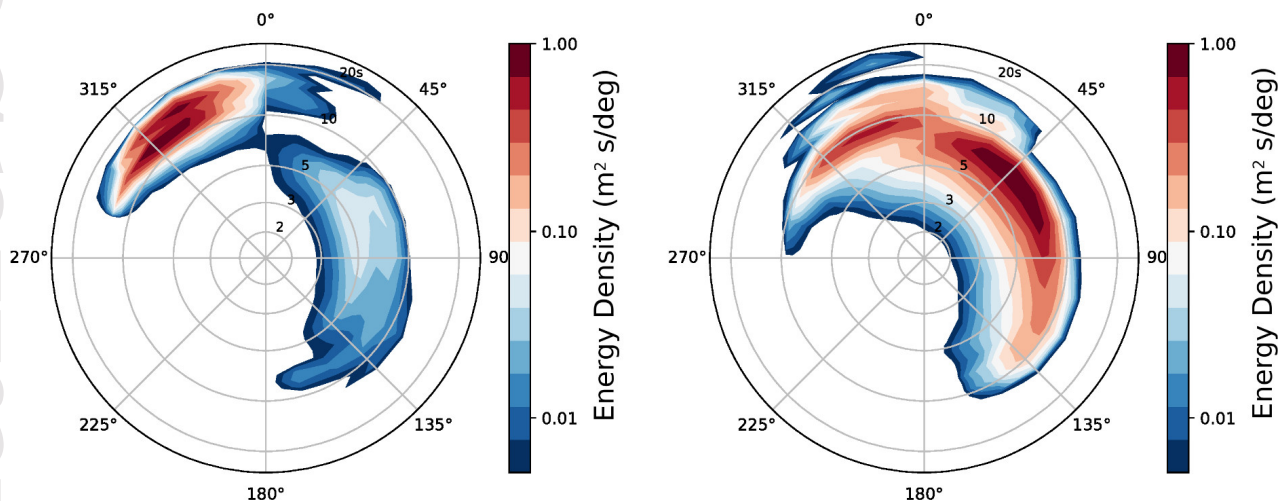


**Figure 13.** (a) Scatter plot of parameterized  $\tau$  ( $\text{Nm}^{-2}$ ) vs.  $U_{10N}$  ( $\text{ms}^{-1}$ ) from WSDF in black and WBF color-coded to denote the corresponding wind-wave angle ( $\theta$ ) on January 8, 2020 at 0600UTC. Note that in the  $z_0$  formulation in WBF assumes  $\theta = 0$ . (b) As in (a) except from  $WBF_\theta$ , where  $\theta$  is treated as a non-zero quantity in the  $z_0$  formulation.

629 Figure 13a compares the parameterized  $\tau$ , color-coded by the angle ( $\theta$ ) be-  
 630 tween the wind direction and peak wave direction in WBF. It shows that the lower  
 631  $\tau$  from WBF compared to WSDF (and also observations) occurs when the swell  
 632 waves are strongly misaligned with winds (e.g.,  $\theta > 60-90^\circ$ ). This indicates that the  
 633 assumption of  $\theta = 0$  in WBF can be attributed to the lower  $\tau$ . When the directional  
 634 misalignment is considered in the roughness length parameterization in COARE3.5  
 635 (Figure 13b),  $\tau$  over the misaligned waves has been effectively elevated as the waves  
 636 opposing the wind increase the surface drag. This is shown to reduce the low  $\tau$  bias  
 637 significantly.

(a) Location: (54°W - 16°N) ;  $\chi=2.1$

(b) Location: (46°W - 6°N) ;  $\chi=1.1$



**Figure 14.** Normalized wave spectrum energy density ( $\text{m}^2\text{sdeg}^{-1}$ ) plotted in period (s) space from (a) one point in the northern part of the domain under swell influence and (b) one point in the center part of the domain on January 08, 2020 at 0600UTC for WBF.

638 Here, the alignment between wind and waves has been defined only by using  
 639 the wave peak direction. Figure 14 compares the normalized wave spectrum energy  
 640 density ( $\text{m}^2\text{sdeg}^{-1}$ ) shown in the period space between one grid point in the north-  
 641 ern part of the domain under swell regime (Fig. 14a) and another grid point in the  
 642 center part of the domain under wind waves regime for WBF. Both are sampled  
 643 on January 08, 2020 at 0600UTC. On the northern grid point where the wave age  
 644 was 2.1, Figure 14a shows the strong swell signal (with the periods of 10-20s) from  
 645 the northwest direction. It does also show a large directional spreading, due to the  
 646 concurrent shorter period wind waves (2-10s) originating from the northeast, east,  
 647 and southeast direction. However, the energy density from the shorter-period waves  
 648 is much weaker. In the center of the domain (Figure 14b), where the sea state is  
 649 dominated by wind-waves and waves near equilibrium (the wave age here is 1.1), the  
 650 directional spreading is also quite large, but with higher energy in the wind waves  
 651 and weaker energy in the swell.

