

Reexamining the Estimation of Tropical Cyclone Radius of Maximum Wind from Outer Size with an Extensive Synthetic Aperture Radar Dataset

ARTHUR AVENAS¹,^{a,b} ALEXIS MOUCHE,^a PIERRE TANDEO,^b JEAN-FRANCOIS PIOLLE,^a DAN CHAVAS,^c RONAN FABLET,^b JOHN KNAFF,^d AND BERTRAND CHAPRON^a

^a Ifremer, Univ. Brest, CNRS, IRD, Laboratoire d'Océanographie Physique et Spatiale (LOPS), IUEM, Plouzané, France

^b IMT Atlantique, Lab-STICC, Université Bretagne Loire, Brest, France

^c Department of Earth, Atmospheric, and Planetary Sciences, Purdue University, West Lafayette, Indiana

^d NOAA/NESDIS Regional and Mesoscale Meteorological Branch, Fort Collins, Colorado

(Manuscript received 14 June 2023, in final form 6 September 2023, accepted 12 October 2023)

ABSTRACT: The radius of maximum wind R_{\max} , an important parameter in tropical cyclone (TC) ocean surface wind structure, is currently resolved by only a few sensors so that, in most cases, it is estimated subjectively or via crude statistical models. Recently, a semiempirical model relying on an outer wind radius, intensity, and latitude was fit to best-track data. In this study we revise this semiempirical model and discuss its physical basis. While intensity and latitude are taken from best-track data, R_{\max} observations from high-resolution (3 km) spaceborne synthetic aperture radar (SAR) and wind radii from an intercalibrated dataset of medium-resolution radiometers and scatterometers are considered to revise the model coefficients. The new version of the model is then applied to the period 2010–20 and yields R_{\max} reanalyses and trends that are more accurate than best-track data. SAR measurements corroborate that fundamental conservation principles constrain the radial wind structure on average, endorsing the physical basis of the model. Observations highlight that departures from the average conservation situation are mainly explained by wind profile shape variations, confirming the model's physical basis, which further shows that radial inflow, boundary layer depth, and drag coefficient also play roles. Physical understanding will benefit from improved observations of the near-core region from accumulated SAR observations and future missions. In the meantime, the revised model offers an efficient tool to provide guidance on R_{\max} when a radiometer or scatterometer observation is available, for either operations or reanalysis purposes.

KEYWORDS: Hurricanes/typhoons; Tropical cyclones; Wind; Remote sensing; Satellite observations; Statistical techniques

1. Introduction

Estimating tropical cyclone (TC) ocean surface wind structure is challenging but crucial for several applications. In particular, TC surface wind spatiotemporal distributions are used as input to surface wave studies (Wright et al. 2001; Young 2017; Kudryavtsev et al. 2021), storm surge studies (Irish et al. 2008; Takagi and Wu 2016), or the upper ocean responses to TC passages (Price 1981; Ginis 2002; Kudryavtsev et al. 2019; Combot et al. 2020b). In such studies, the radius of maximum winds (hereinafter R_{\max}) is a critical parameter that significantly affects wave developments, surge estimates, sea surface height, temperature, and salinity variations within the TC wakes. Most parametric surface wind fields, often used for those applications, assume that R_{\max} is known (Holland 1980; Willoughby et al. 2006). Thus, R_{\max} errors cascade into errors for the entire spatial distribution of wind speeds. For instance, a Rankine profile may be defined as

$$V_{\text{Rankine}}(r) = \begin{cases} V_{\min} + (V_{\max} - V_{\min}) \left(\frac{r}{R_{\max}} \right) & \text{if } r \leq R_{\max} \\ V_{\min} + (V_{\max} - V_{\min}) \left(\frac{R_{\max}}{r} \right) & \text{if } r > R_{\max} \end{cases} \quad (1)$$

Figure 1a shows a comparison between two Rankine profiles for two different R_{\max} values representative of TC Lane, a

northeastern Pacific Ocean hurricane that reached category 5 on the Saffir–Simpson scale in 2018. TC Lane's wind speeds were estimated by a swath of satellite-based SAR observation at 0438 UTC 23 August (Fig. 1b). From the SAR wind speeds, the azimuthally averaged wind profile can be derived (dashed green curve in Fig. 1a). The inferred R_{\max} is 15 km, which is smaller by approximately a factor of 2–3 than the 37-km value interpolated to the SAR acquisition time in the best-track data (Knapp et al. (2010); hereinafter IBTrACS). Such a mismatch between best-track and SAR R_{\max} estimates is representative of what has been reported in the literature (Combot et al. 2020a). In the present case (Fig. 1a), this discrepancy results in a mean absolute error (MAE) as high as 28 m s^{-1} near the eyewall region when using subsequent Rankine profile estimates.

To date, airborne Stepped Frequency Microwave Radiometer (SFMR) surface winds (Uhlhorn et al. 2007) provide means to estimate R_{\max} . Yet, airborne measurements have limited azimuthal coverage and are operated over only few ocean regions and events. From a satellite perspective, high-spatial-resolution estimates of TC ocean surface wind field are now more systematically carried out, especially from SAR dedicated acquisitions (Mouche et al. 2017; Combot et al. 2020a). More reliable R_{\max} estimates are then obtained for all ocean basins, though with limited spatiotemporal sampling. Presently, the most often available spaceborne observing systems, capable of probing the ocean surface during TC conditions, are the combined capabilities from active

Corresponding author: Arthur Avenas, rthur.avenas@ifremer.fr

DOI: 10.1175/MWR-D-23-0119.1

© 2023 American Meteorological Society. This published article is licensed under the terms of the default AMS reuse license. For information regarding reuse of this content and general copyright information, consult the AMS Copyright Policy (www.ametsoc.org/PUBSReuseLicenses).

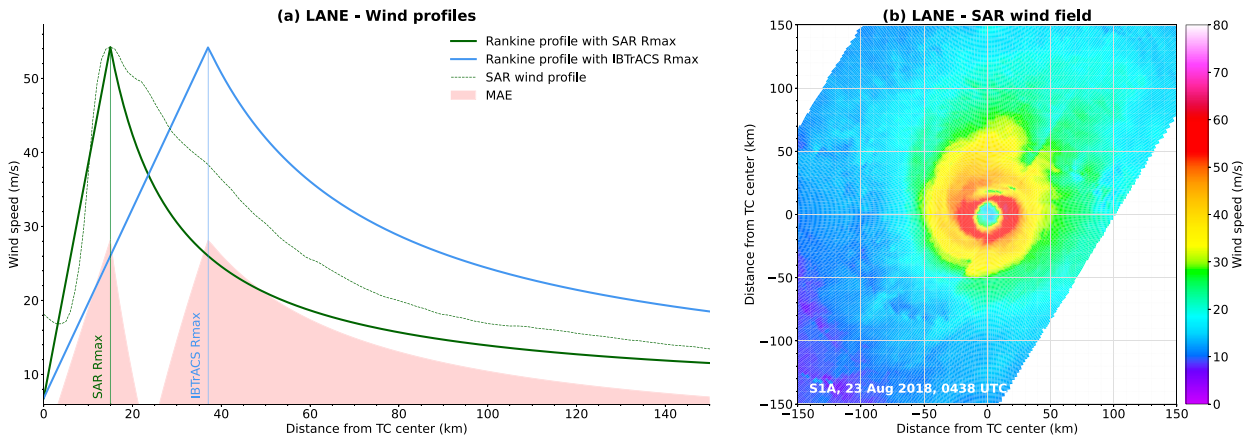


FIG. 1. (a) Comparison between two Rankine profiles inspired by (b) the SAR acquisition over TC Lane at 0438 UTC 23 Aug 2018. Rankine profiles are defined with SAR R_{\max} (15 km; solid green) or IBTrACS R_{\max} (37 km; solid blue) and the same V_{\max} (54 m s^{-1}) and V_{\min} (7 m s^{-1}), consistently with the SAR azimuthally averaged profile (dashed green). MAE between the two Rankine profiles is shaded in red.

scatterometers and passive radiometers (Quilfen et al. 2007). Relative to radiometers, scatterometers generally have an improved medium spatial resolution. Yet, the strong gradients of the surface wind existing at scales of a few kilometers may still be too smoothed to precisely locate the wind maxima, and the position of the center (Quilfen et al. 1998). In addition, scatterometers, especially those operating at Ku band and higher microwave frequencies, can suffer from rain contamination. Signal sensitivity at high winds (above hurricane force wind: 33 m s^{-1}) has also been questioned (Donnelly et al. 1999; Mouche et al. 2019). Radiometer measurements may be less impacted by rain, especially those operating at L band (Reul et al. 2012, 2017), and are demonstrated to still be highly sensitive above hurricane force winds. However, actual spaceborne radiometers operating at L or C band have a lower spatial resolution. High wind speed gradients near the R_{\max} region for the most-intense TCs are then generally indistinct. Direct estimates of R_{\max} using scatterometers or radiometers are thus difficult to perform and are possibly limited to particular large-storm cases.

More indirect means to infer R_{\max} were also considered. Both Mueller et al. (2006) and Kossin et al. (2007) used geostationary infrared satellite data. For the cases where a clear eye is well-defined on the infrared image, using linear regression to estimate R_{\max} results in an MAE of only $\sim 5 \text{ km}$ when compared with aircraft-based estimates. Under less favorable conditions, R_{\max} can still be estimated via multiple linear regression in combination with a principal components analysis, but leads to a degraded MAE of $\sim 20 \text{ km}$. Notably, for the clear-eye case, Tsukada and Horinouchi (2023) trained the linear regression with available SAR R_{\max} estimates and improved the method, decreasing the MAE to $\sim 2 \text{ km}$.

In the absence of infrared data, a rough R_{\max} estimate can also be obtained, considering the storm intensity and latitude known, as evidenced by Willoughby et al. (2006) and Vickery and Wadhera (2008). Indeed, following the angular momentum

conservation, R_{\max} must decrease when the intensity increases. On average, such a physical constraint agrees well with observations [see for instance Fig. 9 in Combet et al. (2020a)]. In addition, it is also known that R_{\max} increases with latitude (e.g., Willoughby and Rahn 2004), another consequence of angular momentum conservation along with the decrease of intensity with latitude. Solely using intensity and latitude to predict R_{\max} yields a root-mean-square error of order $\sim 20 \text{ km}$. Results from Vickery and Wadhera (2008) show that in several cases, the observed R_{\max} is inconsistent with the general principle of angular momentum conservation. This suggests that R_{\max} natural variability can hardly be captured by such simple statistical models.

More recently, Chavas and Knaff (2022)—hereinafter CK22—suggested to use information on the TC outer size in combination with latitude and intensity. In the CK22 framework, R_{\max} is estimated from the TC intensity V_{\max} , the radius of gale R_{34} [i.e., the maximum radial extent of the 34-kt winds ($1 \text{ kt} \approx 0.51 \text{ m s}^{-1}$)] and the Coriolis parameter, defined as $f = 2\Omega \sin(\phi)$, where $\Omega = 7.292 \times 10^{-5} \text{ s}^{-1}$ is Earth's angular velocity and ϕ is the latitude of the TC center. Such an approach is practical, especially because R_{34} is well estimated by satellite scatterometers and radiometers (Brennan et al. 2009; Chou et al. 2013; Reul et al. 2017). In fact, R_{34} estimates are routinely produced for every TC and are included in IBTrACS.

The CK22 framework is based on physical understanding of the radial wind structure (Emanuel 2004; Emanuel and Rotunno 2011) and phrased in terms of absolute angular momentum $M(r) = rV + (1/2)fr^2$, where f , r , and V are the Coriolis parameter, the radius, and the tangential wind speed of an air parcel, respectively. If the ratio

$$\frac{M_{\max}}{M_{34}} := \frac{M(r = R_{\max})}{M(r = R_{34})}$$

is prescribed, one can then estimate R_{\max} provided estimates for the three abovementioned parameters using

$$R_{\max} = \frac{V_{\max}}{f} \left(\sqrt{1 + \frac{2fM_{\max}}{V_{\max}^2}} - 1 \right). \quad (2)$$

CK22 fitted a log-linear regression model to estimate the ratio M_{\max}/M_{34} with the two predictors $X_{34}^{(1)} := (V_{\max} - 17.5 \text{ m s}^{-1})$ and $X_{34}^{(2)} := (V_{\max} - 17.5 \text{ m s}^{-1})[(1/2)/R_{34}]$.

It is tempting to use this framework in combination with best-track data. CK22 used best-track estimates (in a region west of 50°W) of R_{\max} , V_{\max} , R_{34} , and latitude to fit the log-linear regression model. As a result, their model inherited best-track biases. In particular, the reported R_{\max} overestimation in best tracks relative to SAR (Combrot et al. 2020a) translated into an overestimation of the ratio M_{\max}/M_{34} during the regression training, further leading to overestimated R_{\max} values.

The quality of R_{34} best-track estimates has also already been questioned (Sampson et al. 2017). This parameter is re-analyzed and compiled in IBTrACS since 2004 for the North Atlantic and northeastern Pacific Oceans and since 2016 for the northwestern Pacific (Knaff et al. 2021). Yet, surveying specialists who produce best tracks in the Atlantic Ocean (Landsea and Franklin 2013) are on average much less confident in their wind radii estimates ($\sim 25\%$ – 50% of relative uncertainty) than in their intensity estimates ($\sim 10\%$ – 20%).

In addition, best tracks may also suffer from temporal and spatial heterogeneities (Schreck et al. 2014; Wang and Toumi 2021). Indeed, the reanalysis methodology depends on the available data at each reanalysis time: best-track estimates of TC events covered by aircraft data are for instance more trustworthy (Landsea and Franklin 2013). Reanalysis is also subjective, with individual specialists from each agency or regional specialized meteorological center conducting their own weighting of the available observations. Furthermore, best tracks are finalized annually and are not updated with evolving reanalysis methodology, creating a temporal discontinuity in the final IBTrACS database.

A possible limitation of the CK22 approach is the arbitrary choice of the outer wind radius R_{34} . Indeed, their model could well be trained using R_{50} or R_{64} . In CK22, the choice of R_{34} was motivated by the fact that best-track estimates of R_{50} and R_{64} are generally more uncertain than R_{34} estimates. With more-reliable R_{50} and R_{64} estimates, possibly obtained from radiometers or scatterometers, one could assess whether using these wind radii would improve the CK22 model.

The physical basis for wind structure relationships such as CK22 is a long-running issue. The assumption that an outer wind radius partly constrains the wind structure dates back to Riehl (1963). Riehl (1963) used a two-layer conceptual model constrained by an angular momentum conservation in the outflow and a potential vorticity (PV) conservation in the inflow layer. Riehl (1963) could then derive a relationship between R_{\max} , V_{\max} , f , and an outer radius R_{out} , corresponding to a distance at which the outflow velocity vanishes. Later, Kalashnik (1994) considered the Holland parametric profile (Holland 1980) within a theoretical framework, to analyze the dependence of the near-core wind structure on the wind profile. Emanuel and Rotunno (2011) also derived an analytical

solution for the near-core wind profile based on an assumption on the outflow temperature.

While these studies offer theoretical guidance, these theoretical inferences of R_{\max} are difficult to apply in practice. Indeed, most actual sensors fail to capture the wind profile shape used in Kalashnik (1994), while the model of Emanuel and Rotunno (2011) relies on parameters that are difficult to evaluate. Following Riehl (1963), the theoretical outer radius R_{out} is unknown and cannot be specified to correspond to a given surface wind speed.

