

1 **The hydrometeor partitioning and microphysical processes over the Pacific Warm Pool in**
2 **numerical modeling**

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27 September 2016

28 Submitted for publication for Atmospheric Research
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Abstract

Numerical modeling is conducted to study the hydrometeor partitioning and microphysical source and sink processes during a quasi-steady state of thunderstorms over the Pacific Warm Pool by utilizing the microphysical model WISCDYMM to simulate selected storm cases. The results show that liquid-phase hydrometeors dominate thunderstorm evolution over the Pacific Warm Pool. The ratio of ice-phase mass to liquid-phase mass is about 41% : 59%, indicating that ice-phase water is not as significant over the Pacific Warm Pool as the liquid water compared to the larger than 50% in the subtropics and ~ 80% in US High Plains in a previous study. Sensitivity tests support the dominance of liquid-phase hydrometeors over the Pacific Warm Pool.

The major rain sources are the key hail sinks: melting of hail and shedding from hail; whereas the crucial rain sinks are evaporation and accretion by hail. The major snow sources are Bergeron-Findeisen process, transfer of cloud ice to snow and accretion of cloud water; whereas the foremost sink of snow is accretion by hail. The essential hail sources are accretions of rain, cloud water, and snow; whereas the critical hail sinks are melting of hail and shedding from hail. The contribution and ranking of sources and sinks of these precipitates are compared with the previous study.

Hydrometeors have their own special microphysical processes in the development and depletion over the Pacific Warm Pool. Microphysical budgets depend on atmospheric dynamical and thermodynamical conditions which determine the partitioning of hydrometeors. This knowledge would benefit the microphysics parameterization in cloud models and cumulus parameterization in global circulation models.

70 1. Introduction

71 Poor precipitation performance in general circulation models (GCMs) has been a serious
72 problem in the atmospheric science community. The parameterizations of cloud species have
73 close links to convection and precipitation modeling, which is closely related to radiation and
74 large-scale circulation in GCMs, and thus has large impacts on climate modeling as well.
75 Successful numerical modeling of precipitation in GCMs depends very much on model
76 resolution and cloud parameterization. It was suggested that if the resolution of GCMs can
77 resolve convective systems realistically, the problems on the numerical modeling of precipitation
78 and tropical cyclones will meliorate (Hourdin et al. 2006; Randall et al. 2007).

79 If the convective cloud parameterization in GCM can include the effect of the
80 Microphysical budget of convective clouds in different geographic zones, the predictions of
81 convective cloud properties in GCM would be greatly improved. The microphysical budgets
82 such as hydrometeor partitioning, microphysical processes, etc. in the quasi-steady state in
83 mature convective storms in the US High Plains and some subtropical locations have been
84 previously investigated by Lin et al. (2005, hereafter LWS05). Their results show that
85 partitioning of hydrometeors in deep convective clouds is sensitive to geographical location. At
86 the present stage, it is impossible to include detailed cloud microphysical processes in a GCM
87 and hence it is useful to parameterize the cloud partitioning in deep convective clouds that take
88 the geographic location into account.

89 One of the regions of great importance to climate studies is the Pacific Warm Pool (Wang
90 and Xie, 1998; Barlow et al., 2002). The Pacific Warm Pool plays a significant role in global
91 heat and water vapor transport; hence, the understanding of the cloud properties including how
92 hydrometeors are partitioned is of great importance. There has been some research related to the
93 hydrometeor partitioning over the Pacific Warm Pool. Chen and Yin (2011) showed the
94 hydrometeor partitioning in a case study in Darwin using a three-dimensional nonhydrostatic
95 convective cloud model with a double-moment bulk microphysics scheme and explicit droplet
96 activation from cloud condensation nuclei. Based on cloud budget employing a two-dimensional
97 cloud-resolving model and simulation data during TOGA COARE, Li et al. (2011) obtained net
98 condensation and hydrometeor change/convergence. With the simulations triggered by large-
99 scale forcing data, Li and Shen (2013) examined the rain microphysical processes connected
100 with precipitation in the tropical deep convection using a 2-D cloud model. Based on the
101 measurements of the shape and electric charge of hydrometeors by a videosonde, Takahashi
102 (2004) constructed a conceptual model of hydrometeor mass, number, and space charge
103 evolution for convective clouds around Tiwi Islands.

104 In the above studies, the hydrometeor partitioning was only about a certain sounding, and
105 there is no study of microphysical processes over the Pacific Warm Pool. In this study, we utilize
106 a numerical cloud model to examine the dynamic and thermodynamic fields, and the details of
107 hydrometeors and microphysical processes during the quasi-steady state of some thunderstorms
108 over the Pacific Warm Pool. We also study the differences in hydrometeor partitioning, and
109 microphysical processes among the thunderstorms in the Pacific Warm Pool, subtropics and High
110 Plains in LWS05.

111 This paper is organized as follows. Section 2 introduces the cloud model employed in
112 this study. Section 3 describes three Pacific Warm Pool storm cases and verifies the numerical
113 simulations by observations. In section 4 we investigate the time evolution of microphysical
114 budgets of them. Section 5 compares the specialties of microphysical budgets in the storms in
115 Pacific Warm Pool with subtropics and High Plains in LWS05. Summary is given in Section 6.

116

117 2. Model description

118 The numerical model used in this study is the Wisconsin Dynamical/Microphysical
119 Model (WISCDYMM), which is a three-dimensional, quasi-compressible, non-hydrostatic
120 primitive-equation cloud model taken from Anderson et al. (1985) (Straka 1989; Johnson et al.,
121 1993; 1995; Wang, 2003). The microphysics of this model is based on those in Lin et al. (1983)
122 and Cotton et al. (1982, 1986). The hydrometeors have six categories: water vapor, cloud water,
123 cloud ice, rain, snow, and graupel/hail. The bulk method is applied to parameterize 38
124 microphysical processes including autoconversion, nucleation, condensation, accretion,
125 evaporation, freezing, melting, sublimation, deposition, etc. Except for cloud ice which is
126 assumed to be hexagonal plates, hydrometeors are assumed to be spherical. The cloud water is
127 supposed to be monodispersed with a number concentration; whereas cloud ice is assumed to be
128 monodispersed as a function of temperature. The rain, snow, and graupel/hail are assumed to be
129 inverse exponential size distributions (Marshall and Palmer, 1948; Gunn and Marshall, 1958;
130 Federer and Waldvogel, 1975). The precipitation in all types of water species is assumed to fall
131 with mass-weighted mean terminal speed.

132 There are twelve prognostic variables containing velocity in x, y, z directions, pressure,
133 potential temperature, turbulent kinetic energy, mixing ratio of water vapor, bulk cloud water,
134 bulk cloud ice, bulk rain water, bulk snow aggregates, and bulk graupel/hail (Straka 1989). The
135 Arakawa-C staggered grid (Arakawa and Lamb, 1981) is utilized in this model. The domain is
136 $64 \text{ km} \times 64 \text{ km} \times 25 \text{ km}$ with horizontal grid interval 500 m, and vertical interval 200 m. The
137 time interval of integration is 2 seconds. The horizontal wind is subtracted, changing every 30
138 mins, to confine the convective activities within the computational domain. The upstream sixth-
139 order, flux-conservative Crowley scheme is employed in advection terms (Tremback et al., 1987).
140 Predictive variables are filtered every time step to handle nonlinear instability by a fourth-order
141 numerical spatial diffusion operator (Klemp and Wilhelmson, 1978). A time filter is used in
142 prognostic variables in the leapfrog scheme from odd and even time steps (Asselin, 1972). The
143 radiation condition is utilized at the lateral boundaries to let patterns go out of domain without
144 producing disturbances (Klemp and Wilhelmson, 1978). At the top boundary, all variables are
145 given the values of the base state. A Rayleigh sponge layer is imposed to reduce gravity waves
146 close to the top of the domain (Clark, 1977).

147 All the hydrometeor concentrations are shown as mixing ratios. Negative numbers of
148 them do not set to null nor participate in any calculation. Such a design is to avoid false
149 accumulation of water species, which happens when negative hydrometeors mixing ratios are set
150 to zero. The convection is initiated by a thermal bubble, whose center is 2 km above ground
151 level, with the maximum 3.5 K at bubble center and the radius 10 km and depth 4 km in the
152 surrounding of horizontal homogeneousness. The relative humidity of the bubble is set to be the
153 same as outside the thermal bubble by varying mixing ratio. A sensitivity study of the thermal
154 bubble size was conducted and the results show that, while the numbers change somewhat, the
155 main conclusions remain the same (see Sec. 5).

156

157 3. Simulation of the case Darwin

158 The three cases studied in this research with heavy rainfall over the Pacific Warm Pool
159 are January 23, 2006, in Darwin, March 28, 2013, in Singapore, and July 18, 2013, in Medan.
160 The Skew-T log-P soundings and wind fields of these storms are employed as the initial
161 conditions in the numerical simulations (Figure 1). The properties of these soundings including

162 the parameters and indices are summarized in Table 1. The significance of these parameters and
 163 indices will be discussed and compared with those in the storms in subtropics and High Plains in
 164 Section 5. The case of Darwin, 23 January 2006 is used as an example to show the verification
 165 with observation and to demonstrate simulated cloud microphysical features including the
 166 kinematics, thermodynamics, radar reflectivity, and precipitation over the Pacific Warm Pool.