652 The sea state in this region appears to be mixed ubiquitously between wind  
 653 waves and swell in winter, leading to a large wave directional spreading. However,  
 654 since the peak energy density is well separated between the swell (in the northern  
 655 point, Fig. 14a) and the wind waves (in the southern point, Fig. 14b), we antic-  
 656 ipate that the use of waves' direction variance in the bulk formula or the spectrally-  
 657 averaged wave direction in the bulk formula, would yield qualitatively similar re-



sults. For this reason, in the present study, only the peak direction of the waves is used to account for the misaligned wave effect on  $z_0$  in COARE. However, it is possible that by using the peak wave direction we would grossly underrepresent some unresolved processes contributing to the directional spread of waves, and its impact on  $z_0$ .

## 6 Conclusion

This study investigated the role of surface waves in surface roughness length ( $z_0$ ) and surface stress ( $\tau$ ) in the persistent and strong trade winds and swell-dominated Northwestern Tropical Atlantic Ocean during the boreal winter season. The main objective is to evaluate how accurately the air-sea momentum flux is represented in advanced bulk flux algorithms such as COARE3.5 when compared to the direct surface flux measurements. In this investigation, estimated  $z_0$  and  $\tau$  from four different SCOAR ocean-atmosphere-wave coupled model simulations are analyzed. The results show that the estimated  $z_0$  and  $\tau$  differences strongly depend on wind speeds and wave age regimes. Wind sea or fully-developed sea under high winds are characterized by the enhanced wave slope and choppy surface (Figure 5b), which effectively increases the surface drag, and  $\tau$ . The increased surface drag decelerates the near-surface winds (Figure 6c).

However, in the mixed sea condition, where moderate to high wind speeds (10 to  $12 \text{ ms}^{-1}$ ) co-occur with decaying swell, the WBF tends to underestimate  $z_0$  compared to the WSDF and  $\tau$  compared to the measurements. The weak stress then accelerates the near-surface wind speed by 5% over the region of negative change in wind work (Figure 6d). The sea state, in this high wave age region, is strongly misaligned with the local wind (Figure 5d), indicating the presence of remotely-generated swell. However, despite the swell-dominated sea state, the observations suggest that the wind seas in this mixed sea condition should continue to support the momentum flux due to moderate-to-high wind speeds, thereby increasing  $\tau$  with wind speed (Figure 7).

The different approaches were explored in this study to alleviate the low-stress bias in the COARE3.5 WBF under the mixed sea regime. The first approach involves re-defining wave age using the mean period of the waves to more accurately represent the wave period in the mixed sea condition (Figure 4a). The second approach takes advantage of the fully coupled model by considering the directionality of waves with respect to winds (Eq. 12), the vital missing process in the current COARE3.5 WBF and many numerical modeling studies except for a limited number of Large Eddy Simulations (LES) and offshore wind energy studies (See Review by Patton et al., 2019). Our results show that both approaches produce equivalent results by effectively boosting  $z_0$  and  $\tau$  under the misaligned waves under moderate-to-high winds. Since both methods yield equivalent results, accounting for both (peak direction and wave mean period), without more dedicated tuning with the measurements, produces too strong correction for the low bias (not shown). Finally, it is important to note that these improvements are most likely to be site-dependent, as we are only using limited observations in one specific region. Moreover, the improvement of the parameterization is mostly over specific regimes of wind and waves where the original parameterization was deficient.

Our analysis reveals a notable deficiency in the ocean-wave and wave-atmosphere coupling components of the coupled model, which guides the direction of our future investigation. That is, the frequency of swell simulated by the coupled WW3 model is overestimated compared to the in situ observations (Figure 8a), more so with the use of peak wave period but nonetheless noticeable with the use of mean period. Since the wave model provide the parameters required by the WBF, some

709 of the issues described above are a result of inaccurate inputs as well as problems  
710 with the parameterization. The tendency toward the higher wave age indicates that  
711 the model under-represents critical dissipation mechanisms of the swell energy, and  
712 waves in general, which likely have contributed to the low-stress bias. There are at  
713 least two possible factors to consider.