Building on the above considerations, the aim of this study is twofold. First, the CK22 model is revised using SAR R_{\max} estimates, different wind radii (referring hereinafter to R_{34} , R_{50} , and R_{64}) estimated on intercalibrated radiometers and scatterometers, and intensity and latitude best-track estimates. Second, the physical basis of the CK22 model is further assessed through an examination of conservation equations and a thorough analysis of the SAR database.

The data used in the present work are introduced in section 2 and further analyzed in section 3. Then, the CK22 model is revised and its performance is assessed in section 4. The physical basis of the model is discussed with respect to SAR observations in section 5. Concluding remarks and possible routes for future investigations are provided in section 6.

2. Data

In the present work, different radiometer and scatterometer data (Table 1) over the period 2010–20 were used to estimate wind radii (R_{34} , R_{50} , and R_{64}), whereas SAR data (Table 2) were used to estimate the R_{\max} values required to fit the CK22 log-linear model. Furthermore, IBTrACS provided intensity and latitude estimates (V_{\max} and f).

We used different radiometer and scatterometer missions to constitute the most extensive dataset of R_{\max} reanalyses. These sensors rely on different physical principles (passive or active sensors), and have different frequencies (L band, C band, or Ku band) and spatial resolutions. To ensure homogeneity of the wind radii estimates, we used radiometer and scatterometer winds intercalibrated by Portabella et al. (2022). Note that surface wind speed estimates from the Cyclone Global Navigation Satellite System (CYGNSS) do not belong to this intercalibrated dataset. Indeed, even though a level-3 stormcentered gridded wind speed product has recently been developed to improve former CYGNSS wind speed retrievals, its capacity to correctly inform the TC surface wind structure, especially R_{50} and R_{64} , remains to be assessed and validated (Morris and Ruf 2017; Krien et al. 2018; Mayers et al. 2023).

Furthermore, a thorough analysis of this database revealed that the wind profiles issued from Ku-band scatterometer data barely exceed 64 kt, even for the most-intense TCs, as shown in appendix A. Thus, we chose to remove Ku-band scatterometers from the present analysis.

a. Radiometer missions

Because both the foam coverage and bubble surface layer thickness increase with surface wind speed (Reul and Chapron 2003), passive microwave measurements have long been known to

TABLE 1. The radiometer and scatterometer data used in [Portabella et al. \(2022\)](#). The period, spatial resolution, and pixel spacing rows refer to the wind product. The same data were used for the present work except that the Ku-band scatterometers were removed from the analysis.

Radiometer	SMOS	SMAP	AMSR-2	WindSat
Period	2010–20	2015–20	2012–20	2010–19
Spatial resolution	50 km	50 km	50 km	50 km
Pixel spacing	25 km	25 km	25 km	25 km
Frequency	L band	L band	C band; X band	C band; X band
Scatterometer	ASCAT	HSCAT	OSCAT	RSCAT
Period	2010–20 (<i>MetOp-A</i>)	2012–15 (<i>HY-2A</i>)	2010–14 (<i>Oceansat-2</i>)	2014–16
	2012–20 (<i>MetOp-B</i>)	2019/20 (<i>HY-2B</i>)	2017–20 (<i>Scatsat-1</i>)	
	2019/20 (<i>MetOp-C</i>)			
Spatial resolution	25 km	50 km	50 km	50 km
Pixel spacing	12.5 km	25 km	25 km	25 km
Frequency	C-band	Ku-band	Ku-band	Ku-band

display very high sensitivity under extreme wind conditions. With large ~ 1000 -km swaths, satellite-borne radiometers are well suited to monitor TCs. However, they have nominally low spatial resolutions (~ 40 km) that generally prevent accurate retrieval of the extreme surface wind speeds associated with the inner core of the most-intense TCs. The radiometer wind products used in this work are at 50-km spatial resolution with a 25-km grid spacing ([Portabella et al. 2022](#)).

In the present study, four different sources of radiometer data were used. Among them, the L-band (1.4 GHz; 21-cm wavelength) radiometers from the European Space Agency (ESA) Soil Moisture and Ocean Salinity (SMOS) mission and the National Aeronautics and Space Administration (NASA) Soil Moisture Active Passive (SMAP). The ability of L-band radiometers to retrieve ocean surface wind speeds under TCs has been discussed both in the case of SMOS ([Reul et al. 2012, 2016](#)) and SMAP ([Yueh et al. 2016; Meissner et al. 2017](#)). [Reul et al. \(2017\)](#) demonstrated that SMOS, SMAP, and AMSR-2 can be used to estimate wind radii.

The Japan Aerospace Exploration Agency launched the Advanced Microwave Scanning Radiometer 2 (hereinafter AMSR-2) on board the *Global Change Observation Mission—Water* satellite in 2012. This instrument is still operating today and uses seven different frequencies (6.93, 7.3, 10.65, 18.7, 23.8, 36.5, and 89.0 GHz). For TCs, the first three channels (6.93, 7.3, and 10.65 GHz) are used. With two C-band channels, initially intended for radio-frequency interference identification, surface wind estimates are improved. Signals at these two C-band frequencies have similar sensitivity to the sea wind speed but differ in sensitivity to rain by approximately 12%. Accuracy of the AMSR-2-retrieved wind speed in storms is comparable to results

obtained from SMOS and SMAP L-band sensors ([Zabolotskikh et al. 2015; Reul et al. 2017](#)).

WindSat is a polarimetric radiometer on board *Coriolis*, a mission designed by the Naval Research Laboratory and the Air Force Research Laboratory and launched in 2003. The sensor provided data until May 2021. This instrument operates at five different channels (6.8, 10.7, 18.7, 23.8 and 37.0 GHz). To minimize heavy precipitation impacts, the C-band 6.8- and the X-band 10.7-GHz channels are used for TC wind retrieval algorithms. Again, changes in the respective contribution of wind and rain to the signal measured by each channel can be used to better infer and discriminate both quantities ([Klotz and Uhlhorn 2014](#)). Heavy precipitation is still found to complicate surface wind speed retrieval with this sensor ([Quilfen et al. 2007](#)), and more recent studies addressed this issue ([Meissner et al. 2021; Manaster et al. 2021](#)).

b. Scatterometer missions

Scatterometers are active sensors that emit a pulse and measure the signal backscattered by the rough ocean surface with different viewing angles. Because backscatter signals are dependent upon both wind speed and wind direction, ocean surface wind vectors can be retrieved. The achieved nominal spatial resolution (up to ~ 25 km) is higher than satellite-borne radiometers. Actual scatterometers operate at different frequencies (C band or Ku band).

The Meteorological Operational satellite program is a series of three satellites (*Metop-A*, *-B*, and *-C*) launched by ESA (in 2006, 2012, and 2018, respectively) that include scatterometers [i.e., the Advanced Scatterometer (ASCAT)] operating at 5.3 GHz (C band). With three antennas oriented at 45° , 90° , and 135° with respect to the satellite track, the wind direction can be retrieved. ASCAT instruments have two sub-swaths, each having a width of ~ 550 km. At C band, the signal may be influenced by very heavy rain. Backscatter signals also tend to saturate at high winds ([Donnelly et al. 1999](#)), and ASCAT measurements progressively lose sensitivity under high wind speeds ([Soisuvarn et al. 2013; Polverari et al. 2022](#)). The ASCAT wind product used in the present study is at 25-km spatial resolution with a 12.5-km grid spacing ([Stoffelen et al. 2017; Portabella et al. 2022](#)).

TABLE 2. The SAR data used in the present study. The period, spatial resolution, and pixel spacing rows refer to the wind product.

SAR	<i>S1A</i>	<i>S1B</i>	<i>RS2</i>
Period	2016–21	2016–21	2012–21
Spatial resolution	3 km	3 km	3 km
Pixel spacing	1 km	1 km	1 km
Frequency	C band	C band	C band

Scatterometers operating at Ku band (~ 13.5 GHz) usually have larger swaths (~ 1000 km) than C-band scatterometers but suffer more contamination in heavy rainfall conditions [see [Quillen et al. \(2007\)](#) for more details]. The Ku-band scatterometer wind products used in [Portabella et al. \(2022\)](#) were finally removed (see [appendix A](#)). They include the China National Space Administration (CNSA) Haiyang missions (hereinafter HSCAT), the Indian Space Research Organization (ISRO) *OceanSat-2* and *ScatSat-1* satellites (hereinafter OSCAT), and the NASA RapidScat (hereinafter RSCAT) on board the International Space Station ([Table 1](#)).

c. SAR missions

The SAR data used here come from three different missions: ESA *Sentinel-1A* and *Sentinel-1B* (hereinafter *S1A* and *S1B*, respectively), and the Canadian Space Agency (CSA) *RADARSAT-2* (hereinafter *RS2*). The SAR instrument on board these three missions is an active sensor operating at 5.4 GHz (C band). By analyzing the received signal in both co- and cross-polarization, wind speeds can be inferred under TC conditions including at very high wind speeds ([Mouche et al. 2017, 2019](#)). Convincing comparisons with passive radiometers have been performed ([Zhao et al. 2018](#)). The ability of SAR-derived wind speeds to accurately capture the TC ocean surface wind structure, including R_{\max} , has further been demonstrated and discussed by [Combot et al. \(2020a\)](#).

Today, SAR wide-swath acquisitions cannot be continuously performed over oceans. Based on track forecasts, it is still possible to best anticipate when the sensor will overpass a TC, and to trigger a SAR acquisition. ESA started the Satellite Hurricane Observation Campaign (SHOC) in 2016, resulting in more than ~ 500 acquisitions over TCs. The derived wind products ([Mouche et al. 2017](#)) are further interpolated on a regular polar grid based on the TC center [see the appendix in [Vinour et al. \(2021\)](#)]. The product has a 3-km spatial resolution, with a 1-km grid spacing. This spatial resolution approximates a 1-min wind speed as a 50 m s^{-1} wind moves 3 km in a minute. In this study, a certain number of SAR cases have been discarded on a qualitative basis, for example when the detected TC center was judged to be wrong or when the SAR file contained corrupted pixel values.

d. Best tracks

Here, IBTrACS data were used for several purposes: the storm centers (latitude and longitude) allowed to azimuthally average the radiometer and scatterometer wind fields, while the wind radii (R_{34} , R_{50} , and R_{64}) were compared with satellite-based wind radii. Both IBTrACS latitude (to compute f) and maximum sustained wind speed (V_{\max}) were used in the [CK22](#) framework, and the distance to closest land (from the TC center) enabled filtering of the dataset. These parameters were extracted for the period 2010–20.

In IBTrACS, some storm tracks are given on a 6-hourly basis, while others are interpolated and thus given on a 3-hourly basis. To account for this varying sampling time, all tracks and their associated parameters were interpolated on an hourly basis with a monotonic cubic interpolation. Last, because of

varying definitions of the maximum sustained wind speed across the different agencies, we selected only U.S. agencies (i.e., National Hurricane Center, Joint Typhoon Warning Center, and Central Pacific Hurricane Center), which all provide the 1-min maximum sustained wind speed.

e. Data filtering

To further restrain the analysis to well-formed systems, i.e., for which R_{\max} can be well determined from the axisymmetric mean profile, and to best ensure consistency with [CK22](#) for further comparison, the following filters have been applied to our dataset:

- 1) $V_{\max} > 20 \text{ m s}^{-1}$,
- 2) $R_{\max} < 150 \text{ km}$,
- 3) any wind radius must be $> 5 \text{ km}$,
- 4) absolute latitude $< 30^\circ$, and
- 5) distance to closest land $> R_{34}$.

Here and below, V_{\max} , R_{\max} , and “wind radii” refer to their values when estimated on azimuthally averaged wind profiles (see below). Unlike [CK22](#), we did not apply any filter on longitude. Therefore, the method presented here applies in every basin and does not depend on the availability of aircraft analysis.

3. Methods and data analysis

a. Estimation of the [CK22](#) predictors

To apply the [CK22](#) framework to the intercalibrated dataset of radiometer and scatterometer data, estimates of the predictors (V_{\max} , R_{34} , f) were needed for every satellite file.

For the wind radii, an azimuthally averaged wind profile was first computed for every satellite file using the corresponding IBTrACS center linearly interpolated to the acquisition time. For each of the three speed values of interest (i.e., 34, 50, and 64 kt), we then selected the radius where the outer profile matches this value to the closest kilometer. Should there be more than one radius value, the wind radius was defined as the smallest of the radii. During the process, wind radii estimates are affected by IBTrACS linearly interpolated center uncertainties. By comparing SAR-based center estimates (see [Vinour et al. 2021](#)) with IBTrACS-based center estimates over the whole SAR database, the average uncertainty is $\sim 13 \text{ km}$, largely below the radiometer and scatterometer effective spatial resolutions.

Unlike the wind radii, V_{\max} and f cannot be accurately estimated from radiometer and scatterometer data, especially for intense small TCs, but both parameters are systematically reanalyzed in the best tracks. However, IBTrACS V_{\max} definition does not strictly coincide with the axisymmetric view adopted here. In particular, the analysis ([appendix A](#)) highlighted that V_{\max} estimated using SAR azimuthally averaged profiles were, on average, lower than IBTrACS V_{\max} . This can be modeled by applying a linear regression (dashed gray line in [appendix Fig. A1](#)) to IBTrACS V_{\max} estimates. The resulting intensity estimates are denoted by V_{\max}^{REG} and were used (instead of the raw IBTrACS V_{\max}) to ensure the consistency with the wind radii defined on azimuthally averaged wind profiles. The pair of parameters

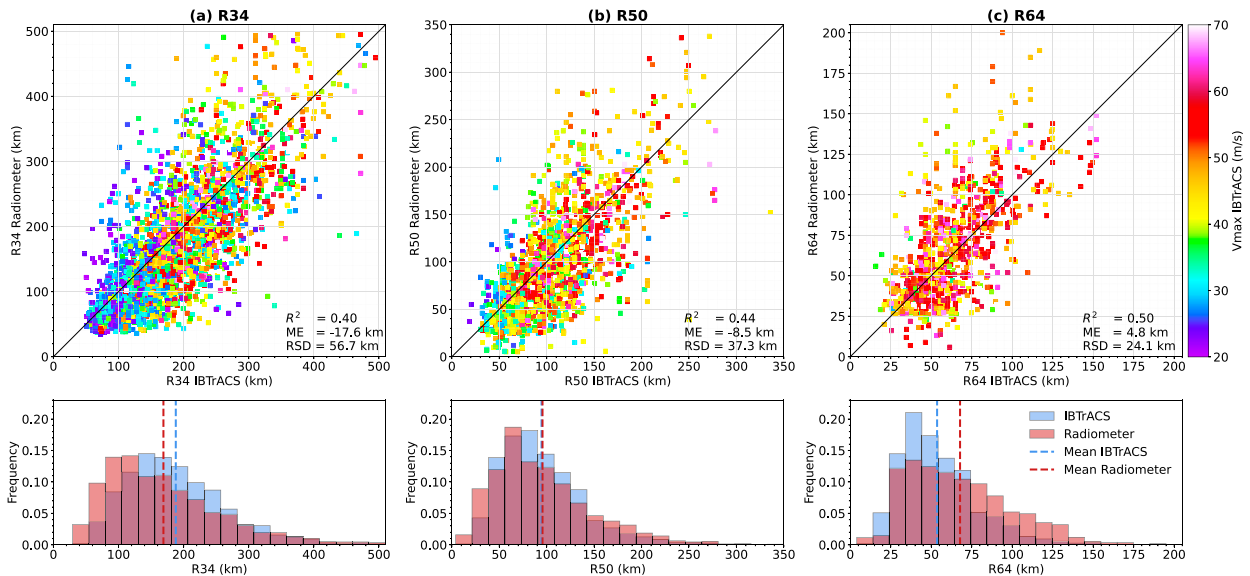


FIG. 2. (top) Comparison between radiometer (y axis) and corresponding IBTrACS (x axis) wind radii [coefficient of determination R^2 , mean error (ME), and residual standard deviation (RSD) are displayed at the bottom left of each panel] and (bottom) corresponding distributions and averages, shown for (a) R_{34} , (b) R_{50} , and (c) R_{64} .