167 After the initiation of a thermal anomaly, the maximum updraft exhibits two peaks with
 168 large vertical velocity due to the spin-up process associated with the initialization. Then the
 169 storm adjusts to a quasi-steady state in the mature stage from 150 min to 240 min. During this
 170 period, the maximum vertical velocity remains in the range of 12.7-28.8 m s⁻¹ (Figure 2a). The
 171 simulated maximum updraft 25.9 m s⁻¹ compares well to both the 90th percentile maximum
 172 vertical velocity of dual-Doppler-retrieved 3-D wind fields (Varble et al. 2014) and a root mean
 173 square (RMS) bias in the dual-Doppler retrieval (Collis et al. 2013) (Figure 2b). The algorithm
 174 for radar reflectivity is the same as that used in Stoelinga (2005). The radar reflectivity around
 175 the surface in the simulation is also close to the observation (Figure 2c). At higher altitudes, the
 176 reflectivity is larger than the observation as in some previous modeling results (Varble et al.
 177 2014). Potential reasons such as size distribution assumption and slow hail mass-weighted fall
 178 speed were proposed to explain such larger upper radar reflectivity.

179 Figs. 3-4 show vertical cross-sections of the simulated hydrometeor mixing ratios,
 180 vertical velocity, and temperature. The mixing ratio of cloud water extends to at least 9 km
 181 altitude in the period from 150 min to 240 min. The corresponding temperature range is about
 182 245-250 K. Cloud base is around 3 km during this period. The greatest cloud ice mixing ratio is
 183 located near, but not lower than, the levels of the largest vertical velocity. The cloud ice can be
 184 as low as 9 km (at about 230 K) or as high as 16 km (at about 190 K) (Figs. 3-4).

185 Typically rain occurs below 6 km. The maximum rain mixing ratios occur mostly close
 186 to local strong downdraft, either higher or lower than the levels of local maximum downdrafts,
 187 but rarely below the maximum updrafts. The maximum rain mixing ratio at surface is
 188 approximately 5.5 g kg⁻¹. The altitudes of snow (9-16 km) are closely linked with cloud ice and
 189 its mixing ratios are modulated by wind fields. The corresponding temperature range is 250-190
 190 K. The wide spread of snow is due to several close convective cells and advection. Hail/graupel
 191 generally does not extend to as high as snow and cloud ice. At 180 min hail may reach a level as
 192 low as 2 km with strong downdraft. The maximum hail mixing ratio occurs at about 5-8 km
 193 height (Figs. 3-4).

194 During the quasi-steady mature stage, rain is distributed below 6 km, hail at about 3-13
 195 km, and cloud ice and snow at about 9-16 km. Supercooled water is common during the mature
 196 stage whereas hail is usually above the altitudes of freezing point. The hail mixing ratio contours
 197 overlap the rain mixing ratio contours, and even extend to rain cores, and vice versa. This implies
 198 that the accretion of rain is a major source of hail, meanwhile, the melting of hail is an important
 199 source of rain (Figs. 3b-d, 4b-c).

200

201 **4. Microphysical budget analysis of the case Darwin**

202 We used the storm occurred at Darwin on 23 January 2006 as an example to demonstrate
 203 the details of Microphysical budget over the Pacific Warm Pool. The portion of liquid phase in
 204 total condensate mass gradually stabilizes after experiencing two minima and subsequently
 205 dominates after 148 min (Figs. 5a-b). The change of the percentage of liquid phase with time is
 206 25%, 39%, 37%, 46%, 52%, 60%, 67%, and 61% of total condensate mass at 30, 60, 90, 120,
 207 150, 180, 210, and 240 min, respectively, by careful examination of the data in Figure 5b. On

208 average during quasi-steady stage over 150-240 min period with the data of every two mins, the
 209 portion of the liquid water phase in total condensate mass is about 62%. Individually, the
 210 percentages for cloud water, rain, cloud ice, snow, and hail in this period are about 25%, 37%,
 211 3%, 18%, and 18% respectively. The mass of liquid water phase and ice phase over this period
 212 are about 2503 and 1559 KT. The individual mass of cloud water, rain, cloud ice, snow, and hail
 213 are about 988, 1515, 116, 725, and 718 KT, respectively (Figure 5). In contrast, the contribution
 214 of the liquid phase is on an average 20% and 46% in the High Plains and subtropics (Table 2 in
 215 LWS05), respectively. Hence, the liquid-phase dominate the thunderstorm hydrometeors in the
 216 case of Darwin.

217 The time evolution of the sources and sinks of rain, snow, and hail/graupel will be
 218 discussed below.

219 *a. Rain*

220 Figs. 6a and 7a show the rain sources and their respective percentage contributions.
 221 Melting from hail (qhmlr) is the leading rain source most of the time. It becomes increasingly
 222 important from the start and reaches a maximum at 60 min (Figs. 6a and 7a). At this time,
 223 abundant hail melts underneath the altitude of 273 K level with downdraft at lower levels (Figs.
 224 3a and 4a). This follows the decrease of vertical velocity (Figure 2a) and peak hail generation at
 225 50 min (Figure 5a). The melting of hail (qhmlr) then decreases (Figs. 6a and 7a) while hail
 226 production becomes smaller (Figure 5a) due to the dissipation of original storm (Figure 2a). The
 227 melting of hail (qhmlr) increases again (Figs. 6a and 7a) with greater hail generation (Figs. 2a
 228 and 5a) after that. Melting of hail (qhmlr) becomes the most significant process of rain
 229 production during 42-78 min and after 88 min (Figs. 6a and 7a). The accretion of cloud water
 230 (qracw) rises very quickly (Figs. 6a and 7a) as new storm develops (Figure 2a), and becomes a
 231 major rain source briefly (Figs. 6a and 7a). The evolution of shedding from hail (qshsr) (Figs. 6a
 232 and 7a) is similar to maximum updraft (Figure 2a).

233 In general, accretion (qracw), shedding of hail (qshsr), and melting of hail (qhmlr) are the
 234 major rain sources subsequent to 18 min. Melting of hail (qhmlr), which contributes mostly
 235 more than 40% of the total rain creation, remains the largest rain source. Shedding of
 236 hail(qshsr), which contributes mostly more than 30%, is the second one (Figs. 6a and 7a).
 237 Accretion of cloud water (qracw), which contributes about 20%, is the third one after 104 min
 238 except it exceeds shedding (qshsr) in 210-212 min. Figs. 6b and 7b show rain sinks and
 239 percentage contributions. For the most part, evaporation of rain (qrcev) and accretion of rain by
 240 hail (qhacr) are the two main rain sinks. Since most of the hydrometeors are liquid in the
 241 beginning, evaporation of rain (qrcev) is the most significant rain sink initially. Soon after the
 242 ice-phase processes develop much more quickly than evaporation (qrcev) and thus dominates
 243 until 68 min. After 80 min the rain depletion rate due to the accretion of rain by hail (qhacr)
 244 increases dramatically just when bountiful hail is generated (Figure 5a). Accretion of rain by
 245 hail (qhacr) is consequential in a large downdraft. After coming to its peak, the production rate
 246 of accretion of rain by hail (qhacr) (Figs. 6b and 7b) is essentially parallel to the hail evolution
 247 (Figure 5a).

248 Evaporation (qrcev) and accretion by hail (qhacr) remain the two major sinks, and
 249 accretion by ice, which contributes less than 12%, is the third one. Accretion by hail (qhacr)
 250 provides at least 31% rain depletion after 100 min. Accretion by snow (qsacr) contributes less
 251 than 3% rain depletion following 24 min, and thus can be neglected. After about 104 min
 252 evaporation (qrcev), accretion by hail (qhacr), and accretion by ice (qiacr) reach a quasi-steady
 253 state and average 50%, 42%, and 7% of the total sink rate, respectively (Figs. 6b, and 7b).

254 *b. Snow*

255 The dominant process of snow production is the Bergeron-Findeisen process, the
 256 transformation of cloud ice to snow (qsfi), which is larger than other processes almost all the
 257 time (Figs. 2a, and 6c). The Bergeron-Findeisen process (qsfi) occurs due to the difference of
 258 saturation vapor pressure over supercooled water and ice (see Wang, 2013) is more efficient than
 259 any other processes in snow production. The fact that the trend of maximum updraft is nearly
 260 parallel to the evolution of the Bergeron-Findeisen process (qsfi) suggests that the updraft moves
 261 cloud water upward to form supercooled water at middle and upper levels (Figs. 2a, and 6c).

262 The time evolution of snow production behaves as follows. Water vapor deposition to
 263 snow (qsdpv) occurs first. The transfer from cloud ice to snow in Bergeron-Findeisen process
 264 (qsfi) contributes no less than 41% of snow production subsequent to 16 min when it becomes
 265 the most significant snow source in simulation except at 12 min when vapor sublimation from
 266 snow (qssbv) surpasses it briefly. Accretion of cloud water (qsacw) continues to be the second
 267 largest source since 14 min. The production rate of deposition (qsdpv) becomes the third place
 268 after 34 min except for the contribution from the accretion of rain by ice (qiacr) overshadows
 269 deposition (qsdpv) briefly during 80-90 min. After 156 min the Bergeron-Findeisen process
 270 (qsfi), accretion of cloud water (qsacw), and deposition (qsdpv) approach a quasi-steady state,
 271 and average 47%, 34%, and 12% of total snow production rate, respectively (Figs. 6c, and 7c).