714 First, the primary loss of swell energy is to the atmosphere in situations where  
715 the swell waves outrun the winds or propagate in the opposite direction to the local  
716 wind (e.g., M. Donelan, 1999; Raschle et al., 2008; Kahma et al., 2016; Liu et al.,  
717 2017). Tropical oceans, including our study region, have many low-wind regimes,  
718 where the wave-driven low-level wind jet (Harris, 1966) and turbulent mixing in  
719 the MABL (Kantha, 2006; Ardhuin & Jenkins, 2006; A. V. Babanin, 2006) consti-  
720 tute important sources for attenuation of the swell energy (Ardhuin et al., 2009;  
721 S. Chen et al., 2019). It is quite possible that the processes related to the upward  
722 flux of momentum and energy over swell are not adequately captured in our coupled  
723 wind-wave model. Previous studies find that the wave-driven wind jet is at heights  
724 of 5-10 m (Sullivan et al., 2008; Smedman et al., 2009). However, our experiments  
725 used the default vertical grid system in WRF, where the wind at the lowest height  
726 of the model is typically 30–50 m. The WRF PBL scheme expects this level to be  
727 within the constant-flux layer, where similarity theory is applied (Aligo et al., 2009;  
728 Shin et al., 2012). Yet, this level can be above the surface layer, especially in the  
729 low-wind and stable boundary layer conditions, as often observed in the northern  
730 part of the ATOMIC domain. If the turbulent mixing between the lowest model  
731 level and the swell at the sea surface is weak, the upward energy and momentum  
732 fluxes from the swell to the wind are likely to be under-represented. This might have  
733 been exacerbated by using a local PBL scheme (MYNN) in our model.

734 Moreover, parameterizations for the so-called negative wind input exist in  
735 standalone WW3 model through the use of the source term packages of wind input  
736 (M. A. Donelan et al., 2006; Ardhuin et al., 2010; A. Babanin, 2011; Rogers et al.,  
737 2012; Liu et al., 2017, 2019). With this, the standalone WW3 model forced with  
738 winds should better capture the loss of energy of swell waves. Yet, it is unclear how  
739 such parameterizations should be incorporated into the coupled model, as they do  
740 not represent the actual gain of momentum by the wind from the swell. Our future  
741 work will focus on adequately representing the near-surface wind responses to swell  
742 waves in the atmospheric model.

743 Secondly, the wave breaking and the induced near-surface mixing would in-  
744 fluence the wave energy growth and attenuation (e.g., Kudryavtsev et al., 2014).  
745 Also, Iyer et al. (2022), using the SWIFT drifters deployed during the ATOMIC  
746 campaign, showed that wave-current interactions can generate significant spatial and  
747 temporal variability in momentum fluxes in this region. However, here, since the  
748 current study does not include wave-ocean coupling, the question about the impacts  
749 of ocean-wave coupling on the skill of the simulated wave fields cannot be addressed.  
750 This is a subject of ongoing efforts.

## 751 **7 Open Research**

752 The observational datasets from the ATOMIC and EUREC<sup>4</sup>A experiments  
753 (Stevens et al., 2021) are available freely on [https://observations.ipsl.fr/  
754 aeris/eurec4a/\#/](https://observations.ipsl.fr/aeris/eurec4a/\#/). ERA5 Atmospheric hourly reanalyses were made avail-  
755 able by the Copernicus Climate Change Service (Hersbach et al., 2018a, 2018b).  
756 Mercator Ocean International daily analyses (Lellouche et al., 2018) were  
757 made available by the Copernicus Marine Environment Monitoring Service on  
758 <https://doi.org/10.48670/moi-00016>. Global 3-hourly spectral wave analy-  
759 ses were made available by Ifremer (Raschle & Ardhuin, 2013) on a FTP server at

760 ftp://ftp.ifremer.fr/ifremer/ww3/HINDCAST/GLOBAL; WaveWatchIII model (The  
 761 WAVEWATCH III Development Group, 2016) is available at [https://github.com/](https://github.com/NOAA-EMC/WW3)  
 762 NOAA-EMC/WW3. WRF model (Skamarock et al., 2008) is available at [https://](https://github.com/wrf-model/WRF)  
 763 [github.com/wrf-model/WRF](https://github.com/wrf-model/WRF). ROMS model (Shchepetkin & McWilliams, 2005) is  
 764 also freely available at <https://github.com/kshedstrom/roms>. The SCOAR (Seo  
 765 et al., 2007) code is available at <https://github.com/hyodae-seo/SCOAR>. Finally,  
 766 the original versions of COARE3.5 (Edson et al., 2013) bulk formula is available at  
 767 <https://github.com/NOAA-PSL/COARE-algorithm>.

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