(V_{\max}^{REG}, f) was then linearly interpolated to the satellite acquisition time for every file.

b. Quality assessment of radiometer and scatterometer wind radii estimates

To assess the quality of the satellite-based wind radii, comparisons were performed with IBTrACS wind radii. A strict comparison cannot be achieved because of varying definitions. In IBTrACS, wind radii are relative to the geographical quadrants and correspond to the maximum radial extent of the associated wind speed in each of the four quadrants. To make IBTrACS values as close as possible to the satellite-based wind radii, the nonzero IBTrACS values were averaged over all the quadrants. Furthermore, both the methodologies and the available observational data can vary across the IBTrACS dataset. Here, the adopted strategy was to compare the whole IBTrACS wind radii dataset (including non-U.S. agencies for this section) to the satellite-based wind radii. Accounting for the differences between the specialists and agencies is beyond the scope of this study. Finally, after removing the Ku-band sensors (see appendix A), we separated radiometer wind radii from the C-band scatterometer wind radii to further investigate possible discrepancies between the remaining sensors.

Figure 2 shows a comparison between radiometer wind radii and IBTrACS values (top) and their corresponding distributions (bottom). While radiometer wind radii look well correlated with IBTrACS values, with R^2 -scores ranging from 0.4 to 0.5, large discrepancies arise, with a residual standard deviation (RSD) as high as 56.7 km for R_{34} . The RSD decreases to 37.3 km for R_{50} , and further to 24.1 km for R_{64} , reflecting the decrease of the mean wind radius, that is, from 181 km for R_{34} to 51 km for R_{64} in IBTrACS. In terms of

relative uncertainties, this leads to $\sim 31\%$, $\sim 36\%$, and $\sim 41\%$ for R_{34} , R_{50} , and R_{64} , respectively. Interestingly, the mean error (ME) is negative for both R_{34} and R_{50} , showing that, on average, these wind radii are lower when extracted from azimuthally averaged radiometer profiles than from IBTrACS. This is likely the result of the differing definition of the wind radii in the satellite data and in IBTrACS. Indeed, on average, a wind radius extracted from an azimuthally averaged profile is expected to be smaller than the maximum radial extent of the same wind speed. Biases due to the differing definition are lower for R_{50} and R_{64} than for R_{34} , because these radii are smaller on average. This definition effect is illustrated on the distribution for R_{34} , where the radiometer R_{34} distribution is biased toward lower values relative to IBTrACS.

Figure 3 shows comparisons between C-band scatterometer wind radii and IBTrACS values. Again, an overall consistency emerges between both data sources for all wind radii. RSD values and R^2 scores are comparable to the previous comparisons between radiometer and IBTrACS. Data and methodology are thus consistent with IBTrACS (which is expected since radiometer and scatterometer data are often used during the reanalysis process), but it also shows that there is a good consistency between the various sensors in terms of wind radii.

For R_{64} , the ME is slightly positive for both radiometer and scatterometer data (Figs. 2c and 3c), with a distribution of R_{64} skewed toward higher values for the satellite data relative to IBTrACS. First, this could be attributed to satellite data limitations, such as low spatial resolution, signal saturation, or rain contamination. Yet, Fig. 4 offers a different explanation. It again shows comparisons between scatterometer wind radii and IBTrACS values, but only over the 3-yr period from 2018

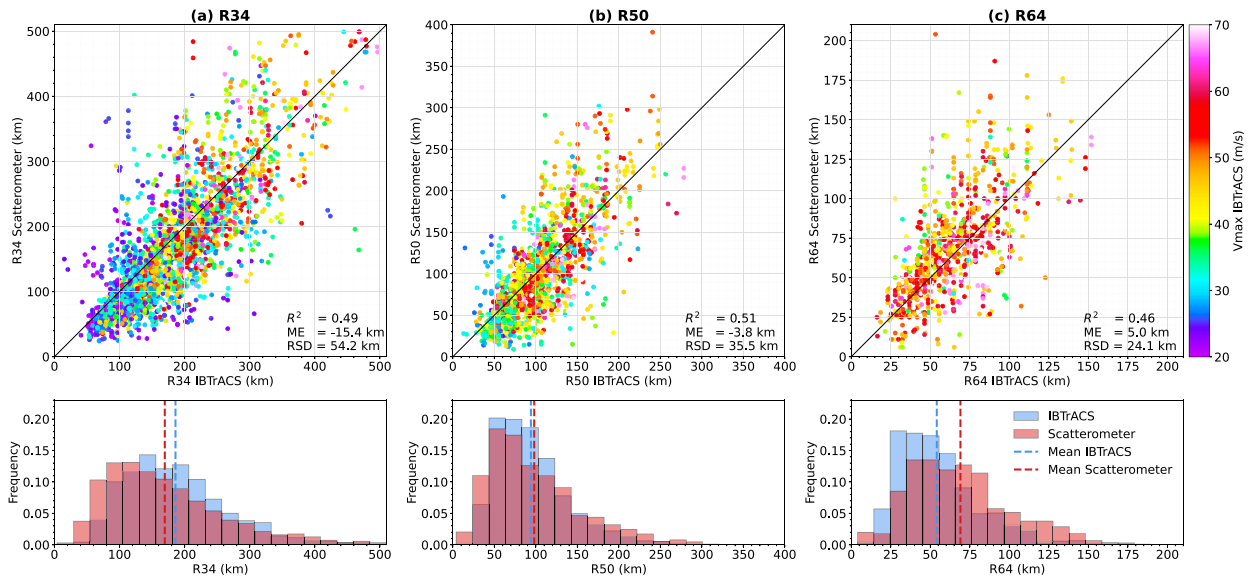


FIG. 3. As in Fig. 2, but for the C-band scatterometer wind radii.

to 2020. For such a period, the computed ME for R_{64} is only 1.5 km (Fig. 4c) and the RSD drops to 19.4 km (as compared with 24.1 km for 2010–20). Consistency between scatterometer and IBTrACS also improves for both R_{34} and R_{50} over the same period (Figs. 4a,b). The positive ME for R_{64} in Fig. 3 likely corresponds to the improving quality of IBTrACS over the years. As mentioned in section 1, wind-radii best-track values were not necessarily reanalyzed depending on the year and the basin. Similar conclusions were obtained with radiometer data (not shown).

To summarize, the comparison between IBTrACS and the intercalibrated dataset shows that radiometers and

scatterometers provide reliable wind radii estimates. Thus, for every radiometer or scatterometer acquisition, we can extract a corresponding set of predictors constituted by a satellite-based wind radius along with IBTrACS V_{\max}^{REG} and f estimates.

c. Collocations of radiometers and scatterometers with SAR

To fit the CK22 model, we also needed an estimate of the predictand R_{\max} for each set of predictors. The latter cannot be directly evaluated from radiometer and scatterometer data but is well observed on SAR data by taking the

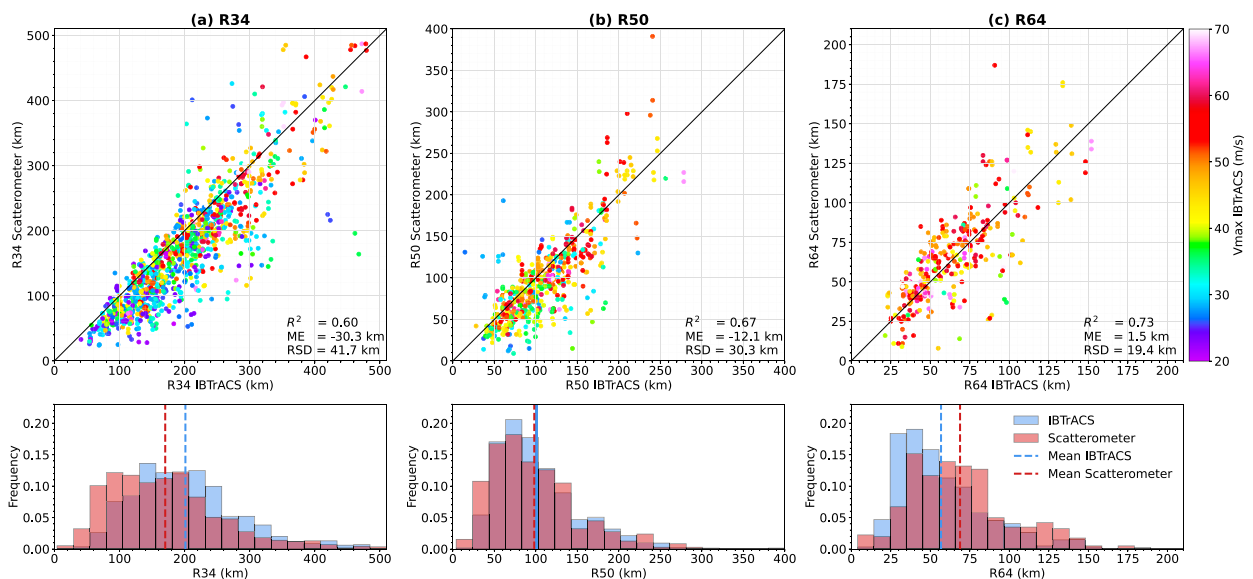


FIG. 4. As in Fig. 3, but only for the 3-yr period 2018–20.

TABLE 3. Number of collocations between SAR and the intercalibrated dataset (radiometer and ASCAT), and corresponding average absolute time difference.

	SMOS	SMAP	AMSR-2	WindSat	ASCAT	TOTAL
Before filtering	106	63	0	100	0	269
After filtering	67	33	0	45	0	145
Avg Δt (min)	12	21		31		19

location of the wind profile maximum. Thus, we looked for collocations between SAR and radiometer or scatterometer TC overpasses. Two images were considered to be collocated if their absolute acquisition time difference is less than 90 min.

For radiometer data (Table 3, first four columns), this procedure resulted in a total of 269 collocations, which further reduced to 145 collocations after applying filters presented in section 2e. Notably, no collocation was found between any of the 3 SAR missions (*SIA*, *SIB*, and *RS2*) and AMSR-2. The average absolute time difference of the found collocations is ~ 19 min, during which we assume the TC wind structure to remain stationary.

For scatterometer data, no collocation was found between SAR and ASCAT (Table 3, penultimate column). In what follows, we thus refer to the dataset obtained by this collocation procedure as the ‘‘SAR–radiometer collocation dataset.’’ It consists of predictors estimated on radiometer data (wind radii) or corresponding IBTrACS values (V_{\max}^{REG} and f), and predictands estimated on SAR (R_{\max}). Note that we could equally have used SAR wind radii estimates to fit the CK22 model, but this approach would have reduced the number of available wind radii estimates because SAR instruments have smaller swaths than do radiometers.

4. Results

a. Fitting CK22 model

As explained in the introduction, the CK22 model relies on the estimation of the ratio M_{\max}/M_{34} via a log-linear regression model, using $[X_{34}^{(1)}, X_{34}^{(2)}]$ as input. Although CK22 used R_{34} in their study, this method is in fact agnostic to the choice of wind radius. Therefore, the ratio M_{\max}/M_{50} can also be estimated using $X_{50}^{(1)} := V_{\max} - 25.7 \text{ m s}^{-1}$ and $X_{50}^{(2)} := (V_{\max} - 25.7 \text{ m s}^{-1})[(1/2)/fR_{50}]$ as input (or M_{\max}/M_{64} using $X_{64}^{(1)} := V_{\max} - 32.9 \text{ m s}^{-1}$ and $X_{64}^{(2)} := (V_{\max} - 32.9 \text{ m s}^{-1})[(1/2)/fR_{64}]$ as input).

CK22 estimated the coefficients of the log-linear regression model based solely on IBTrACS rather than direct observational estimates, and only for the ratio M_{\max}/M_{34} . In the present work, we use observational data not only to obtain improved estimates of the predictors in the CK22 model framework, but also to obtain improved estimates of the model coefficients that relate the parameters to one another. We also extend the CK22 model for the ratios M_{\max}/M_{50} and M_{\max}/M_{64} . A log-linear regression model was fitted for each of the three ratios using the SAR–radiometer collocation dataset previously presented. The following relationships were obtained:

$$M_{\max}/M_{34} = 0.531 \exp\{-0.00214(V_{\max}^{\text{REG}} - 17.5 \text{ m s}^{-1}) - 0.00314(V_{\max}^{\text{REG}} - 17.5 \text{ m s}^{-1})[(1/2)/fR_{34}]\}, \quad (3)$$

$$M_{\max}/M_{50} = 0.626 \exp\{0.00282(V_{\max}^{\text{REG}} - 25.7 \text{ m s}^{-1}) - 0.00724(V_{\max}^{\text{REG}} - 25.7 \text{ m s}^{-1})[(1/2)/fR_{50}]\}, \quad \text{and} \quad (4)$$

$$M_{\max}/M_{64} = 0.612 \exp\{0.00946(V_{\max}^{\text{REG}} - 32.9 \text{ m s}^{-1}) - 0.01183(V_{\max}^{\text{REG}} - 32.9 \text{ m s}^{-1})[(1/2)/fR_{64}]\}. \quad (5)$$

With these formulas, R_{\max} can then be estimated using the steps presented in the introduction [Eq. (2)]. Subsequent estimates will be referred to as $R_{\max}^{\text{CK22-}R_{34}}$, $R_{\max}^{\text{CK22-}R_{50}}$, or $R_{\max}^{\text{CK22-}R_{64}}$ depending on which wind radius is used.

b. Assessment of the resulting R_{\max} estimates

To check the fitting procedure, we compared $R_{\max}^{\text{CK22-}R_{34}}$ estimates and SAR R_{\max} references (Fig. 5a). The consistency between both is reasonably good, with a R^2 -score of 0.41 and an RSD of 10.6 km. A low ME of 3.7 km is observed, which can be related to the distribution of $R_{\max}^{\text{CK22-}R_{34}}$ being slightly skewed toward higher R_{\max} values relative to SAR.

Because R_{50} and R_{64} are closer to R_{\max} than R_{34} , using one or the other wind radii thresholds should improve the quality of the ratio estimate relative to R_{34} . Ideally, an estimate of R_{\max} should be performed with R_{64} if available. If R_{64} is not defined (i.e., if V_{\max} is less than 33 m s^{-1}), R_{50} should be used. R_{34} should only be used if both R_{64} and R_{50} were not defined. Following this procedure, we further estimated R_{\max} using the ‘‘best’’ available wind radius.

Figure 5b shows a comparison between these estimates (hereinafter $R_{\max}^{\text{CK22-BR}}$) and SAR R_{\max} references. The R^2 -score increased to 0.63 and the ME decreased to 0.9 km in comparison with the $R_{\max}^{\text{CK22-}R_{34}}$ methodology, and RSD decreased from 10.6 to 8.8 km. Therefore, use of wind radii closer to R_{\max} does improve the estimate quality. In addition, such a low RSD demonstrates the efficiency of the fitted CK22 relationships [Eqs. (3)–(5)] to provide reliable R_{\max} estimates.