272 Accretion of snow by hail (qhacs) is the largest snow sink except for the first 10 min
 273 (Figure 6d). The peak of snow mass (at about 85 min) leading the peak of accretion by hail
 274 (qhacs) (at about 110 min) indicates the importance of snow depletion due to the accretion of
 275 snow aggregates by hail (qhacs) (Figs. 5a, and 6d). The time evolution of the sublimation of
 276 snow (qssbv) is similar to that of the accretion by hail (qhacs) except that the peak of
 277 sublimation (qssbv) occurs at about 115 min (Figure 6d).

278 Sublimation (qssbv) becomes and remains the second largest sink after 26 min. The
 279 depletion rate due to the melting of snow to rain (qsmr) is similar to the accretion by rain
 280 (qhacs) but becomes larger after 144 minute. Accretion by hail (qhacs), sublimation (qssbv), and
 281 melting (qsmr) come to a quasi-steady state subsequent to 156 min, and average about 77%,
 282 14%, and 7% of total snow sink rate, respectively (Figs. 6d, and 7d).

283 *c. Hail/graupel*

284 We will use hail to represent this category of hydrometeor in the following discussion.
 285 Accretion of rain by hail (qhacr), accretion of cloud water by hail (qhacw), and accretion of snow
 286 aggregates by hail (qhacs) are the top three production processes of hail after 40 min. Ample hail
 287 is produced in the initial spin-up period. Afterwards, the hail mass evolution is approximately
 288 parallel to the accretion of rain (qhacr). It is observed that the curve of accretion of rain (qhacr)
 289 lags hail mass progression. The hail production can be summarized as follows. After 108 min
 290 the accretion of rain (qhacr) becomes the leading mechanism except for the accretion of snow
 291 (qhacs) takes over briefly during 148-152 min. Accretion of rain (qhacr), cloud water (qhacw),
 292 and snow (qhacs) are the major processes and have similar production rates from 90 min on.
 293 These three processes become quasi-steady subsequent to 112 min, and average 37%, 30%, and
 294 26% of total hail/graupel production rate, respectively (Figs. 6e, and 7e).

295 Rain shed from hail (qhshr) and melting of hail (qhmlr) to form rain are the two prime
 296 hail/graupel sinks all the time (Figs. 6f, and 7f). This is consistent with the previous observation
 297 that the shedding (qhshr) and melting (qhmlr) are mainly the utmost rain sources. Rain mass is
 298 the most plentiful of all hydrometeors during 62-80 min and after 114 min. Shedding (qhshr)

299 and melting (qhmlr) reach the quasi-steady state after 112 min, and average 58%, and 41%, of
 300 total hail sink rate, respectively (Figs. 6f, and 7f).

301 A flow chart of the major microphysical processes is shown in Fig. 8. See the appendix
 302 for the microphysical meanings of the acronyms.

303

304 **5. Comparison of microphysical budgets of storms over the Pacific Warm Pool with that** 305 **in the subtropics and High Plains in LWS05**

306 To comprehend the specialties in the microphysical processes over the Pacific Warm
 307 Pool, we will compare the microphysical budgets of precipitates over the Pacific Warm Pool with
 308 those in the subtropics and High Plains in LWS05 during the quasi-steady stage in this section.

309 The altitude for the melting of hails to start after they fall below the freezing level can be
 310 approximated by wet-bulb zero height. The wet-bulb zero height represents the minimum depth
 311 that a hail would fall without melting appreciably (Miller and McGinley 1977). The wet-bulb
 312 zero height over the Pacific Warm Pool is on the whole higher than in the subtropics and much
 313 higher than in the High Plains in LWS05. This shows that Pacific Warm Pool storms have lower
 314 potential to produce hail. Even existing hail would have melted in thick and warm low-level
 315 clouds before reaching the ground. The sub-cloud average mixing ratio of water vapor over the
 316 Pacific Warm Pool, in general, is more abundant in spite of lower surface temperatures than in
 317 the subtropics and much more abundant than in the High Plains in LWS05. The Warm Pool
 318 storms have lower and warmer cloud bases, implying that liquid-phase hydrometeors are more
 319 crucial. The lifting condensation levels (LCLs) and levels of free convection (LFCs) of the
 320 soundings over the Pacific Warm Pool are lower, but the equilibrium levels (ELs) are higher than
 321 in the subtropical and High Plains cases in LWS05. Therefore, the convective available potential
 322 energies (CAPEs) over the Pacific Warm Pool are larger than in the subtropics and High Plains in
 323 LWS05. In short, the characteristics of storm environment over the Pacific Warm Pool are close
 324 to those in the subtropics in LWS05, but they have more abundant moisture and are more prone
 325 to deep convection (see Table 1 and also Table 1 in LWS05).

326 On average the percentage of precipitation particles (rain, snow, and graupel/hail) in the
 327 total hydrometeor mass in storms over the Pacific Warm Pool (77.7%) is much less than in the
 328 subtropics (84.3%) and High Plains (84.6%) in LWS05. Even the largest percentage over the
 329 Pacific Warm Pool is less than the smallest one in the subtropics and High Plains in LWS05. In
 330 detail, the percentage of cloud water over the Pacific Warm Pool is much larger, but the
 331 percentage of hail is smaller than in the subtropics and High Plains in LWS05 due to higher
 332 temperature. Ice-phase hydrometeors (cloud ice, snow, and graupel/hail) over the Pacific Warm
 333 Pool (about 41%) are smaller than in the subtropics (about 54%) and High Plains (about 80%) in
 334 LWS05 (Table 2 and Table 2 in LWS05). In a nutshell, liquid-phase hydrometeors dominate in
 335 thunderstorms development over the Pacific Warm Pool.

336 To study how the hydrometeor partitioning would be sensitive to the bubble parameters
 337 over the Pacific Warm Pool, we conducted sensitivity tests to understand the effects of bubble
 338 width, height, position, and temperature perturbation. In the case of Singapore, the ratio of
 339 liquid-phase hydrometeors is about 61%. The thermal bubble is 2 km above ground level, with
 340 the maximum 3.5 K at bubble center and the radius 10 km and depth 4 km. When the thermal
 341 bubble is 4 km above ground, the ratio of liquid-phase hydrometeors is about 66%. The ratios
 342 for the maximum 2 K and 5 K cases are about 62% and 59%, respectively. The ratios for the
 343 depth 2 km and 6 km cases are about 56% and 60%, correspondingly. The ratios for the radius 5
 344 and 20 km is 69% and 56%, respectively. In short, the sensitivity tests demonstrate that the

345 dominance of liquid-phase hydrometeors in thunderstorm development over the Pacific Warm
 346 Pool is robust.

347 *a. Rain*

348 Melting of hail to rain (qhmlr) and accretion of cloud water by rain (qracw) are the two
 349 most significant rain sources over the Pacific Warm Pool (Table 3). The temperature over the
 350 Pacific Warm Pool is warmer and the updraft at lower levels is weaker than in the subtropics and
 351 High Plains in LWS05. Typically hail contours strongly overlap that of rain. The contours of 0.5
 352 g kg^{-1} hail/graupel extend below the levels of 285 K and are close to the rain cores (Figs. 3-4).
 353 These indicate that melting of hail (qhmlr) is a very important source of rain, and more
 354 significant than shedding from hail (qhshr) when hail falls that is the difference in vertical cross-
 355 sections of hydrometeors between over the Pacific Warm Pool and in the subtropics (Figs. 4-5 in
 356 LWS05). Hence, melting of hail (qhmlr) over the Pacific Warm Pool (about 40%) on average
 357 contributes more than in the subtropics (about 30%); whereas the contribution of shedding from
 358 hail (qhshr) over the Pacific Warm Pool (about 31%) is between in the subtropics (about 45%)
 359 and High Plains (about 20%) in LWS05 (Figs. 8a-b, and Figs. 16a-b in LWS05). Although
 360 melting is more important than shedding in the High Plains, the conditions of hail separating
 361 from rain, no strong wet growth, and few raindrops being collected (LWS05) are very different
 362 from the Pacific Warm Pool. Because of high temperature and large humidity in tropical Pacific
 363 Warm Pool, cloud water tends to develop instead of freezing, and thus, has more chance to
 364 undergo collision and coalescence to form rain. Therefore, accretion of cloud water by rain
 365 (qracw) over the Pacific Warm Pool (about 26%) contributes more than in the subtropics (about
 366 20%) and High Plains (about 5%) (Figure 9c and Figure 16c in LWS05).

367 Accretion of rain by hail (qhacr) and evaporation of rain (qrcev) are the two dominant
 368 rain sinks over the Pacific Warm Pool (Table 3). These two processes contribute more than 88%
 369 of rain depletion in all three cases. Generally, rain stays at low levels. The humidity at lower
 370 level determines the dominant sinks (Figure 1 and Table 3). Because of larger updraft and more
 371 significant overlap of hail and rain regions in the subtropical case CaPE (0729) in LWS05 (Figs.
 372 3-5 in LWS05), accretion by hail (qhacr) is more important in rain depletion in the subtropics
 373 (about 65%) than over the Pacific Warm Pool (about 50%). In addition to accretion by hail
 374 (qhacr), the updraft in a quasi-steady state of CaPE (0729) in LWS05 (Figure 3 in LWS05) is
 375 larger than over the Pacific Warm Pool (Figure 2a). Therefore, more rain is brought aloft in the
 376 subtropics in LWS05 that leads to the larger accretion of supercooled water by ice and snow
 377 (qiacr, and qsacr) than over the Pacific Warm Pool. The accretions of supercooled water by ice
 378 and snow (qiacr, and qsacr) in the subtropics in LWS05 account for more rain depletion than
 379 over the Pacific Warm Pool (Figs. 9c-d and Figs. 17c-d in LWS05). Conversely, the contribution
 380 of evaporation (qrcev) in the subtropics in LWS05 (about 20%) is less than over the Pacific
 381 Warm Pool (about 40%) on average (Figure 10b and Figure 17b in LWS05).