In their paper, the R_{\max} predicted by CK22 had a systematic bias that could be bias adjusted in postprocessing to improve the model. Here we find that our model does not require a bias adjustment, which may be an indication of the benefit of using direct observational data for R_{\max} (SAR).

While the method is successful on average, it is remarkable that errors can be large (more than ~ 10 km), even for cases where R_{64} predictors are used (see for instance Kong-Rey and Mangkhut in Fig. 5b). Before discussing how to explain these large uncertainties, a single TC life cycle was chosen to illustrate the potential of the present methodology.

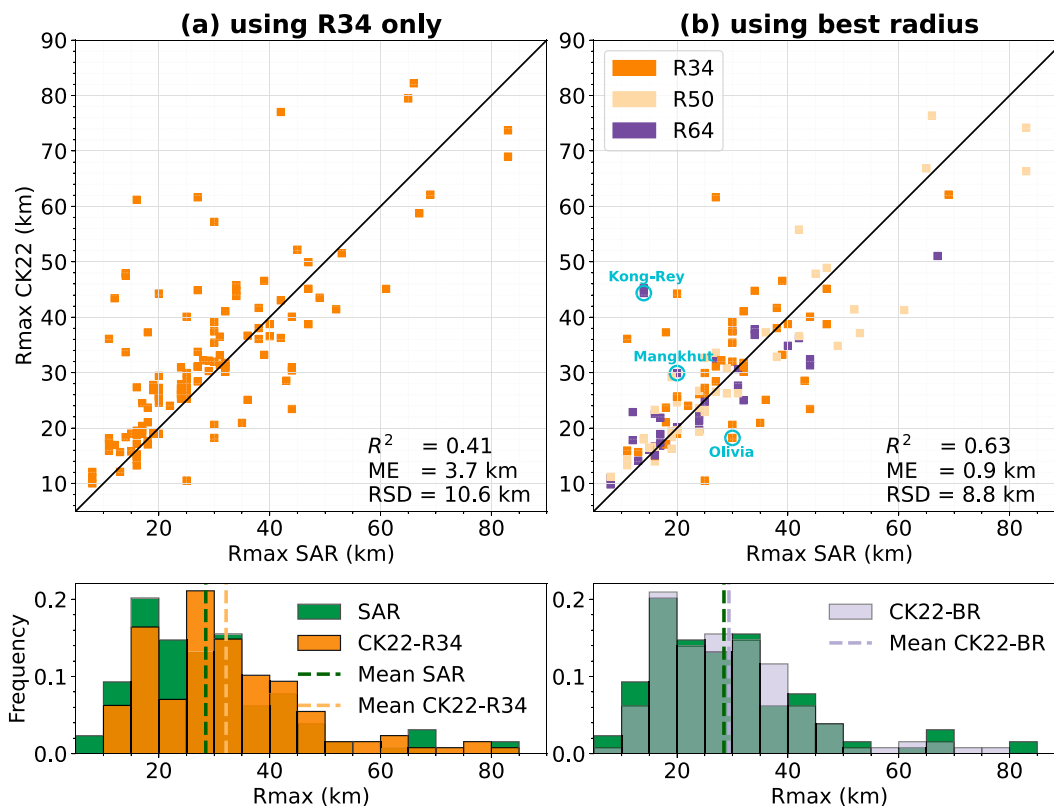


FIG. 5. (top) Comparison between R_{\max} estimates using the CK22 model and SAR R_{\max} and (bottom) corresponding distributions for (a) $R_{\max}^{\text{CK22-R34}}$ and (b) $R_{\max}^{\text{CK22-BR}}$. For analysis purposes, color reveals which radius was used to define $R_{\max}^{\text{CK22-BR}}$ for each case.

c. Application to TC Kilo life cycle

Producing $R_{\max}^{\text{CK22-BR}}$ estimates every time a radiometer or a scatterometer TC overpass is available can be an efficient tool for characterizing the time evolution of R_{\max} for any given TC. Figure 6 shows TC Kilo R_{\max} and V_{\max}^{REG} time series between 27 August and 10 September 2015, a period over which V_{\max}^{REG} was larger than 20 m s^{-1} . TC Kilo evolved in the Pacific Ocean, reaching category 4 on the Saffir–Simpson scale. It intensified from 20 to 49 m s^{-1} between 27 August and 30 August before entering a weakening phase. In the meantime, R_{\max} first varied between 55 and 15 km according to IBTrACS, then stagnated at 37 km between 30 August and 2 September, before varying again after these dates. Stagnation phases of R_{\max} from IBTrACS are likely not physical according to the V_{\max}^{REG} variations during that time interval (see section 5) and the two eyewall replacement cycles (ERCs) suggested by passive microwave observations (not shown). The $R_{\max}^{\text{CK22-BR}}$ estimates show much more pronounced variations during those phases, with an increasing trend between 30 August and 8 September. This particular phase corresponds to an overall decrease of V_{\max}^{REG} and an overall increase of R_{64} in our data (not shown), both of which would be expected to be associated with an increase in R_{\max} .

For reference, three SAR R_{\max} estimates were available during TC Kilo’s life cycle (green stars). The first SAR R_{\max}

(10 km) on 27 August does not match with our first estimate of R_{\max} (35 km). This illustrates the limitations of our proposed methodology and is discussed herein. The second and third SAR R_{\max} estimates are in better agreement with the $R_{\max}^{\text{CK22-BR}}$ estimates, especially if we account for the overall R_{\max} trend given by our estimates. A large eye is also depicted in passive microwave data during this period (not shown), supporting the robustness of the $R_{\max}^{\text{CK22-BR}}$ estimates.

Notably, there is more spread in the CK22 estimates on the last two days of the study period. Despite this increasing uncertainty, the increase of R_{\max} is well depicted, suggesting R_{\max} would significantly increase before 8 September in contradiction with the IBTrACS trend.

In summary, every time a radiometer or scatterometer wind profile is available, a subsequent $R_{\max}^{\text{CK22-BR}}$ estimate can be obtained, using the proposed objective method. In such a way, one can estimate R_{\max} trends that are more realistic than IBTrACS, less impacted from spatial or temporal heterogeneities. Such a framework could also be used operationally.

5. Discussion

The previous section demonstrated the potential of the CK22 model fitted with SAR, when used in combination with intercalibrated medium-resolution radiometer and scatterometer data. Still, $R_{\max}^{\text{CK22-BR}}$ estimates can display rather large

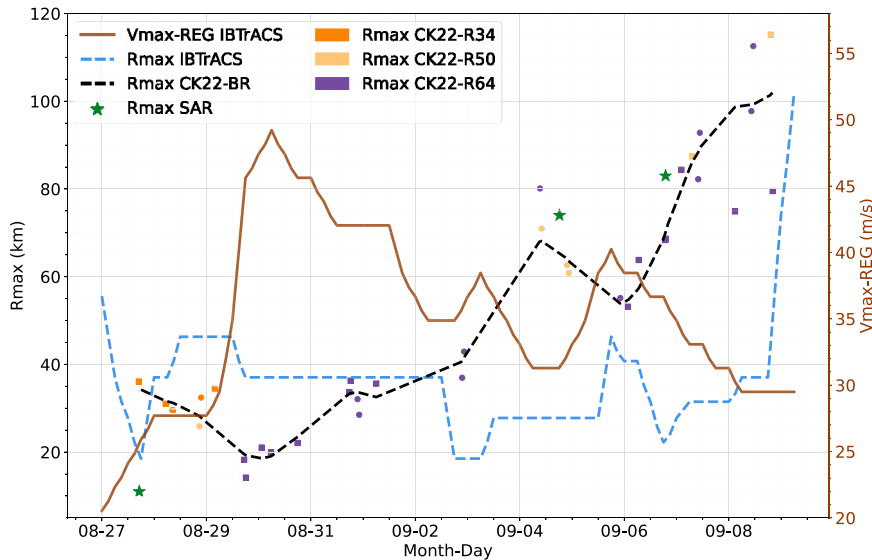


FIG. 6. TC Kilo (2015) time series of IBTrACS R_{\max} (left axis; dashed blue), radiometer- and scatterometer-based $R_{\max}^{\text{CK22-BR}}$ (left axis; dashed black), along with IBTrACS-based V_{\max}^{REG} (right axis; solid brown). Also displayed are radiometer (squares) and scatterometer (circles) $R_{\max}^{\text{CK22-BR}}$ estimates (color reveals which radius was used to define $R_{\max}^{\text{CK22-BR}}$ for each observation), and SAR R_{\max} estimates (green stars). The dashed black line was used to define $R_{\max}^{\text{CK22-BR}}$ estimates.

uncertainties, despite the expected improved use of R_{64} as predictor. To better understand the sources of such uncertainties, three other case studies (cyan circles on Fig. 5a) were considered before examining theoretical aspects and drawing a picture of the average situation.

a. Case studies from the SAR-radiometer dataset

The first case (Fig. 7, left column) is TC Olivia in 2018, an eastern Pacific Ocean hurricane that reached category 4 on the Saffir–Simpson scale. It reached a first intensity peak ($\sim 56 \text{ m s}^{-1}$) on 5 September, then weakened before re-strengthening ($\sim 59 \text{ m s}^{-1}$) during the night between 6 and 7 September. On 8 September, both *RS2* at 1510 UTC and *WindSat* at 1533 UTC overflew Olivia (Figs. 7a,d). Its eyewall, depicted by the high-resolution SAR observation, was clearly defined though asymmetric. With its rather low spatial resolution, the radiometer failed to map the inner core areas with high wind speed gradients, and eyewall asymmetries. These differences between SAR and radiometer two-dimensional observations translate into differences in the azimuthally averaged wind profiles. From the SAR wind profile, Olivia's R_{\max} was 30 km at that time, with a V_{\max} of 32 m s^{-1} (Fig. 7g). Notably, *WindSat* failed to estimate V_{\max} correctly, with a negative bias of almost 10 m s^{-1} when compared with SAR V_{\max} and V_{\max}^{REG} , which are in good agreement at that time. This bias is largely attributable to sensor spatial averaging effects. In fact, the entire azimuthally averaged wind profile is negatively biased, leading to an underestimation of R_{34} , further reflected in $R_{\max}^{\text{CK22-BR}}$. This case illustrates how wind radii uncertainties translate into $R_{\max}^{\text{CK22-BR}}$ uncertainties. Note that in other cases uncertainties on V_{\max}^{REG} could also affect $R_{\max}^{\text{CK22-BR}}$ uncertainties.

The second case (Fig. 7, middle column), Mangkhut, was a super typhoon (category 5 on Saffir–Simpson scale), causing considerable damage in the western Pacific region in 2018. It reached its peak intensity ($\sim 80 \text{ m s}^{-1}$) on 12 September. On 11 September, both *SIB* at 2048 UTC and *WindSat* at 2126 UTC overflew Mangkhut (Figs. 7b,e). According to the SAR observation, Mangkhut had a clearly defined symmetric eyewall at that time. Note that the eyebrow shape in the high winds to the left of the eyewall (Fig. 7b) is probably due to rain contamination [for discussion about such a feature see Mouche et al. (2019)]. The extent of high winds was seemingly well captured by the radiometer sensor, but the eye was not resolved. Nevertheless, a very good agreement between *SIB* and *WindSat* wind outer-profiles is obtained for this case (Fig. 7h), with only ~ 3 -km difference between R_{64} estimates from the two sensors. Still, the estimate given by $R_{\max}^{\text{CK22-BR}}$ (~ 30 km) largely overestimates the actual SAR R_{\max} (~ 20 km). Note that in this case the clear eye depicted by infrared data and the ring captured by passive microwave sensors are both small (not shown), supporting the small SAR R_{\max} estimate. With its large R_{64} and small R_{\max} at that time, Mangkhut illustrates the high variability that occurs in nature. Such a case is likely to depart from any statistical relationship (like CK22) that links a wind radius to R_{\max} .

The last case study (Fig. 7, right column), Kong-Rey, in 2018, was a super typhoon reaching category 5 on the Saffir–Simpson scale, also evolving over the western North Pacific Ocean. Following an $\sim 72 \text{ m s}^{-1}$ peak intensity on 2 October Kong-Rey experienced an ERC and entered its weakening phase. Kong-Rey was captured on 2 October by both *SIA* at 2111 UTC and *SMAP* at 2133 UTC (Figs. 7c,f). The SAR observation depicts a well-defined symmetric eyewall, with a secondary ring of

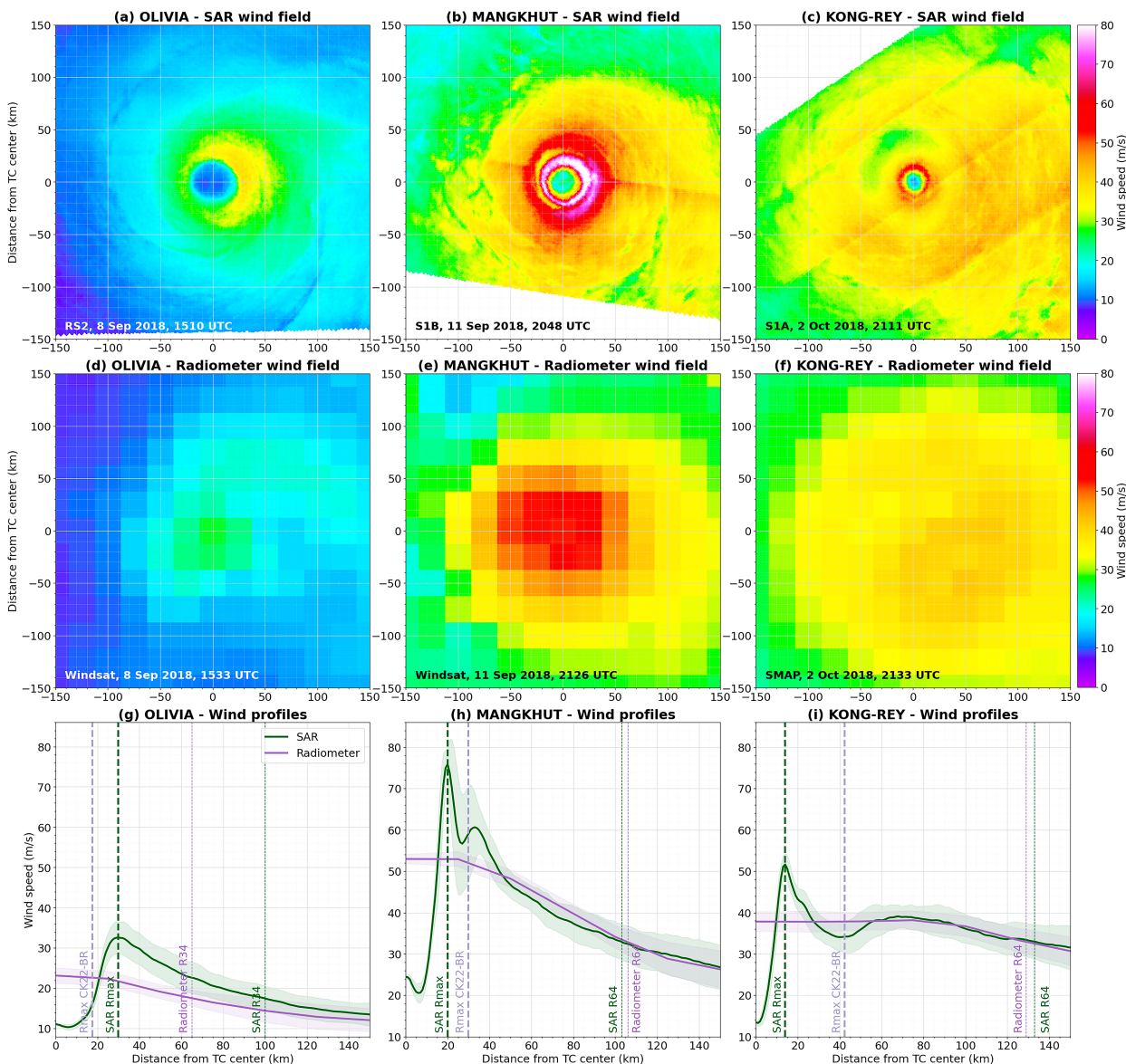


FIG. 7. Comparison of (a)–(c) SAR and (d)–(f) radiometer wind fields (TCs are translating toward the top of each panel), along with (g)–(i) corresponding wind profiles, for (left) Olivia, (center) Mangkhut, and (right) Kong-Rey.

maximum winds farther out from the TC center. In fact, Kong-Rey exhibited two eyewalls in 89-GHz imagery at this time (not shown). These two high wind regions were not well captured by the radiometer. The radiometer wind profile saturates in the 80-km inner-part of the TC, while the SAR wind profile exhibits two wind speed local maxima (Fig. 7i). Despite the inability of the radiometer sensor to capture the dual wind maxima observed at this time, the outer parts of the azimuthally averaged wind profiles match well, both yielding an R_{64} estimate of ~ 128 km. However, $R_{\max}^{\text{CK22-BR}}$ is 42 km, far from the 14 km of SAR R_{\max} . Note, however, that $R_{\max}^{\text{CK22-BR}}$ lies between the two SAR wind maxima. The complex shape of Kong-Rey during its ERC is the main cause to explain such a huge discrepancy. Indeed,

the R_{64} estimate is pushed to an outer radius due to the existence of secondary wind maxima.

b. Structural aspects

From these examples, we see that neither the use of high-quality data (SAR) to train the algorithm nor the use of a radius that is very close to R_{\max} (i.e., R_{64}) precludes large uncertainties of R_{\max} estimates using the CK22 framework. Underlying CK22, the use of an outer wind radius (e.g., R_{34} ¹) to estimate R_{\max} is justified by the angular momentum conservation

¹ In this section we chose R_{34} as outer wind radius for clarity, but the reasoning well applies to any other wind radius (e.g. R_{50} and R_{64}).

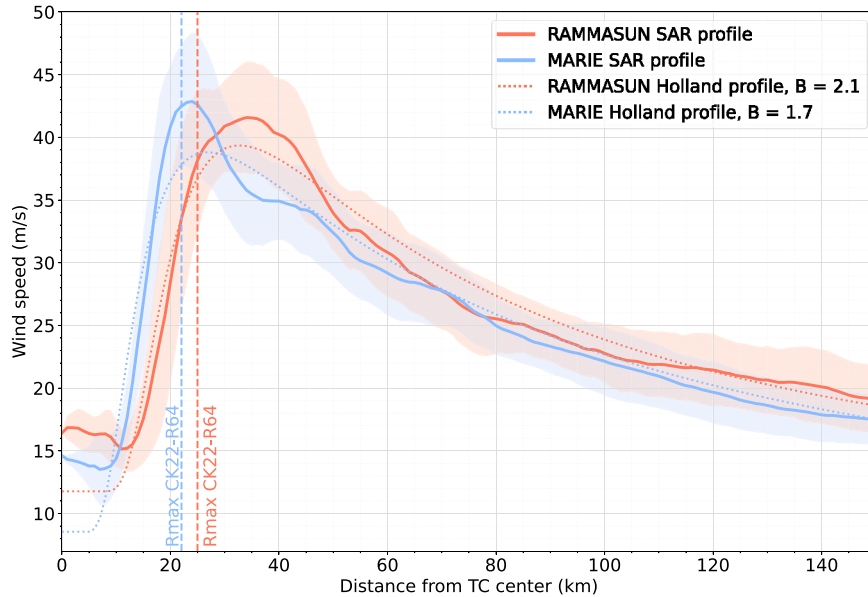


FIG. 8. SAR wind profiles for Rammasun (solid red) and Marie (solid blue) and associated Holland best-fit profiles (dotted curves) fitted on $0 \leq r \leq 500$ km.

principle: an air parcel, advected from the outer radii to the innermost radii, must lose angular momentum due to surface friction. The ratio M_{\max}/M_{34} thus represents the ability for an air parcel to keep its angular momentum while being advected from R_{34} to R_{\max} . In the log-linear framework, this ratio solely depends on V_{\max} , R_{34} , and f .

The use of these three parameters to estimate M_{\max}/M_{34} was discussed in Chavas et al. (2015) and Chavas and Lin (2016). In these studies, the ability of a radial parametric wind profile to represent the variability of observational data was tested. In brief, the radial parametric profile geometrically merges an inner-part profile with an outer-part profile, previously anticipated from theoretical studies (Emanuel and Rotunno 2011; Emanuel 2004). Chavas and Lin (2016) concluded that the ratio M_{\max}/M_0 between the angular momentum at R_{\max} and at an outer radius R_0 solely depends on four parameters: V_{\max} , fR_0 , C_k/C_d , and W_{cool}/C_d , where C_k and C_d are the heat and momentum exchange coefficients and W_{cool} models the radiative-subsidence rate in the free troposphere of the outer-part model. Considering $R_0 = R_{34}$, a log-linear dependence of M_{\max}/M_{34} on (V_{\max}, R_{34}, f) thus neglects the variations of both C_k/C_d and W_{cool}/C_d .

Besides, the axisymmetric and steady-state theory of Emanuel and Rotunno (2011) invokes a direct relationship between M_{\max}/M_{34} and C_k/C_d that can be stated as

$$\frac{M_{\max}}{M_{34}} = \pi \left(\frac{C_k}{C_d} \right), \quad (6)$$

with $\pi(x) := [(1/2)x]^{1/(2-x)}$ being a monotonically increasing function [see their Eq. (38)]. This relationship assumes the TC is in steady-state and the Richardson number in the outflow is slightly below 1. The latter implies the outflow is self-stratified

by small-scale turbulence. Using numerical simulations that resolved convection, Emanuel and Rotunno (2011) showed that such an assumption was satisfied in an outflow region near R_{\max} . This assumption might then not hold true farther out. Chavas et al. (2015) suggested that the optimal merging radius between the inner and outer part of the model was $\sim 2-3R_{\max}$ when fitting the complete parametric profile to observational data. While not strictly corresponding to the region where the theoretical developments of Emanuel and Rotunno (2011) could remain valid, it identifies the region where the inner part of the model is most likely to apply to the observations.

When writing Eq. (6), one assumes that the model of Emanuel and Rotunno (2011) is still valid at R_{34} , which largely exceeds $3R_{\max}$ in nature. This might be a strong approximation, but it offers an instructive relationship between the rate of conservation of angular momentum (left-hand side) to a function of C_k/C_d , characterizing the balance between energy generation and friction loss (right-hand side). Most important, C_k/C_d controls the shape of the parametric radial wind profile, with higher values corresponding to more peaked profiles. In practice, unlike C_k/C_d values, this shape of the near-peak radial wind profile is more easily quantifiable using SAR data.

To highlight these considerations, we present TC cases that have the same CK22 predictors (V_{\max} , R_{64} , and f) but different wind profile shapes near their peak intensities. Figure 8 is representative of such a situation. SAR acquisitions over TC Rammasun (western Pacific; red curve) and TC Marie (eastern Pacific; blue curve), occurred at 1027 UTC 17 July 2014 and at 1419 UTC 3 October 2020, respectively. Both storms display similar outer-core profiles, with almost the same R_{64} (~ 52 and ~ 49 km), V_{\max} (~ 42 and ~ 43 m s $^{-1}$) and f (~ 4.3 and ~ 4.6 s $^{-1}$). Applying CK22 to these cases (vertical dashed

lines) thus leads to almost the same $R_{\max}^{\text{CK22-BR}}$ value (~ 25 and ~ 22 km). However, SAR-derived wind profiles provide different estimates, R_{\max} (~ 34 and ~ 24 km, respectively).

Comprehensively, the CK22 model cannot fully adjust to peculiar local wind profiles. To quantify the wind profile shapes, a Holland parametric profile (Holland 1980) was adjusted to each SAR azimuthally averaged wind profile:

$$V_{\text{Holland}}(r) = V_{\min} + \sqrt{(V_{\max} - V_{\min})^2 \left(\frac{R_{\max}}{r}\right)^B e^{1-(R_{\max}/r)^B} + \left(\frac{rf}{2}\right)^2} - \frac{rf}{2}. \quad (7)$$

This parametric formulation is useful to quantify variations in the shape of observed wind profiles. In particular, the empirical B parameter controls the rate of radial decay of the tangential winds, with higher and smaller values respectively corresponding to narrower and broader vortices. In addition, this parameter was found to be sensitive to TC intensity and size while independent of R_{\max} (Knaff et al. 2011).

Note that Holland’s profiles were designed for gradient-level wind and are not necessarily suited for surface wind profiles with nonzero wind speeds at the TC center, well captured using SAR observations. A complementary degree of freedom V_{\min} was thus included in Eq. (7) to cope with the existence of nonzero minimum wind speeds.

Using the full extent of the wind profile, a solution for V_{\min} , V_{\max} , R_{\max} , and B can be estimated via least squares. Applied to TCs Rammasun and Maria, the fitting procedure results in two different B values, ~ 2.1 and ~ 1.7 , respectively (Fig. 8). Such a difference quantifies the remaining variability of the near-core wind profile for comparable outer-core wind profiles.

c. Analysis framework

The shape of the near-core wind profile is generally associated with the radial gradient of absolute angular momentum and thus the loss of angular momentum when an air parcel is advected from R_{34} to R_{\max} . To guide the analysis, we recall the equation of angular momentum conservation for an axisymmetric vortex:

$$\frac{\partial M}{\partial t} + u \frac{\partial M}{\partial r} + w \frac{\partial M}{\partial z} = \frac{r \partial \tau_{\theta z}}{\rho}, \quad (8)$$

with u and w as the radial and vertical velocities, respectively; $\tau_{\theta z}$ is a tangential stress component; and ρ is the density. The continuity equation links u and w as

$$\frac{1}{r} \frac{\partial(ru)}{\partial r} + \frac{\partial w}{\partial z} = 0. \quad (9)$$

Under steady state condition, Eq. (8) can be integrated from the surface to a boundary layer height, h , where the stress vanishes:

$$\int_0^h u \frac{\partial M}{\partial r} dz + \int_0^h w \frac{\partial M}{\partial z} dz = -\frac{r\tau_{\theta s}}{\rho} = -C_d rV^2, \quad (10)$$

with $\tau_{\theta s} \approx C_d \rho V^2$ the surface stress, C_d a drag coefficient and V is the tangential surface wind component. Assuming $w(z = 0) = 0$ and the use of the continuity equation [Eq. (9)], the second term of the left-hand side in Eq. (10) is integrated by parts, following developments presented by Kalashnik (1994), to obtain

$$\int_0^h u \frac{\partial M}{\partial r} dz + (wM)|_{z=h} + \int_0^h \frac{M \partial(ru)}{r \partial r} dz = -C_d rV^2. \quad (11)$$

Grouping the two integrals yields

$$\frac{1}{r} \frac{d}{dr} \left(r \int_0^h u M dz \right) + (wM)|_{z=h} = -C_d rV^2. \quad (12)$$

Defining $\bar{u} := (1/h) \int_0^h u dz$ we can approximate the integral $\int_0^h u M dz \approx h \bar{u} M|_{z=h}$ and rewrite the continuity equation $w|_{z=h} = -[(h/r)(d/dr)(r\bar{u})]$. Rearrangement finally yields

$$rV^2 \approx -\frac{h\bar{u}}{C_d} \frac{dM}{dr}, \quad (13)$$

where dM/dr is the radial gradient of absolute angular momentum at the top of the boundary layer. Assuming the latter is closely related to its value at the surface, Eq. (13) then explicitly links the shape of the wind profile dM/dr to rV^2 .

Using SAR measurements, both quantities can be accurately estimated and the validity of Eq. (13) can be assessed. Figure 9a represents $R_{\max} V_{\max}^2$ (y axis) as a function of $R_{34} V_{34}^2$ (x axis) and colored by the fitted B values.² On average—that is, $B \simeq 1.8$ —a relationship emerges when comparing $R_{\max} V_{\max}^2$ and $R_{34} V_{34}^2$. Departures from a one-to-one relationship, related to conservation of the rV^2 parameter, are seemingly well explained by B values. Large B , corresponding to very peaked wind profiles near V_{\max} , leads to larger $R_{\max} V_{\max}^2$ for a given $R_{34} V_{34}^2$. For broader wind profiles, corresponding to smaller B , smaller $R_{\max} V_{\max}^2$ are generally found.

Moreover, the space spanned in the $(R_{34} V_{34}^2, R_{\max} V_{\max}^2)$ plane is still apparently large, even at constant B . From Eq. (13), this increased variability is possibly associated with the factor $h\bar{u}/C_d$. Overall, these results suggest that the variability encountered in nature does not solely depend on the three predictors (V_{\max} , R_{34} , and f).

To further illustrate this diagnosis, Fig. 9b displays the same $(R_{34} V_{34}^2, R_{\max} V_{\max}^2)$ plane, but using the radiometer and scatterometer database, and corresponding V_{\max}^{REG} , R_{34} , and $R_{\max}^{\text{CK22-BR}}$ estimates. As expected, the variability captured by using R_{34} or R_{50} to estimate R_{\max} via CK22 is poor. While using R_{64} increases this variability, the overall spread is reduced relative to Fig. 9a, suggesting that the variability of the wind profile shapes associated with the $R_{\max}^{\text{CK22-BR}}$ estimates is low.

Note that the average situation $rV^2 \approx \text{constant}$ that is depicted in our study because of the SAR database has already

² Here B , as a scalar value, was used instead of a criterion based on dM/dr to describe the shape of the wind profile.