382 *b. Snow*

383 The ranking of snow sources over the Pacific Warm Pool is similar to the subtropics
 384 (Table 4 and Table 4 in LWS05). Bergeron-Findeisen process of cloud ice to snow (qsfi) is the
 385 chief process and is on average more efficient in the subtropics in LWS05 (about 70%) than over
 386 the Pacific Warm Pool (about 57%) (Figure 11a, and Figure 18a in LWS05). Due to larger snow
 387 concentration and lower cloud base over the Pacific Warm Pool (Figs. 3-4 and Figs. 4-5 in
 388 LWS05) accretion of cloud water (qsacw) is more dominant over the Pacific Warm Pool (about
 389 25%) than in the subtropics in LWS05 (about 15%) (Figure 11b, and Figure 18b in LWS05).
 390 Both Bergeron-Findeisen process of cloud ice to snow (qsfi) and accretion of cloud water

391 (qsacw) generally account for about 83% of snow production during the quasi-steady stage over
 392 the Pacific Warm Pool. Because of higher levels of snow over the Pacific Warm Pool (Figs. 3-4
 393 and Figs. 4-5 in LWS05) less water vapor accompanied by larger snow concentration gives rise
 394 to more essential vapor deposition to snow (qsdpv) (about 9%) than in the subtropics in LWS05
 395 (about 4%). Accretion of cloud ice (qsaci) and accretion of rain by ice (qiacr) play about equal
 396 role over the Pacific Warm Pool and in the subtropics in LWS05 (Figure 11c, and Figs. 11c, and
 397 18c in LWS05).

398 The ranking of snow sinks over the Pacific Warm Pool is close to that in the subtropics
 399 and High Plains in LWS05 (Table 4 and Table 4 in LWS05). The contribution of about 80% of
 400 snow depletion due to accretion by hail (qhacs) over the Pacific Warm Pool is similar to that in
 401 the subtropics and High Plains in LWS05 as well (Figure 12a, and Figure 19a in LWS05). Vapor
 402 sublimation of snow (qssbv) accounts for 15% of snow depletion over the Pacific Warm Pool,
 403 roughly the same as those in the subtropics and High Plains in LWS05 (Figure 12b, and Figure
 404 19b in LWS05). The significance of accretion of snow by rain to form snow or hail (qracs) over
 405 the Pacific Warm Pool (about 3%) is also similar to those in the subtropics in LWS05 (Figure
 406 12c, and Figure 19c in LWS05).

407 *c. Hail/graupel*

408 The order of ranking of hail sources over the Pacific Warm Pool is similar to the
 409 subtropics in LWS05 (Table 5 and Table 5 in LWS05), but the percentage contributions of the
 410 major three hail sources, accretion of rain (qhacr), accretion of cloud water (qhacw), and
 411 accretion of snow (qhacs), over the Pacific Warm Pool are a bit different from in the subtropics
 412 in LWS05. Larger updraft in the subtropics (Figure 2a, and Figure 3 in LWS05) enables more
 413 accretion of supercooled water (about 50%) than over the Pacific Warm Pool (about 40%)
 414 (Figure 13a, and Figure 20a in LWS05). The accretion of rain by hail (qhacr) is more efficient
 415 than the accretion of cloud water by hail (qhacw) (about 30%) (Figure 13b, and Figure 20b in
 416 LWS05). Larger snow concentration over the Pacific Warm Pool results in more accretion of
 417 snow (qhacs) (about 22%) than in the subtropics in LWS05 (about 10%) (Figure 13c, and Figure
 418 20c in LWS05). Due to larger updraft the accretions of supercooled by ice (qiacr) and snow
 419 (qsacr) to form snow or hail in the subtropics in LWS05 (about 7% and 3%, respectively) are
 420 larger than over the Pacific Warm Pool (about 5% and 1%, respectively) (Figs. 12d-e, and Figs.
 421 20d-e in LWS05).

422 The ranking of hail sinks over the Pacific Warm Pool is the same as the High Plains in
 423 LWS05 (Table 5 and Table 5 in LWS05). Both melting of hail (qhmlr) and shedding from hail
 424 (qhshr) contribute about 99% of hail depletion over the Pacific Warm Pool (Figure 14). Melting
 425 of hail (qhmlr) (about 55%) is more important than shedding from hail (qhshr) (about 44%) over
 426 the Pacific Warm Pool that is very different from the subtropics where shedding from hail
 427 (qhshr) (about 60%) is more important than melting of hail (qhmlr) (about 40%) (Figs. 13a-b,
 428 and Figs. 21a-b in LWS05). In the High Plains in the LWS05 melting of hail (qhmlr) (about
 429 75%) is much more efficient than shedding from hail (qhshr) (about 20%). The microphysical
 430 conditions in High Plains that hail region separates from rain, no strong wet growth happens, and
 431 little rain is collected as hail falls (LWS05) are completely different from the Pacific Warm Pool.
 432 Due to higher temperature and moisture, water vapor sublimation from hail (qhsbv) over the
 433 Pacific Warm Pool (0.8%) is less efficient than in the subtropics (1%) and High Plains (3%) in
 434 LWS05.

435

436 **6. Summary**

437 In this study, we have employed a microphysical model WISCDYMM to investigate the
438 dynamics, thermodynamics, the details of hydrometeors, and the development microphysical
439 budgets during quasi-steady stage over the Pacific Warm Pool. In addition, we have examined
440 the differences in hydrometeor partitioning and cloud microphysical processes between the
441 Pacific Warm Pool, the subtropics, and High Plains in LWS05.

442 During the mature stage, supercooled water is common. The hail mixing ratio contours
443 overlap significantly with that of rain. Furthermore, the hail mixing ratio contours get as far as
444 rain cores and the rain mixing ratio contours extend to hail cores, which shows that the melting
445 of hail is an important source of rain and that the accretion of rain is a major source of hail.

446 Microphysical budgets depend on atmospheric dynamical and thermodynamical
447 conditions and give the partitioning information of hydrometeors. Liquid-phase hydrometeors
448 govern thunderstorm development over the Pacific Warm Pool. The ratio of ice-phase mass to
449 liquid-phase mass is about 41% : 59% (Table 2) over the Pacific Warm Pool, compared with
450 54% : 46% in the subtropics and 80% : 20% in the High Plains (Table 2 in LWS05). The
451 sensitivity tests on bubble width, height, position, and temperature perturbation indicate that the
452 domination of liquid-phase hydrometeors in thunderstorm development over the Pacific Warm
453 Pool is potent.

454 The principal microphysical processes over the Pacific Warm Pool are as follows. The
455 main rain sources are also the key hail sinks, namely, melting of hail and shedding from hail;
456 whereas the crucial rain sinks are accretion by hail and evaporation. The major snow sources are
457 Bergeron-Findeisen process, transfer of cloud ice to snow, and accretion of cloud water; whereas
458 the foremost snow sink is accretion by hail. The essential hail sources are accretions of rain,
459 cloud water, and snow; whereas the most important hail sinks are melting of hail and shedding
460 from hail. The contribution and ranking of sources and sinks of these precipitates have been
461 compared with those in the subtropics and High Plains in LWS05.

462 Microphysical budgets are closely associated with cloud radiative properties. Cloud-
463 radiation interactions would lead to the changes in climate. Therefore, a clear understanding of
464 such microphysical processes is not only important in estimating the performance of cloud
465 models and products of passive satellite measurements but also helps to develop a proper
466 parameterization of moist processes in climate models.

467 The sounding data decide atmospheric dynamical and thermodynamical conditions, and
468 thus the hydrometeor partitioning and cloud microphysical processes. Microphysical processes
469 are highly nonlinear in themselves. A comprehensive and thorough study is necessary to
470 understand clear connections between them. This knowledge will ultimately benefit the
471 microphysics parameterization and cumulus parameterization in global circulation and climate
472 models.

473

474 Acknowledgments

475 This study is partially supported by NOAA GOES-R contract 9636, Risk
476 Reduction Award NA10NES4400013, National Science Foundation Grant AGS-1219586
477 and AGS- 1633921, and research fund support from Academia Sinica and Ministry of Science
478 and Technology, Taiwan. Any opinions, findings and conclusions or recommendations expressed
479 in this material are those of the authors and do not necessarily reflect the views of the National
480 Science Foundation (NSF). The authors also thank the helpful comments of two anonymous
481 reviewers that led to improvements of the paper.
482

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Table Captions

- 559 **Table 1.** Characteristics and indices of the observed soundings over the Pacific Warm Pool
560 **Table 2.** Time-averaged partitioning of the domain-integrated hydrometeor masses during the
561 quasi-steady stage of each simulated storm
562 **Table 3.** Rankings of the domain-integrated individual sources and sinks of rain
563 **Table 4.** As in Table 3, but for snow
564 **Table 5.** As in Table 3, but for hail

565 **Table 1**
 566 Characteristics and indices of the observed soundings over the Pacific Warm Pool

		Pacific Warm Pool		
		Darwin	Singapore	Medan
		1/23/2006	3/ 28/2013	7/18/2013
Sub-cloud average mixing ratio (g/kg)		20	18	17
Surface	P (m b)	998	995	996
	T (° C)	28.1	27.5	27.4
LCL	P (m b)	970.34	964.23	937.29
	T (° C)	25.66	24.07	21.88
LFC	P (m b)	888.70	922.53	833.51
	T (° C)	22.55	22.44	17.49
EL	P (m b)	100.73	102.81	119.45
	T (° C)	-82.35	-84.82	-79.80
CAPE (J/kg)		3941	3725	3667
CIN (J/kg)		6	2	13
Total totals index		47	46	48
Lift index		-5	-5	-7
Showalter index		-2.9	-3.5	-2.1
K index		40	41	41
Deep convective index		44.5	42.7	39.1
Severe weather threat index		343	257	246
Wet bulb zero _m (AGL)		5130	4824	4149

567

568 **Table 2**
 569 Time-averaged partitioning of the domain-integrated hydrometeor masses during the quasi-
 570 steady stage of each simulated storm

	Pacific Warm Pool (%)		
	Darwin	Singapore	Medan
Cloud water	25	16	15
Rain	37	45	39
Cloud ice	3	4	4
Snow	18	15	21
Hail	17	20	21
Precipitates	72	80	81
Non- precipitates	28	20	19
Snow + Hail	35	35	42
Ice phases	38	39	46
Liquid phases	62	61	54

571

572 **Table 3**

573 Rankings of the domain-integrated individual sources and sinks of rain

		Pacific Warm Pool		
		Darwin	Singapore	Medan
Production	qhshr	2	3	2
	qhmlr	1	2	1
	qracw	3	1	3
	qrcnw	4	4	4
	qsmr	5	5	5
	qsacw	6	6	6
Depletion	qhacr	2	1	1
	qrcev	1	2	2
	qiacr	3	3	3
	qsacr	4	4	4
	qrfrz	5	5	5

574

575 **Table 4**
 576 As in Table 3, but for snow

		Pacific Warm Pool		
		Darwin	Singapore	Medan
Production	qsfi	1	1	1
	qsacw	2	2	2
	qsdpv	3	4	3
	qsaci	5	5	5
	qiacr	4	3	4
	qsacr	6	6	6
	qseni	7	7	7
	qraci	9	9	9
	qrfrz	8	8	8
	qsfw	10	10	10
Depletion	qhacs	1	1	1
	qssbv	2	2	2
	qracs	4	3	3
	qsmr	3	4	4
	qhens	5	5	5

577

578 **Table 5**
 579 As in Table 3, but for hail

		Pacific Warm Pool		
		Darwin	Singapore	Medan
Production	qhacr	1	1	1
	qhacw	2	2	2
	qhacs	3	3	3
	qiacr	4	4	4
	qsacr	5	5	5
	qhdpv	6	6	6
	qhaci	8	7	7
	qracs	7	8	8
	qhens	9	9	9
	qrfrz	10	10	10
	qraci	11	11	11
Depletion	qhshr	2	2	2
	qhmlr	1	1	1
	qhsbv	3	3	3

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Figure Captions

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Figure 1. Skew T – log P diagram at (a) Darwin, Northern Territory, Australia at 1200 UTC on 23 January 2006; (b) Singapore at 0000 UTC on 28 March 2013; (c) Medan, Indonesia at 1200 UTC on 18 July 2013.

Figure 2. (a) Time evolution of the maximum vertical velocity (m s^{-1}), (b) the 50th percentile (blue), and 90th percentile (red) of maximum vertical velocity (m s^{-1}), and (c) the 50th percentile (blue), 90th percentile (red), and 99th percentile (green) maximum radar reflectivity (dBZ) in the simulation for the case of Darwin, 23 January 2006.

Figure 3. Simulated hydrometeor mixing ratios, vertical velocity (shaded), and temperature (gray) for the case of Darwin on 23 January 2006 in X-Z vertical cross-sections through the maximum updraft as of (a) 150 min, (b) 180 min, (c) 210 min, and (d) 240 min. The minimum values of cloud water (green), rain (pink), cloud ice (cyan), snow (blue) and graupel/hail (black) are 1, 0.5, 0.5, 0.5, and 0.5 g kg^{-1} , respectively. The intervals of these hydrometeors are 0.5, 0.5, 0.2, 0.2, and 0.5 g kg^{-1} , respectively.

Figure 4. As in Figure 3 but in Y-Z vertical cross-sections.

Figure 5. Time evolution of simulated (a) total condensate mass and (b) percentage contribution, integrated over the entire simulation domain ($64 \times 64 \times 18.4 \text{ km}^3$), for cloud water (blue), rain (red), cloud ice (green), snow (purple), and hail/graupel (cyan) in the case of Darwin, 23 January 2006.

Figure 6. Sources and sinks of each precipitating hydrometeor class for the Darwin, 23 January 2006 storm: (a) rain sources, (b) rain sinks, (c) snow sources, (d) snow sinks, (e) hail/graupel sources, (f) hail/graupel sinks. See the appendix for the microphysical meanings of the acronyms.

Figure 7. As in Figure 6 but that the precipitating hydrometeor sources and sinks are plotted by percentage contributions.

Figure 8. A flow chart of all microphysical processes. Red capital letters are the major ones. See the appendix for the microphysical meanings of the acronyms.

Figure 9. Superimposed time series in all three storm simulations for the rain sources: (a) qhshr, (b) qhmlr, (c) qracw. See the appendix for the microphysical meanings of the acronyms.

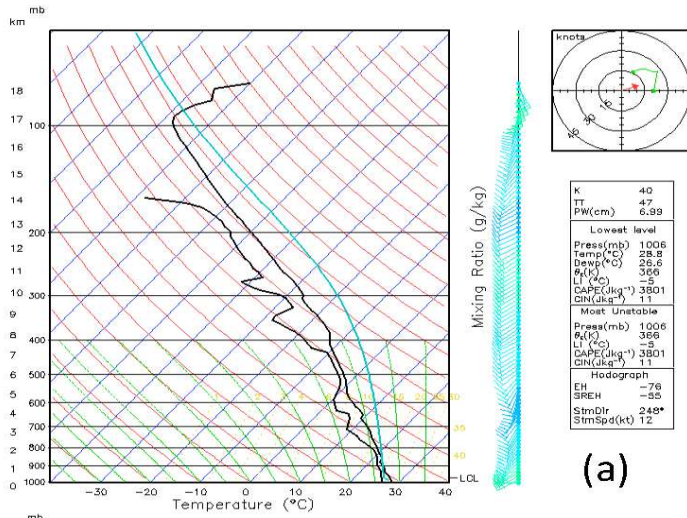
Figure 10. Superimposed time series in all three storm simulations for the rain sinks: (a) qhacr, (b) qrcev, (c) qiacr, (d) qsacr, (e) qrfrz. See the appendix for the microphysical meanings of the acronyms.

Figure 11. Superimposed time series in all three storm simulations for the snow sources: (a) qsf, (b) qsacw, (c) qsdpv, (d) qsaci, (e) qiacr. See the appendix for the microphysical meanings of the acronyms.

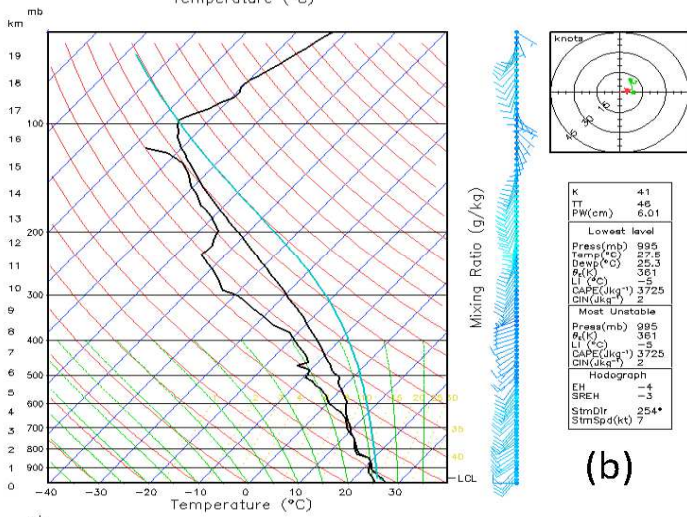
Figure 12. Superimposed time series in all three storm simulations for the snow sinks : (a) qhacs, (b) qssbv, (c) qracs. See the appendix for the microphysical meanings of the acronyms.

Figure 13. Superimposed time series in all three storm simulations for the hail/graupel sources: (a) qhacr, (b) qhacw, (c) qhacs, (d) qiacr, (e) qsacr. See the appendix for the microphysical meanings of the acronyms.