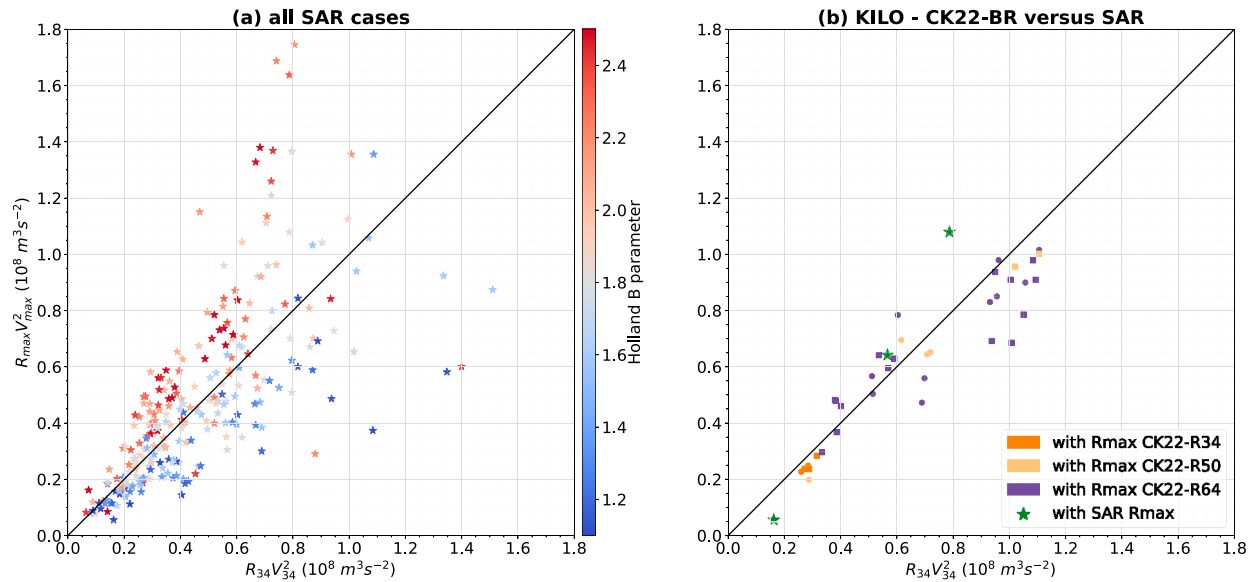


FIG. 9. Evaluation of the PV conservation assumption (a) in the SAR dataset and (b) for Kilo's life cycle using R_{34} estimated on radiometer and scatterometer data along with corresponding $R_{\max}^{\text{CK22-BR}}$ estimates and V_{\max}^{REG} . The three SAR cases (green stars) are also displayed for reference in (b).

been discussed by Riehl (1963) when he argued that PV is conserved within the inflow layer. The PV conservation implies that the vertical component of the curl of the frictional force is zero, or

$$\frac{r}{\rho} \frac{\partial \tau_{\theta z}}{\partial z} = \text{constant}. \quad (14)$$

Integrating this equation over the boundary layer height yields (assuming constant density)

$$\frac{r \tau_{\theta s}}{\rho} = C_d r V^2 = \text{constant}. \quad (15)$$

Thus, for a constant or slowly varying drag coefficient C_d , PV conservation leads to $rV^2 \approx \text{constant}$ (Riehl 1963). Mentioned above, such a relationship is, on average, consistent with the SAR estimates. However, for this relationship, the only source of variability comes from C_d . From arguments raised above [Eq. (13)], h and \bar{u} should also be further considered.

Last, one limitation of our observational analysis is that SAR V_{\max} is an estimate of the maximum total wind speed rather than the maximum tangential wind speed. Knowing how the total wind speed is distributed between its tangential and radial component near the eyewall region would allow to better estimate the impact of \bar{u} on PV conservation and its variability.

d. Comparison of $R_{\max}^{\text{CK22-BR}}$ with existing R_{\max} estimates

With these results in mind, we assessed how much $R_{\max}^{\text{CK22-BR}}$ estimates improved existing R_{\max} estimates. Figure 10 displays density contours of (V_{\max}, R_{\max}) joint distributions using IBTrACS R_{\max} (dashed blue) or $R_{\max}^{\text{CK22-BR}}$ estimates (solid black).

For comparison, the same density contours are shaded for the SAR dataset (green).

We remind readers that SAR wide-swath acquisitions cannot be continuously performed over the ocean. As a consequence, not only does the SAR dataset contain much fewer cases, it is also biased toward higher intensities. Indeed, acquisition orders are most often requested to observe higher intensity systems. Thus, for the lowest V_{\max} (less than $\sim 30 \text{ m s}^{-1}$), possible inconsistencies in R_{\max} densities arise when comparing SAR to radiometer and scatterometer or IBTrACS. The density contours suggest that both IBTrACS R_{\max} and $R_{\max}^{\text{CK22-BR}}$ estimates are larger than SAR R_{\max} , while, in fact, this is just a consequence of the lack of SAR data at these intensities.

Nevertheless, and more importantly for high surface winds, discrepancies in R_{\max} densities are observed. Indeed, on average IBTrACS density contours are centered on a higher R_{\max} ($\sim 30 \text{ km}$) than SAR (progressively decreasing to $\sim 20 \text{ km}$). Confirming the efficacy of the revised model, radiometer- and scatterometer-based density contours display an average $R_{\max}^{\text{CK22-BR}}$ ($\sim 20 \text{ km}$) that is consistent with SAR R_{\max} . Depicted by the R_{\max} density curves (right panel), for low R_{\max} , IBTrACS density is lower than both SAR and $R_{\max}^{\text{CK22-BR}}$ values.

For further comparison, we computed R_{\max} estimates from R_{34} on the radiometer and scatterometer data using Eq. (7) of Chavas and Knaff (2022). The corresponding density curve (dotted red) shows only a minor improvement relative to IBTrACS at low R_{\max} .

Last, the density contours of the radiometer and scatterometer dataset with $R_{\max}^{\text{CK22-BR}}$ span a larger space than IBTrACS in the (V_{\max}, R_{\max}) plane (cf. for instance the 20% contours in Fig. 10, i.e., the black and blue outermost contours). This shows that the former captures more variability than best-track data. This is

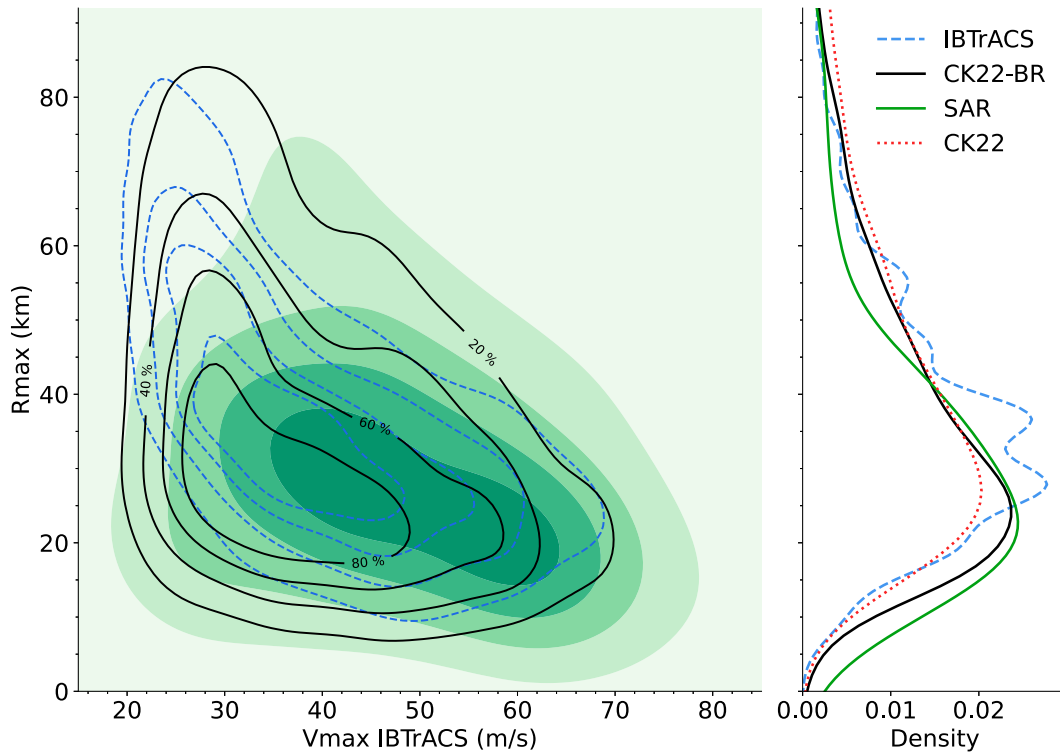


FIG. 10. Density contours of (V_{\max}, R_{\max}) joint distribution for the SAR dataset (shaded green; V_{\max} is based on IBTrACS), for the dataset based upon radiometers and scatterometers with $R_{\max}^{\text{CK22-BR}}$ (solid black) and based on corresponding IBTrACS R_{\max} values (dashed blue). All contours correspond to isoproportions (with 20% increments; see the black contour labels) in density obtained by two-dimensional Gaussian kernel density estimation. For instance, the area outside the 80% contour contains 80% of the probability mass. To the right, the corresponding R_{\max} density curves are displayed, along with R_{\max} estimates obtained by applying Eq. (7) of Chavas and Knaff (2022) to the radiometer- and scatterometer-based dataset (dotted red).

likely due to the use of R_{64} in the regression, a result already suggested by Fig. 9b. Even though the datasets do not have the same V_{\max} distributions, Fig. 10 also suggests that the radiometer and scatterometer density contours span less space than SAR observations in the (V_{\max}, R_{\max}) plane. While this is consistent with the above analysis, more SAR cases are needed to properly interpret Fig. 10.

6. Conclusions and perspectives

Understanding TC intensity changes certainly remains an observationally challenging problem. As expressed during the Tenth International Workshop on Tropical Cyclones (IWTC-10, recommendation 4), both the operational and research communities recognize the need for more homogeneous and standardized datasets for TC wind structure parameters, such as R_{\max} and the wind radii. The fact that R_{34} was not systematically reanalyzed in all basins, and that R_{\max} is still not reanalyzed today (best-track R_{\max} value typically stems from its operational estimate) hampers the consolidation of such a dataset. Systematic and standardized wind radii are needed when using, and further improving, a semiempirical model such as CK22. Although satellite sensors have their limitations, especially regarding the intercalibration of different

missions and sensors, resulting multimodal observations shall serve for such a systematic and global approach, at least for wind radii estimation.

More specifically and because of high-resolution (SAR) data, it is now possible to more systematically estimate R_{\max} . Fitted with SAR estimates and used in conjunction with the closest wind radius to R_{\max} , our study proposed a revised CK22 model. It is shown to be an efficient tool to provide improved reliable estimates, with an average uncertainty of ~9 km. Because outer-core wind radii can be estimated from radiometer or C-band scatterometer data, the developed framework thus allows to produce a more extensive dataset of reanalyzed R_{\max} estimates. The resulting time series are generally more realistic than those obtained from best-track R_{\max} estimates. The method can also be used to provide operational guidance on the location of the maximum intensity every time a radiometer or C-band scatterometer overflies the TC, as long as its intensity and location are also estimated, noting that such estimates are routinely available from operational centers. In fact, the developed framework is relevant to any situation where V_{\max} and an outer size are known and R_{\max} is biased or unknown. This includes low-resolution weather and climate modeling applications where the outer core (i.e., near R_{34}) is better resolved than the inner core (i.e., near R_{\max}),

and risk modeling with synthetic TCs (Gori et al. 2022) where V_{\max} and an outer size are commonly used as input, while R_{\max} must be predicted in conjunction with a wind profile model. The proposed method could also be used to guide the best-tracking process when no reliable R_{\max} observation is available.

The efficacy of the semiempirical CK22 model stems from fundamental conservation principles. Indeed, the high-resolution SAR database highlights that TCs, on average, conserve their PV, with a resulting approximation $rV^2 \approx \text{constant}$. Accordingly, the use of CK22 to retrieve R_{\max} , based on an outer-radius wind observation coupled with an intensity estimate is, on average, justified. Single cases can still depart considerably from the PV conservation assumption, especially those at very high intensity (V_{\max}) or with large inner (R_{\max}) or outer (R_{34}) size. In addition, to first order, those deviations are well explained by variations of the observed wind profile shapes.

While the use of R_{64} can account for some of the deviations due to the radial gradient of absolute angular momentum, the CK22 model seems to fail to capture the remaining variability observed in the SAR database. Large variability is apparently still occurring near the TC core. To further advance our understanding, there continues to be a need for spaceborne SAR and airborne SFMR sensors as these are the only tools that resolve surface winds in this area. Both sensors however suffer from a lack of spatiotemporal sampling, and airborne measurements suffer from a lack of azimuthal coverage. The future is bright with the recently launched RADARSAT Constellation Mission (RCM) operated by CSA, which should improve the satellite SAR spatiotemporal sampling. RCM has already proved useful by providing significantly more R_{\max} estimates than anticipated for the 2022/23 season. Also, increasing the number of available SAR cases will certainly allow us to better understand how absolute angular momentum gradients are constrained in the near-core region. Objective estimates of TC eye sizes or core sizes at intermediate levels are also routinely performed with spaceborne infrared or passive microwave data (Knapp et al. 2018; Cossuth 2014). While such information may complement SAR or SFMR surface observations in a multimodal approach, there still is a need to better understand how they relate to the TC wind structure.

Furthermore, the integrated equations show that both the boundary layer depth h , the average radial inflow \bar{u} , and the drag C_d also impact the relationship between PV conservation and the near-core wind profile shape. While the C_d behavior under very high winds is still actively debated (Powell et al. 2003; Bell et al. 2012; Donelan 2018; Curcic and Haus 2020), measurements of both h and \bar{u} may be facilitated by the Doppler-based motions derived from the Imaging Wind and Rain Airborne Profiler (IWRAP) instrument (Sapp et al. 2022). For the radial inflow, improved estimates at the surface, in the near-core region, shall be made possible with the future Harmony mission (ESA 2022), the ESA Earth Explorer 10. This mission will augment Sentinel-1D observations with two satellite companions, providing azimuth diversity from these bistatic observations. In addition, the Second Generation Meteorological Operational satellite program (Metop-SG) will operate in both co- and cross-polarization. Unlike the current spaceborne instruments, ASCAT, which have only co-polarization measurements, the

higher sensitivity of cross-polarized signals to ocean breaking waves may thus improve the ocean surface wind vectors measured by scatterometers, approaching the TC core regions. Also, the coming Copernicus Imaging Microwave Radiometer (CIMR) promises to offer large swath with improved resolution, low uncertainty observation capabilities, combining L-, C-, and X-band frequencies. The presence of 1.4 GHz L-band channel on board CIMR will open up the possibility to further interpret the high-resolution C- and X-band measurements, to provide improved surface wind vector estimates under extreme conditions (Kilic et al. 2018).

In the absence of high-resolution observations, the shape of the near-core wind profile may also be indirectly estimated. Given the relation $rV^2 \approx \text{constant}$ under a steady-state assumption, a departure from this relation can help one to understand the temporal variations of absolute angular momentum. Estimates of these temporal variations may then be used to evaluate how much the near-core wind profile shape departs from the average relationship. The wind profile shape is also linked to the drag coefficient [see for instance the steady-state view of Emanuel and Rotunno (2011)], which modulates asymmetries in the boundary layer response (Shapiro 1983; Kepert 2001). Asymmetries possibly captured by medium- or low-resolution observations (scatterometers or radiometers), may thus help to infer boundary layer frictional drag terms and to quantify the resulting shape of the wind profile.

Acknowledgments. This work was financially supported by the ERC Synergy Project 856408-STUOD, the ANR projects OceaniX and ISblue, and the ESA Marine Atmosphere eXtreme Satellite Synergy project (MAXSS). The radiometer and scatterometer dataset used in this report is part of the MAXSS project. The SAR database was obtained from IFREMER/CyclObs and produced with the SAR wind processor codeveloped by IFREMER and CLS. Author Knaff thanks NOAA/Center for Satellite Applications and Research for providing the time to work on this subject. The views, opinions, and findings contained in this report are those of the authors and should not be construed as an official National Oceanic and Atmospheric Administration or U.S. government position, policy, or decision. We also acknowledge three anonymous reviewers for their comments and especially for the enriching suggestions about infrared and passive microwave data.

Data availability statement. The data used in this study are freely available online for both the dataset of radiometer and scatterometer winds (<https://www.odatis-ocean.fr/donnees-et-services/acces-aux-donnees/catalogue-complet/#/metadata/6c56bcde-050f-42eb-92b8-8e882e1f4db9>) and the SAR database (<https://cyclobs.ifremer.fr/>).