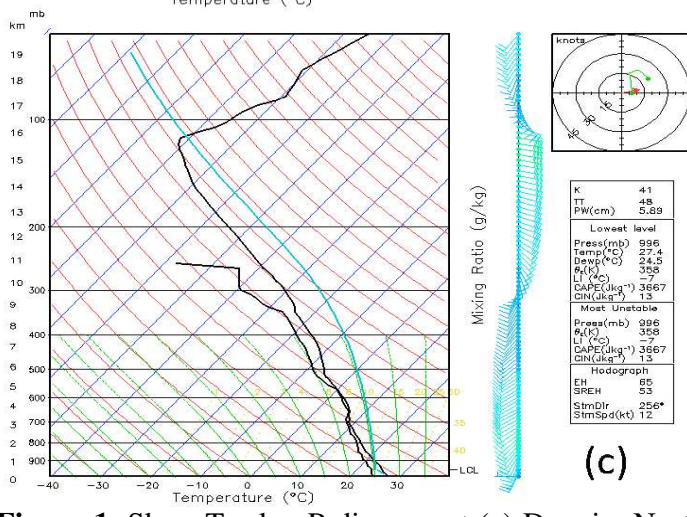
Figure 14. Superimposed time series in all three storm simulations for the hail/graupel sinks: (a) qhshr, (b) qhmlr, (c) qhsbv. See the appendix for the microphysical meanings of the acronyms.



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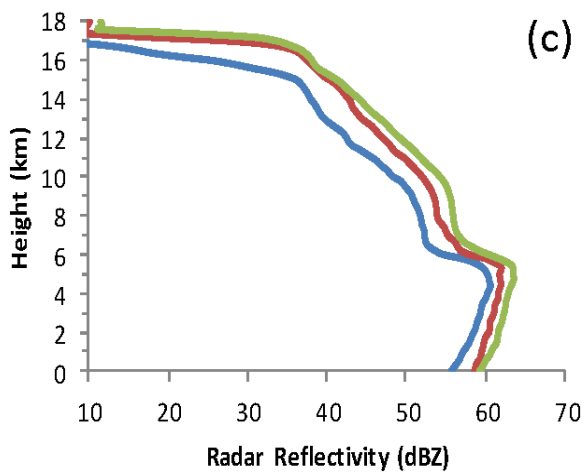
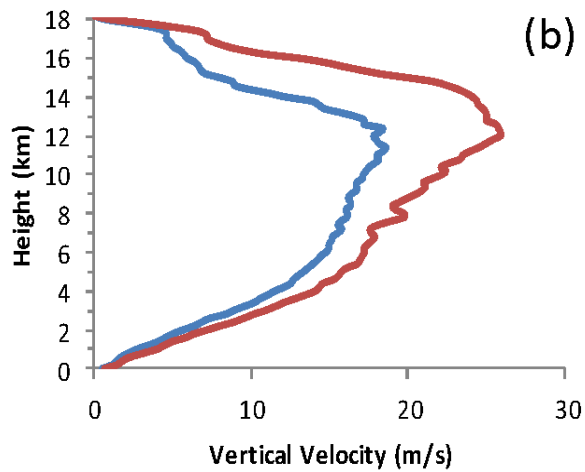
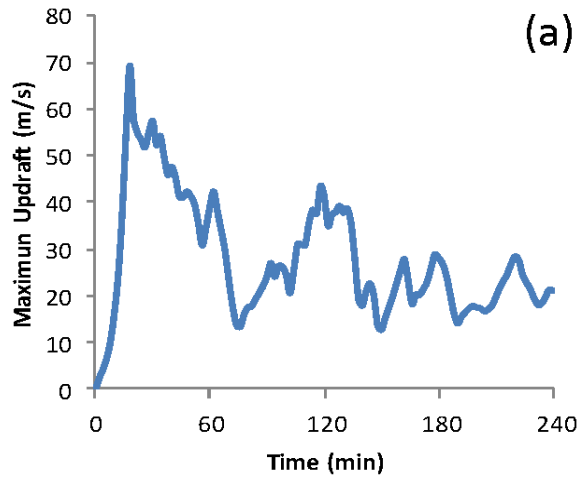


(b)



(c)

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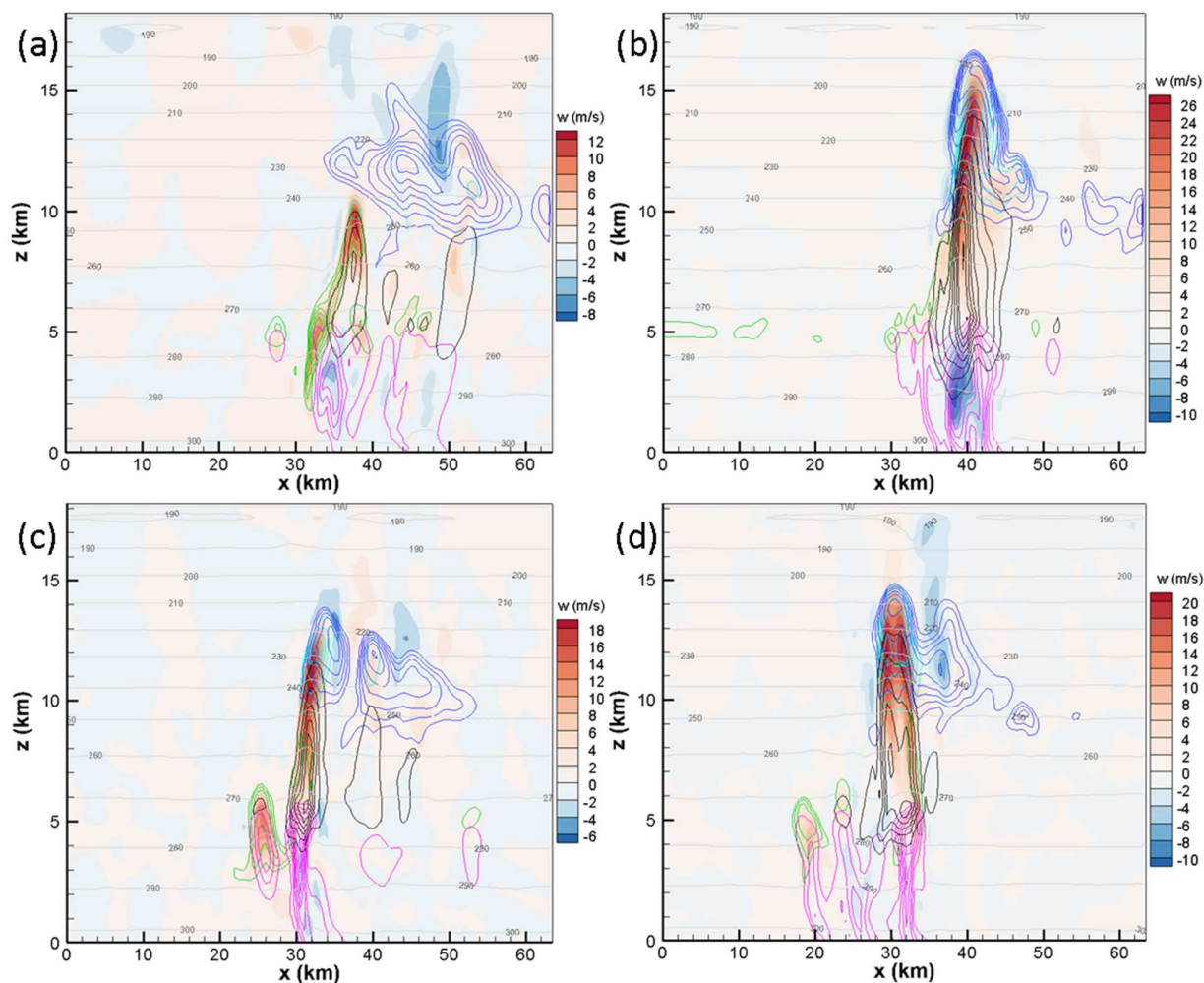
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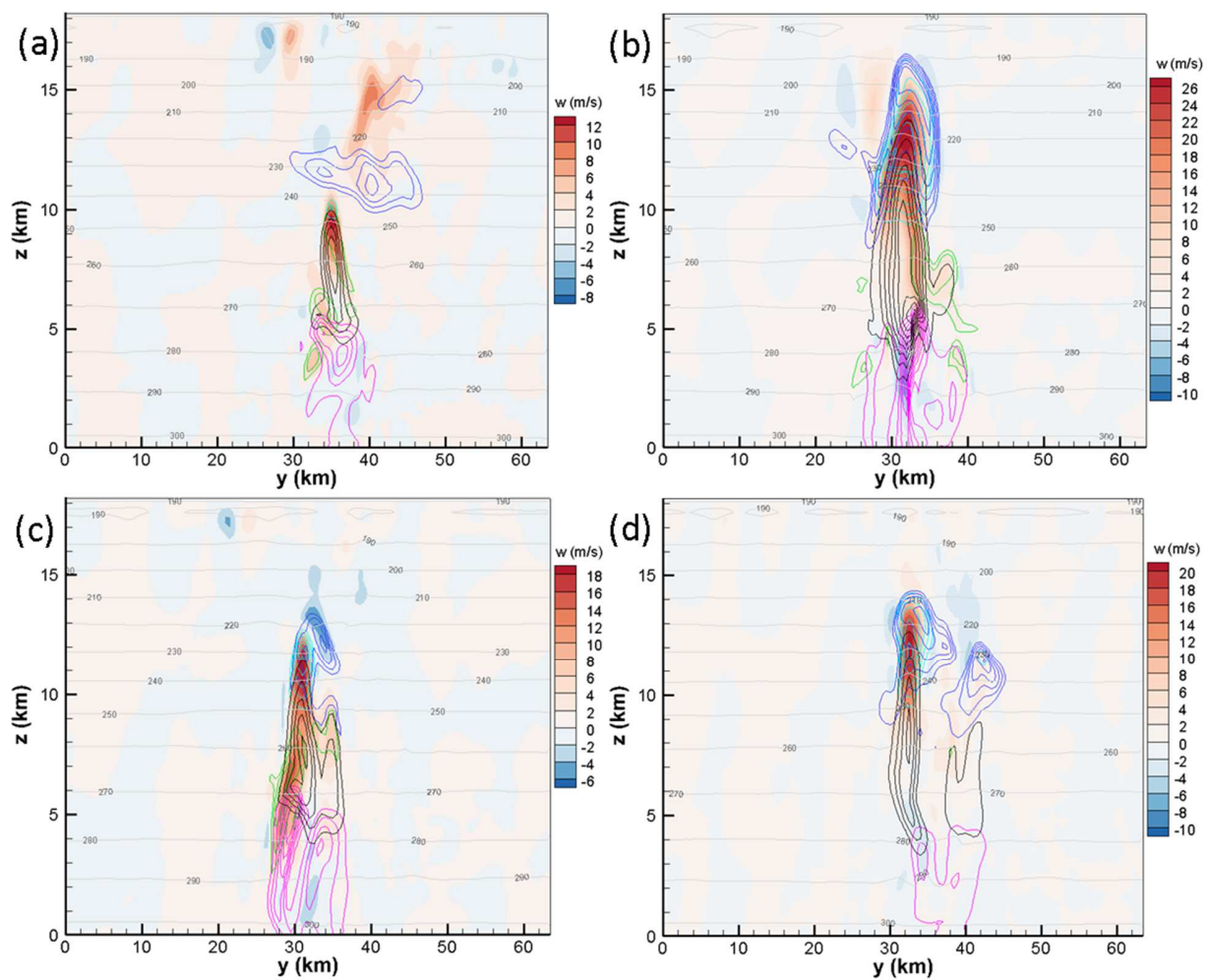
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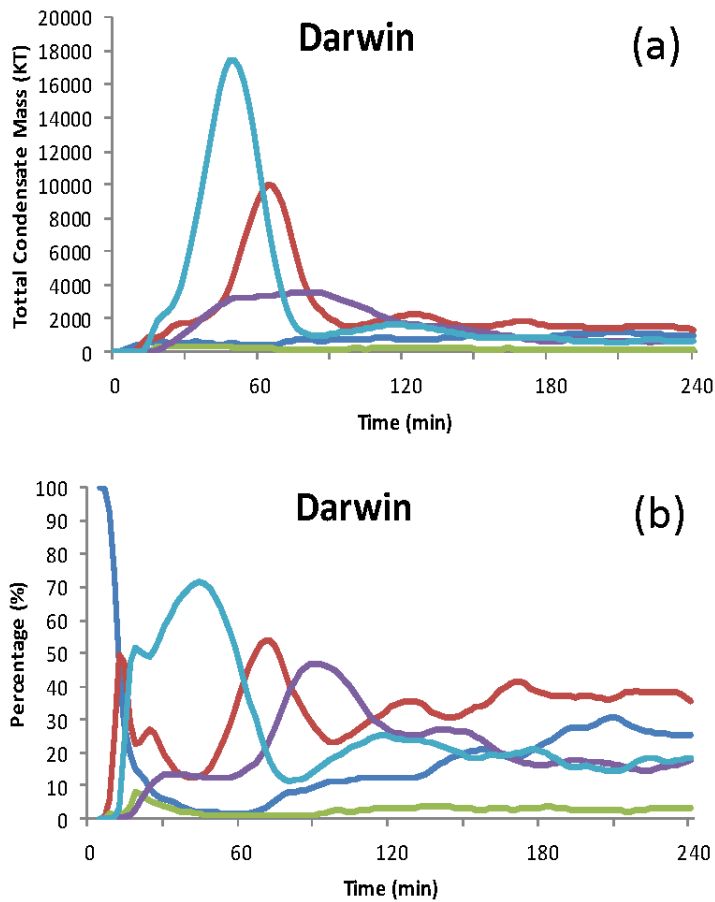


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 642 0.2, 0.2, and 0.5 g kg^{-1} , respectively.
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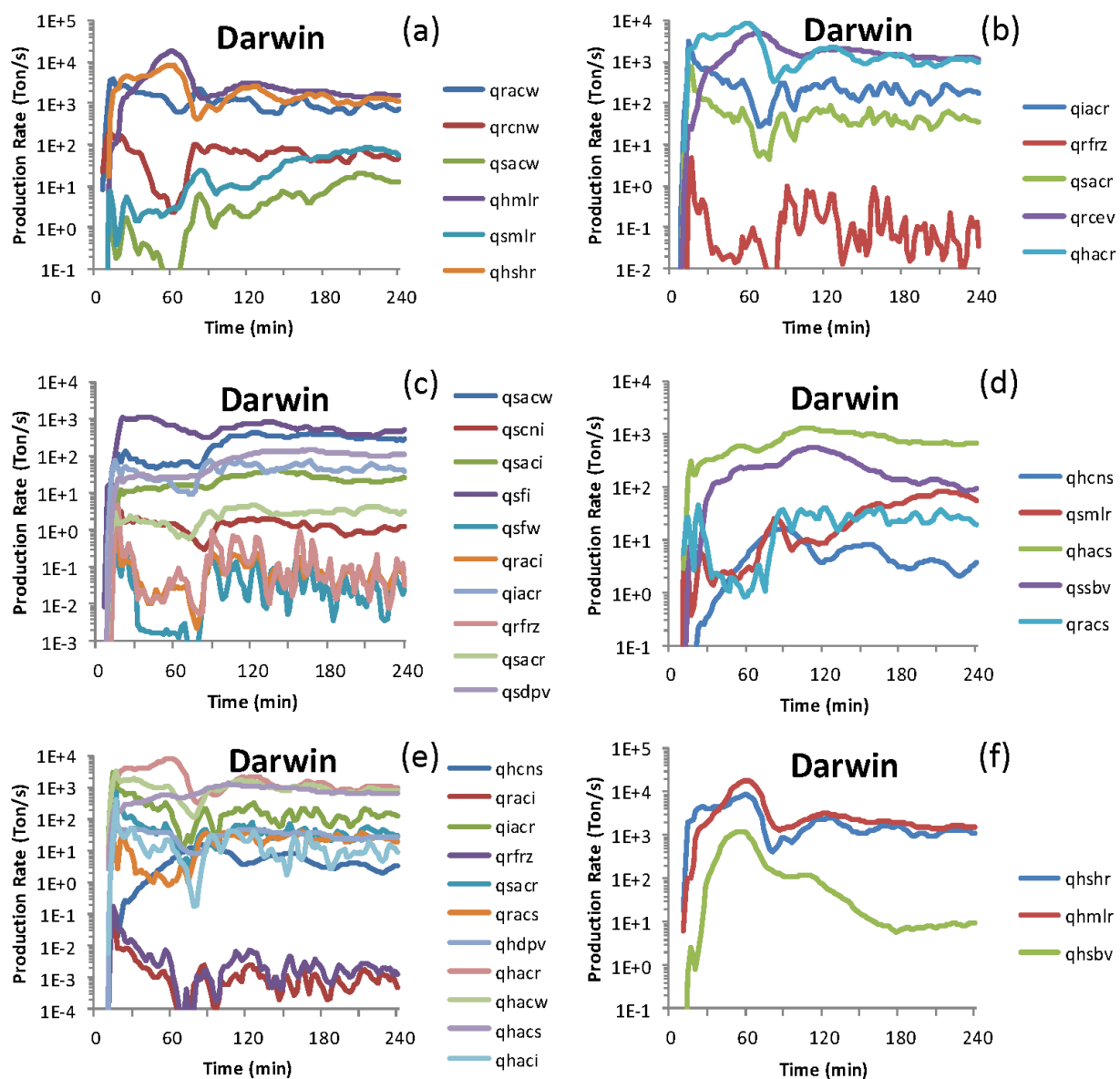


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Figure 4. As in Figure 3 but in Y-Z vertical cross-sections.