APPENDIX A

Scatterometer Wind Speed Estimates

As explained in section 2, the wind speed estimates from different radiometer and scatterometer sensors have been intercalibrated prior to our study. During this process, the C-band ASCAT missions were calibrated using a 25-km resolution,

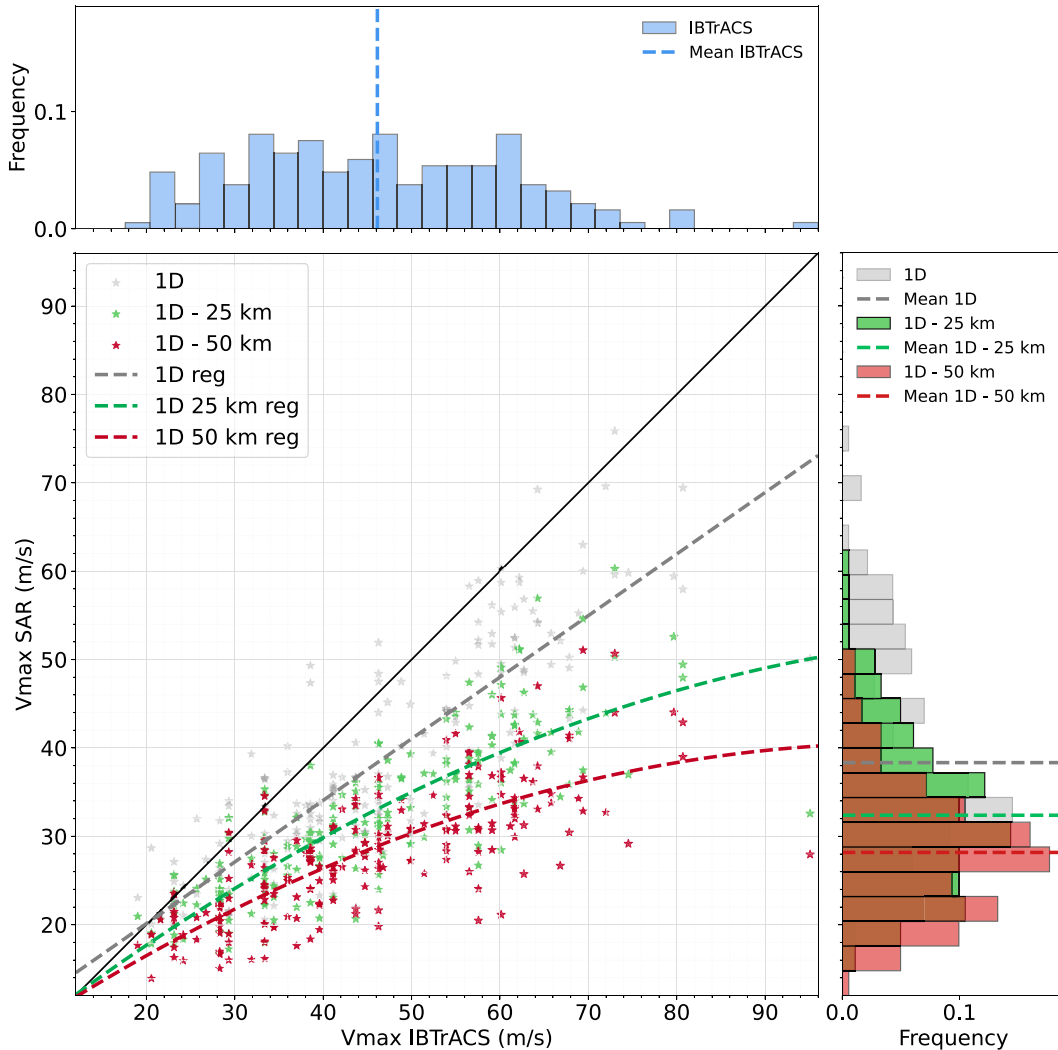


FIG. A1. Comparison between SAR (y axis) and IBTrACS (x axis) V_{\max} for the raw dataset (gray) and when degraded at 25-km (green) or 50-km (red) resolution. Dashed lines represent best linear fit for the raw dataset (gray) and best second-order polynomial fits for the 25-km (green) and 50-km (red) datasets. A solid black line represents identity. The V_{\max} distributions and averages are displayed to the right for the different SAR samples and at the top for corresponding IBTrACS values.

while the Ku-band scatterometer sensors were calibrated using a 50-km resolution. Spatial resolution was already demonstrated to impact how well TCs intensities are resolved in numerical models (Davis 2018) and observations (Quilfen et al. 1998). Here, we expect discrepancies between the C- and Ku-band observational wind products.

To quantify this resolution effect, SAR wind fields were degraded to both 25- and 50-km spatial resolution and then azimuthally averaged. The V_{\max} values estimated from these degraded wind profiles were then compared with IBTrACS V_{\max} , as represented by the green (25 km) and red (50 km) stars of Fig. A1. Here, SAR V_{\max} refers to the maximum found in an azimuthally averaged wind profile. We thus expect slight discrepancies with IBTrACS V_{\max} , whose definition does not strictly coincide with a wind profile maximum. The comparison

between SAR azimuthal means and IBTrACS is indicated by the gray stars and is modeled by a linear fit (gray dashed line in Fig. A1) that defines V_{\max}^{REG} :

$$V_{\max}^{\text{REG}} = 0.6967V_{\max}^{\text{IBTrACS}} + 6.1992. \quad (\text{A1})$$

The green and red scatters in Fig. A1 should be compared with this regression line (gray dashed) rather than the 1:1 line. The 25- and 50-km simulated V_{\max} values show that as spatial resolution decreases V_{\max} also decreases, and the decreasing tendency is more pronounced as intensity increases. On average, a V_{\max} of $\sim 38 \text{ m s}^{-1}$ observed at the full-resolution azimuthally averaged wind profile (i.e., the raw SAR wind profile) would yield $\sim 32 \text{ m s}^{-1}$ when observed at a 25-km spatial resolution and $\sim 28 \text{ m s}^{-1}$ at a 50-km spatial resolution. Second-order polynomial fits were constructed to model this spatial resolution effect.

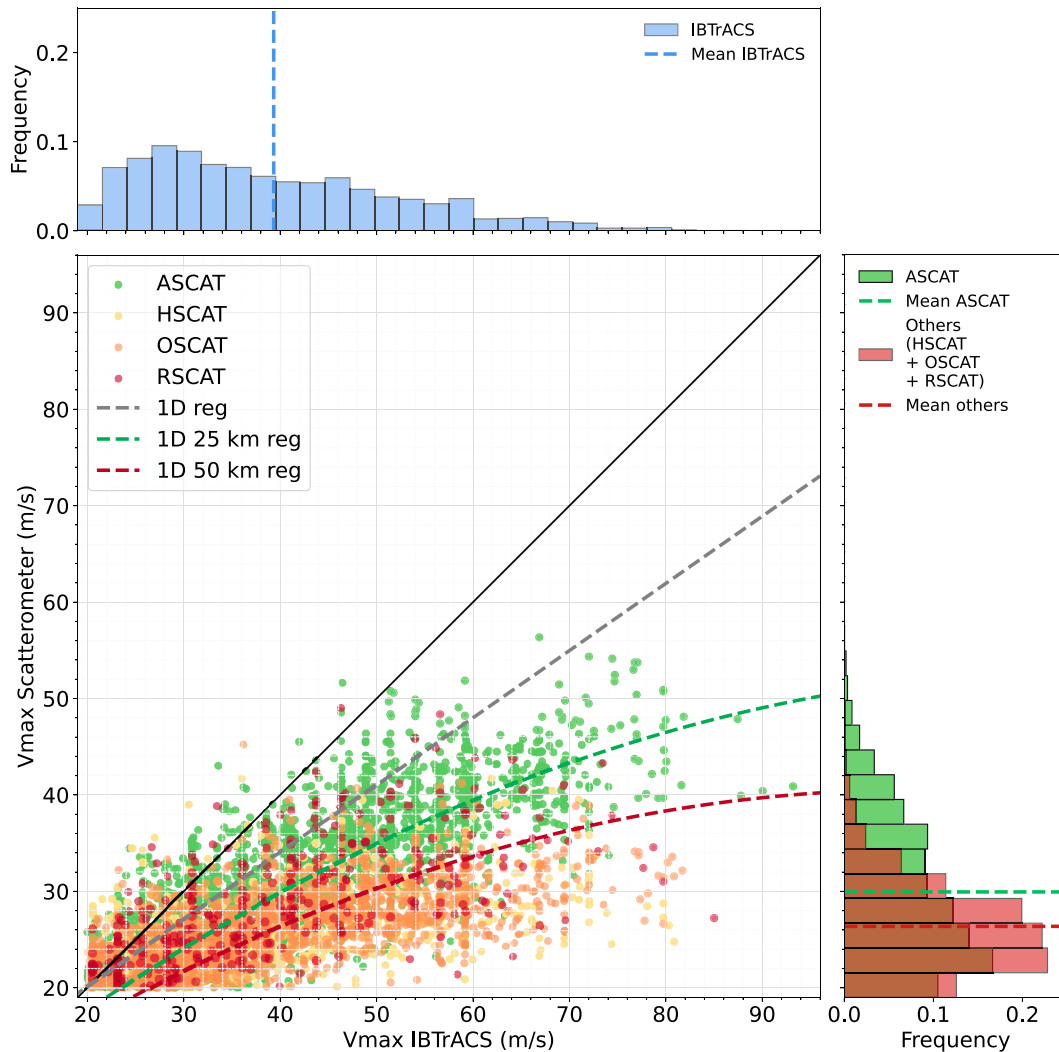


FIG. A2. Comparison between scatterometer (y axis) and IBTrACS (x axis) V_{\max} for ASCAT (green), HSCAT (yellow), OSCAT (orange), and RSCAT (red). Solid and dashed lines are identical to those in Fig. A1. The V_{\max} distributions and averages are displayed to the right for the different scatterometer datasets and at the top for corresponding IBTrACS values.

Using these linear and polynomial fits as reference, we then compared C-band and Ku-band scatterometer V_{\max} values with IBTrACS in Fig. A2. It shows that C-band scatterometer V_{\max} values are consistent with the 25-km spatial resolution polynomial model (green dashed curve). In contrast, Ku-band scatterometer V_{\max} are still underestimated when compared with IBTrACS values following the correction for their 50-km resolution (red dashed curve). In particular, Ku-band scatterometer V_{\max} estimates rarely exceed 64 kt (33 m s^{-1}), precluding their use to estimate wind radii in our analysis.

REFERENCES

- Bell, M. M., M. T. Montgomery, and K. A. Emanuel, 2012: Air-sea enthalpy and momentum exchange at major hurricane wind speeds observed during CBLAST. *J. Atmos. Sci.*, **69**, 3197–3222, <https://doi.org/10.1175/JAS-D-11-0276.1>.
- Brennan, M. J., C. C. Hennon, and R. D. Knabb, 2009: The operational use of QuikSCAT ocean surface vector winds at the National Hurricane Center. *Wea. Forecasting*, **24**, 621–645, <https://doi.org/10.1175/2008WAF2222188.1>.
- Chavas, D. R., and N. Lin, 2016: A model for the complete radial structure of the tropical cyclone wind field. Part II: Wind field variability. *J. Atmos. Sci.*, **73**, 3093–3113, <https://doi.org/10.1175/JAS-D-15-0185.1>.
- , and J. A. Knaff, 2022: A simple model for predicting the tropical cyclone radius of maximum wind from outer size. *Wea. Forecasting*, **37**, 563–579, <https://doi.org/10.1175/WAF-D-21-0103.1>.
- , N. Lin, and K. Emanuel, 2015: A model for the complete radial structure of the tropical cyclone wind field. Part I: Comparison with observed structure. *J. Atmos. Sci.*, **72**, 3647–3662, <https://doi.org/10.1175/JAS-D-15-0014.1>.
- Chou, K.-H., C.-C. Wu, and S.-Z. Lin, 2013: Assessment of the ASCAT wind error characteristics by global dropwindsonde

- observations. *J. Geophys. Res. Atmos.*, **118**, 9011–9021, <https://doi.org/10.1002/jgrd.50724>.
- Combot, C., A. Mouche, J. Knaff, Y. Zhao, Y. Zhao, L. Vinour, Y. Quilfen, and B. Chapron, 2020a: Extensive high-resolution Synthetic Aperture Radar (SAR) data analysis of tropical cyclones: Comparisons with SFMR flights and best track. *Mon. Wea. Rev.*, **148**, 4545–4563, <https://doi.org/10.1175/MWR-D-20-0005.1>.
- , Y. Quilfen, A. Mouche, J. Gourrion, C. de Boyer Montégut, B. Chapron, and J. Tournadre, 2020b: Space-based observations of surface signatures in the wakes of the 2018 eastern Pacific tropical cyclones. *J. Oper. Oceanogr.*, **13** (Suppl. 1), S132–S137.
- Cossuth, J. H., 2014: Exploring a comparative climatology of tropical cyclone core structures. Ph.D. dissertation, The Florida State University, 201 pp.
- Curcic, M., and B. K. Haus, 2020: Revised estimates of ocean surface drag in strong winds. *Geophys. Res. Lett.*, **47**, e2020GL087647, <https://doi.org/10.1029/2020GL087647>.
- Davis, C. A., 2018: Resolving tropical cyclone intensity in models. *Geophys. Res. Lett.*, **45**, 2082–2087, <https://doi.org/10.1002/2017GL076966>.
- Donelan, M. A., 2018: On the decrease of the oceanic drag coefficient in high winds. *J. Geophys. Res. Oceans*, **123**, 1485–1501, <https://doi.org/10.1002/2017JC013394>.
- Donnelly, W. J., J. R. Carswell, R. E. McIntosh, P. S. Chang, J. Wilkerson, F. Marks, and P. G. Black, 1999: Revised ocean backscatter models at C and Ku band under high-wind conditions. *J. Geophys. Res.*, **104**, 11 485–11 497, <https://doi.org/10.1029/1998JC900030>.
- Emanuel, K., 2004: Tropical cyclone energetics and structure. *Atmospheric Turbulence and Mesoscale Meteorology: Scientific Research Inspired by Doug Lilly*, E. Fedorovich, R. Rotunno, and B. Stevens, Eds., Cambridge University Press, 165–192.
- , and R. Rotunno, 2011: Self-stratification of tropical cyclone outflow. Part I: Implications for storm structure. *J. Atmos. Sci.*, **68**, 2236–2249, <https://doi.org/10.1175/JAS-D-10-05024.1>.
- ESA, 2022: Report for mission selection: Earth Explorer 10 candidate mission harmony. ESA Tech. Rep., 369 pp.
- Ginis, I., 2002: Tropical cyclone-ocean interactions. *Atmosphere-Ocean Interactions*, M. Rahman, R. Verhoeven, and C. A. Brebbia, Eds., Advances in Fluid Mechanics Series, Vol. 33, WIT Press, 83–114.
- Gori, A., N. Lin, D. Xi, and K. Emanuel, 2022: Tropical cyclone climatology change greatly exacerbates U.S. extreme rainfall-surge hazard. *Nat. Climate Change*, **12**, 171–178, <https://doi.org/10.1038/s41558-021-01272-7>.
- Holland, G. J., 1980: An analytic model of the wind and pressure profiles in hurricanes. *Mon. Wea. Rev.*, **108**, 1212–1218, [https://doi.org/10.1175/1520-0493\(1980\)108<1212:AAMOTW>2.0.CO;2](https://doi.org/10.1175/1520-0493(1980)108<1212:AAMOTW>2.0.CO;2).
- Irish, J. L., D. T. Resio, and J. J. Ratcliff, 2008: The influence of storm size on hurricane surge. *J. Phys. Oceanogr.*, **38**, 2003–2013, <https://doi.org/10.1175/2008JPO3727.1>.
- Kalashnik, M., 1994: On the maximum wind velocity in the tropical cyclone. *Izv. Akad. Nauk SSSR. Fiz. Atmos. I Okeana*, **30**, 26–30.
- Keper, J., 2001: The dynamics of boundary layer jets within the tropical cyclone core. Part I: Linear theory. *J. Atmos. Sci.*, **58**, 2469–2484, [https://doi.org/10.1175/1520-0469\(2001\)058<2469:TDOBLJ>2.0.CO;2](https://doi.org/10.1175/1520-0469(2001)058<2469:TDOBLJ>2.0.CO;2).
- Kilic, L., and Coauthors, 2018: Expected performances of the Copernicus Imaging Microwave Radiometer (CIMR) for an all-weather and high spatial resolution estimation of ocean and sea ice parameters. *J. Geophys. Res. Oceans*, **123**, 7564–7580, <https://doi.org/10.1029/2018JC014408>.
- Klotz, B. W., and E. W. Uhlhorn, 2014: Improved Stepped Frequency Microwave Radiometer tropical cyclone surface winds in heavy precipitation. *J. Atmos. Oceanic Technol.*, **31**, 2392–2408, <https://doi.org/10.1175/JTECH-D-14-00028.1>.
- Knaff, J. A., C. R. Sampson, P. J. Fitzpatrick, Y. Jin, and C. M. Hill, 2011: Simple diagnosis of tropical cyclone structure via pressure gradients. *Wea. Forecasting*, **26**, 1020–1031, <https://doi.org/10.1175/WAF-D-11-00013.1>.
- , and Coauthors, 2021: Estimating tropical cyclone surface winds: Current status, emerging technologies, historical evolution, and a look to the future. *Trop. Cyclone Res. Rev.*, **10**, 125–150, <https://doi.org/10.1016/j.tcr.2021.09.002>.
- Knapp, K. R., M. C. Kruk, D. H. Levinson, H. J. Diamond, and C. J. Neumann, 2010: The International Best Track Archive for Climate Stewardship (IBTrACS) unifying tropical cyclone data. *Bull. Amer. Meteor. Soc.*, **91**, 363–376, <https://doi.org/10.1175/2009BAMS2755.1>.
- , C. S. Velden, and A. J. Wimmers, 2018: A global climatology of tropical cyclone eyes. *Mon. Wea. Rev.*, **146**, 2089–2101, <https://doi.org/10.1175/MWR-D-17-0343.1>.
- Kossin, J. P., J. A. Knaff, H. I. Berger, D. C. Herndon, T. A. Cram, C. S. Velden, R. J. Murnane, and J. D. Hawkins, 2007: Estimating hurricane wind structure in the absence of aircraft reconnaissance. *Wea. Forecasting*, **22**, 89–101, <https://doi.org/10.1175/WAF985.1>.
- Krien, Y., and Coauthors, 2018: Can we improve parametric cyclonic wind fields using recent satellite remote sensing data? *Remote Sens.*, **10**, 1963, <https://doi.org/10.3390/rs10121963>.
- Kudryavtsev, V., A. Monzikova, C. Combot, B. Chapron, and N. Reul, 2019: A simplified model for the baroclinic and barotropic ocean response to moving tropical cyclones: 2. Model and simulations. *J. Geophys. Res. Oceans*, **124**, 3462–3485, <https://doi.org/10.1029/2018JC014747>.
- , M. Yurovskaya, and B. Chapron, 2021: Self-similarity of surface wave developments under tropical cyclones. *J. Geophys. Res. Oceans*, **126**, e2020JC016916, <https://doi.org/10.1029/2020JC016916>.
- Landsea, C. W., and J. L. Franklin, 2013: Atlantic hurricane database uncertainty and presentation of a new database format. *Mon. Wea. Rev.*, **141**, 3576–3592, <https://doi.org/10.1175/MWR-D-12-00254.1>.
- Manaster, A., L. Ricciardulli, and T. Meissner, 2021: Tropical cyclone winds from WindSat, AMSR2, and SMAP: Comparison with the HWRF model. *Remote Sens.*, **13**, 2347, <https://doi.org/10.3390/rs13122347>.
- Mayers, D. R., C. S. Ruf, and A. M. Warnock, 2023: CYGNSS storm-centric tropical cyclone gridded wind speed product. *J. Appl. Meteor. Climatol.*, **62**, 329–339, <https://doi.org/10.1175/JAMC-D-22-0054.1>.
- Meissner, T., L. Ricciardulli, and F. J. Wentz, 2017: Capability of the SMAP mission to measure ocean surface winds in storms. *Bull. Amer. Meteor. Soc.*, **98**, 1660–1677, <https://doi.org/10.1175/BAMS-D-16-0052.1>.
- , —, and A. Manaster, 2021: Tropical cyclone wind speeds from WindSat, AMSR and SMAP: Algorithm development and testing. *Remote Sens.*, **13**, 1641, <https://doi.org/10.3390/rs13091641>.
- Morris, M., and C. S. Ruf, 2017: Determining tropical cyclone surface wind speed structure and intensity with the CYGNSS satellite constellation. *J. Appl. Meteor. Climatol.*, **56**, 1847–1865, <https://doi.org/10.1175/JAMC-D-16-0375.1>.

- Mouche, A., B. Chapron, B. Zhang, and R. Husson, 2017: Combined co-and cross-polarized SAR measurements under extreme wind conditions. *IEEE Trans. Geosci. Remote Sens.*, **55**, 6746–6755, <https://doi.org/10.1109/TGRS.2017.2732508>.
- , —, J. Knaff, Y. Zhao, B. Zhang, and C. Combot, 2019: Copolarized and cross-polarized SAR measurements for high-resolution description of major hurricane wind structures: Application to Irma category 5 hurricane. *J. Geophys. Res. Oceans*, **124**, 3905–3922, <https://doi.org/10.1029/2019JC015056>.
- Mueller, K. J., M. DeMaria, J. Knaff, J. P. Kossin, and T. H. Vonder Haar, 2006: Objective estimation of tropical cyclone wind structure from infrared satellite data. *Wea. Forecasting*, **21**, 990–1005, <https://doi.org/10.1175/WAF955.1>.
- Polverari, F., M. Portabella, W. Lin, J. W. Sapp, A. Stoffelen, Z. Jelenak, and P. S. Chang, 2022: On high and extreme wind calibration using ASCAT. *IEEE Trans. Geosci. Remote Sens.*, **60**, 1–10, <https://doi.org/10.1109/TGRS.2021.3079898>.
- Portabella, M., A. S. Rabaneda, and G. Grieco, 2022: MAXSS: Algorithm theoretical baseline document for SFMR-based satellite-derived extreme wind recalibration, version 2.0. 43 pp., <https://digital.csic.es/bitstream/10261/287615/1/Por.pdf>.
- Powell, M. D., P. J. Vickery, and T. A. Reinhold, 2003: Reduced drag coefficient for high wind speeds in tropical cyclones. *Nature*, **422**, 279–283, <https://doi.org/10.1038/nature01481>.
- Price, J. F., 1981: Upper ocean response to a hurricane. *J. Phys. Oceanogr.*, **11**, 153–175, [https://doi.org/10.1175/1520-0485\(1981\)011<0153:UORTAH>2.0.CO;2](https://doi.org/10.1175/1520-0485(1981)011<0153:UORTAH>2.0.CO;2).
- Quilfen, Y., B. Chapron, T. Elfouhaily, K. Katsaros, and J. Tournaire, 1998: Observation of tropical cyclones by high-resolution scatterometry. *J. Geophys. Res.*, **103**, 7767–7786, <https://doi.org/10.1029/97JC01911>.
- , C. Prigent, B. Chapron, A. A. Mouche, and N. Houti, 2007: The potential of QuikSCAT and WindSat observations for the estimation of sea surface wind vector under severe weather conditions. *J. Geophys. Res.*, **112**, C09023, <https://doi.org/10.1029/2007JC004163>.
- Reul, N., and B. Chapron, 2003: A model of sea-foam thickness distribution for passive microwave remote sensing applications. *J. Geophys. Res.*, **108**, 3321, <https://doi.org/10.1029/2003JC001887>.
- , J. Tenerelli, B. Chapron, D. Vandemark, Y. Quilfen, and Y. Kerr, 2012: SMOS satellite L-band radiometer: A new capability for ocean surface remote sensing in hurricanes. *J. Geophys. Res.*, **117**, C02006, <https://doi.org/10.1029/2011JC007474>.
- , B. Chapron, E. Zabolotskikh, C. Donlon, Y. Quilfen, S. Guimbard, and J.-F. Piolle, 2016: A revised L-band radiobrightness sensitivity to extreme winds under tropical cyclones: The five year SMOS-storm database. *Remote Sens. Environ.*, **180**, 274–291, <https://doi.org/10.1016/j.rse.2016.03.011>.
- , and Coauthors, 2017: A new generation of tropical cyclone size measurements from space. *Bull. Amer. Meteor. Soc.*, **98**, 2367–2385, <https://doi.org/10.1175/BAMS-D-15-00291.1>.
- Riehl, H., 1963: Some relations between wind and thermal structure of steady state hurricanes. *J. Atmos. Sci.*, **20**, 276–287, [https://doi.org/10.1175/1520-0469\(1963\)020<0276:SRBWAT>2.0.CO;2](https://doi.org/10.1175/1520-0469(1963)020<0276:SRBWAT>2.0.CO;2).
- Sampson, C. R., E. M. Fukada, J. A. Knaff, B. R. Strahl, M. J. Brennan, and T. Marchok, 2017: Tropical cyclone gale wind radii estimates for the western North Pacific. *Wea. Forecasting*, **32**, 1029–1040, <https://doi.org/10.1175/WAF-D-16-0196.1>.
- Sapp, J., Z. Jelenak, P. Chang, C. Shoup, and J. Carswell, 2022: Processing of high-resolution Hurricane Ida boundary layer winds from the IWRAP instrument on the NOAA WP-3D aircraft. *IGARSS 2022—2022 IEEE Int. Geoscience and Remote Sensing Symp.*, Kuala Lumpur, Malaysia, Institute of Electrical and Electronics Engineers, 7286–7289, <https://doi.org/10.1109/IGARSS46834.2022.9883259>.
- Schreck, C. J., III, K. R. Knapp, and J. P. Kossin, 2014: The impact of best track discrepancies on global tropical cyclone climatologies using IBTrACS. *Mon. Wea. Rev.*, **142**, 3881–3899, <https://doi.org/10.1175/MWR-D-14-00021.1>.
- Shapiro, L. J., 1983: The asymmetric boundary layer flow under a translating hurricane. *J. Atmos. Sci.*, **40**, 1984–1998, [https://doi.org/10.1175/1520-0469\(1983\)040<1984:TABLEFU>2.0.CO;2](https://doi.org/10.1175/1520-0469(1983)040<1984:TABLEFU>2.0.CO;2).
- Soisuvarn, S., Z. Jelenak, P. S. Chang, S. O. Alsweiss, and Q. Zhu, 2013: CMOD5.H—A high wind geophysical model function for C-band vertically polarized satellite scatterometer measurements. *IEEE Trans. Geosci. Remote Sens.*, **51**, 3744–3760, <https://doi.org/10.1109/TGRS.2012.2219871>.
- Stoffelen, A., J. A. Verspeek, J. Vogelzang, and A. Verhoef, 2017: The CMOD7 geophysical model function for ASCAT and ERS wind retrievals. *IEEE J. Sel. Top. Appl. Earth Obs. Remote Sens.*, **10**, 2123–2134, <https://doi.org/10.1109/JSTARS.2017.2681806>.
- Takagi, H., and W. Wu, 2016: Maximum wind radius estimated by the 50 kt radius: Improvement of storm surge forecasting over the western North Pacific. *Nat. Hazards Earth Syst. Sci.*, **16**, 705–717, <https://doi.org/10.5194/nhess-16-705-2016>.
- Tsukada, T., and T. Horinouchi, 2023: Strong relationship between eye radius and radius of maximum wind of tropical cyclones. *Mon. Wea. Rev.*, **151**, 569–588, <https://doi.org/10.1175/MWR-D-22-0106.1>.
- Uhlhorn, E. W., P. G. Black, J. L. Franklin, M. Goodberlet, J. Carswell, and A. S. Goldstein, 2007: Hurricane surface wind measurements from an operational Stepped Frequency Microwave Radiometer. *Mon. Wea. Rev.*, **135**, 3070–3085, <https://doi.org/10.1175/MWR3454.1>.
- Vickery, P. J., and D. Wadhera, 2008: Statistical models of Holland pressure profile parameter and radius to maximum winds of hurricanes from flight-level pressure and H*wind data. *J. Appl. Meteor. Climatol.*, **47**, 2497–2517, <https://doi.org/10.1175/2008JAMC1837.1>.
- Vinour, L., S. Jullien, A. Mouche, C. Combot, and M. Mangeas, 2021: Observations of tropical cyclone inner-core fine-scale structure, and its link to intensity variations. *J. Atmos. Sci.*, **78**, 3651–3671, <https://doi.org/10.1175/JAS-D-20-0245.1>.
- Wang, S., and R. Toumi, 2021: Recent tropical cyclone changes inferred from ocean surface temperature cold wakes. *Sci. Rep.*, **11**, 22269, <https://doi.org/10.1038/s41598-021-01612-9>.
- Willoughby, H. E., and M. E. Rahn, 2004: Parametric representation of the primary hurricane vortex. Part I: Observations and evaluation of the Holland (1980) model. *Mon. Wea. Rev.*, **132**, 3033–3048, <https://doi.org/10.1175/MWR2831.1>.
- , R. Darling, and M. E. Rahn, 2006: Parametric representation of the primary hurricane vortex. Part II: A new family of sectionally continuous profiles. *Mon. Wea. Rev.*, **134**, 1102–1120, <https://doi.org/10.1175/MWR3106.1>.
- Wright, C. W., and Coauthors, 2001: Hurricane directional wave spectrum spatial variation in the open ocean. *J. Phys. Oceanogr.*, **31**,

- 2472–2488, [https://doi.org/10.1175/1520-0485\(2001\)031<2472:HDWSSV>2.0.CO;2](https://doi.org/10.1175/1520-0485(2001)031<2472:HDWSSV>2.0.CO;2).
- Young, I. R., 2017: A review of parametric descriptions of tropical cyclone wind-wave generation. *Atmosphere*, **8**, 194, <https://doi.org/10.3390/atmos8100194>.
- Yueh, S. H., A. G. Fore, W. Tang, A. Hayashi, B. Stiles, N. Reul, Y. Weng, and F. Zhang, 2016: SMAP L-band passive microwave observations of ocean surface wind during severe storms. *IEEE Trans. Geosci. Remote Sens.*, **54**, 7339–7350, <https://doi.org/10.1109/TGRS.2016.2600239>.
- Zabolotskikh, E. V., L. M. Mitnik, N. Reul, and B. Chapron, 2015: New possibilities for geophysical parameter retrievals opened by GCOM-W1 AMSR2. *IEEE J. Sel. Top. Appl. Earth Obs. Remote Sens.*, **8**, 4248–4261, <https://doi.org/10.1109/JSTARS.2015.2416514>.
- Zhao, Y., A. A. Mouche, B. Chapron, and N. Reul, 2018: Direct comparison between active C-band radar and passive L-band radiometer measurements: Extreme event cases. *IEEE Geosci. Remote Sens. Lett.*, **15**, 897–901, <https://doi.org/10.1109/LGRS.2018.2811712>.