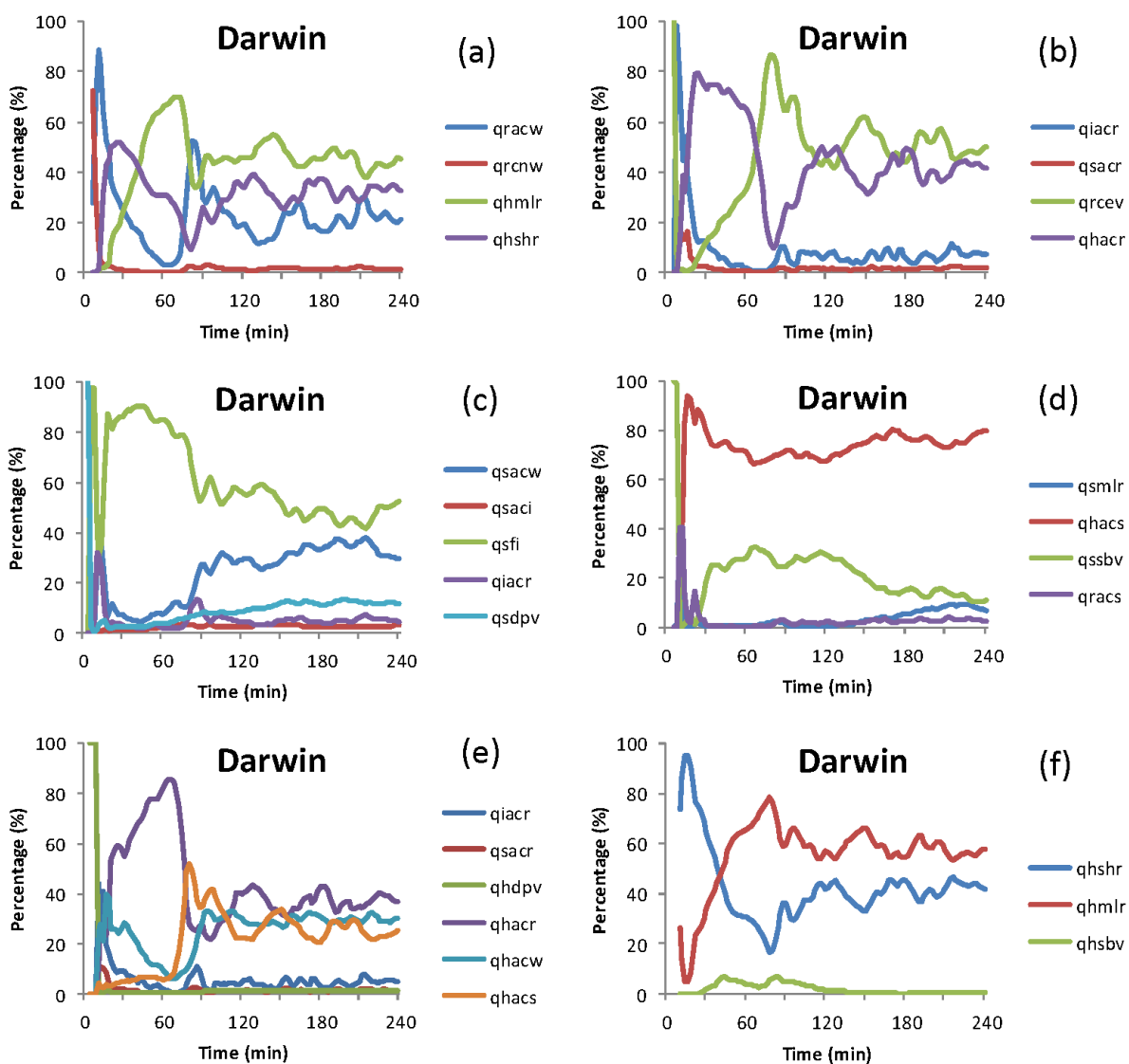


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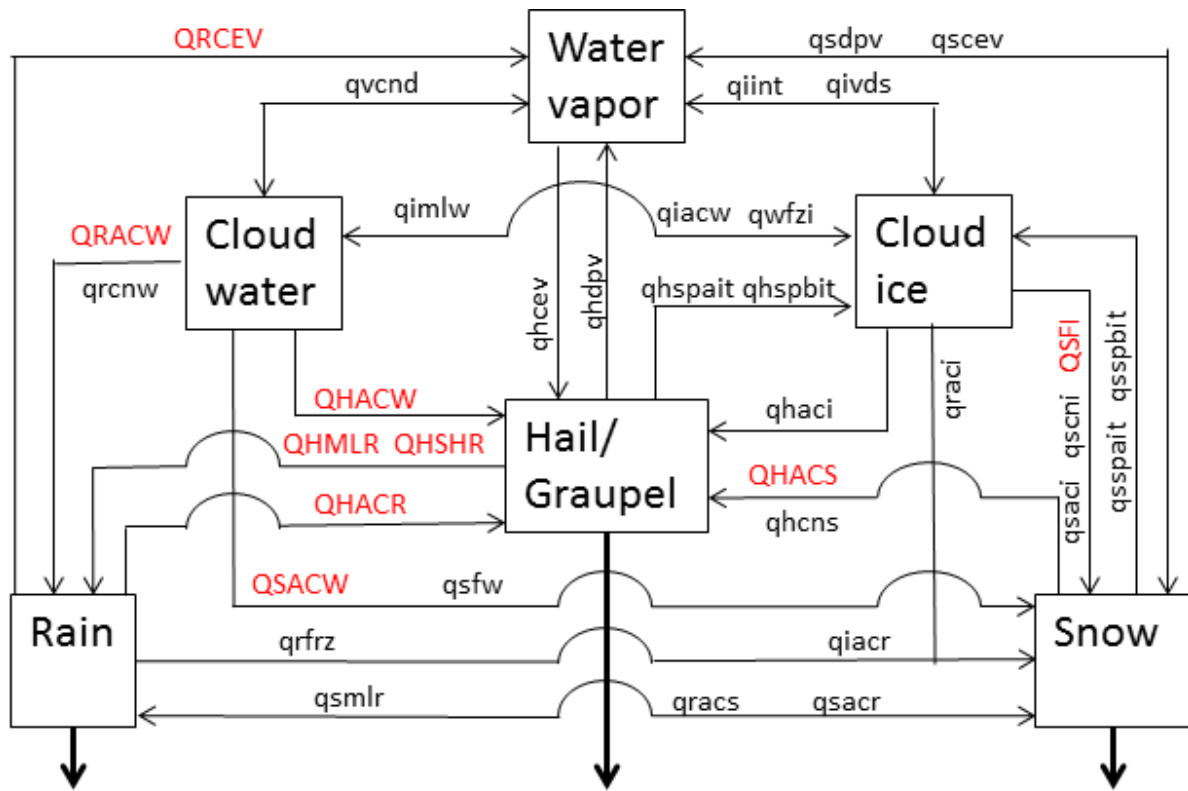


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Figure 6. Sources and sinks of each precipitating hydrometeor class for the Darwin, 23 January 2006 storm: (a) rain sources, (b) rain sinks, (c) snow sources, (d) snow sinks, (e) hail/graupel sources, (f) hail/graupel sinks. See the appendix for the microphysical meanings of the acronyms.



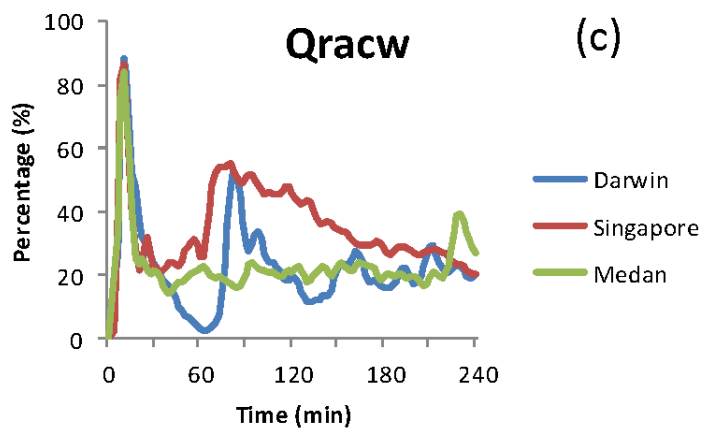
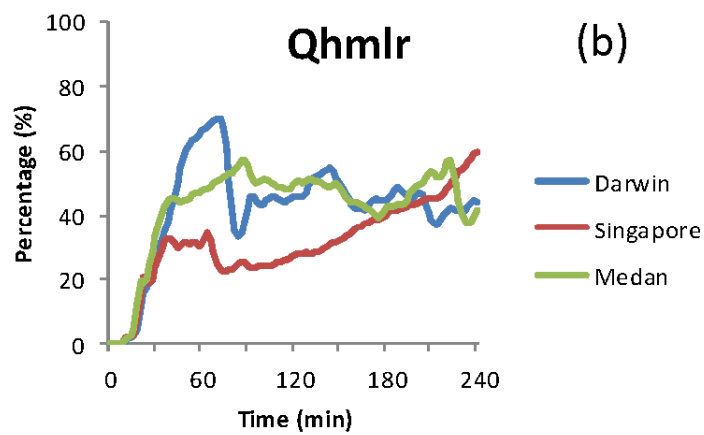
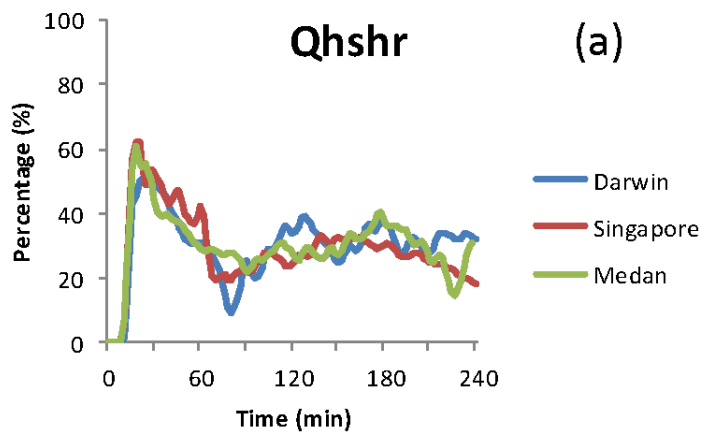
660
 661 **Figure 7.** As in Figure 6 but that the precipitating hydrometeor sources and sinks are plotted by
 662 percentage contributions.
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Figure 8. A flow chart of all microphysical processes. Red capital letters are the major ones. See the appendix for the microphysical meanings of the acronyms.

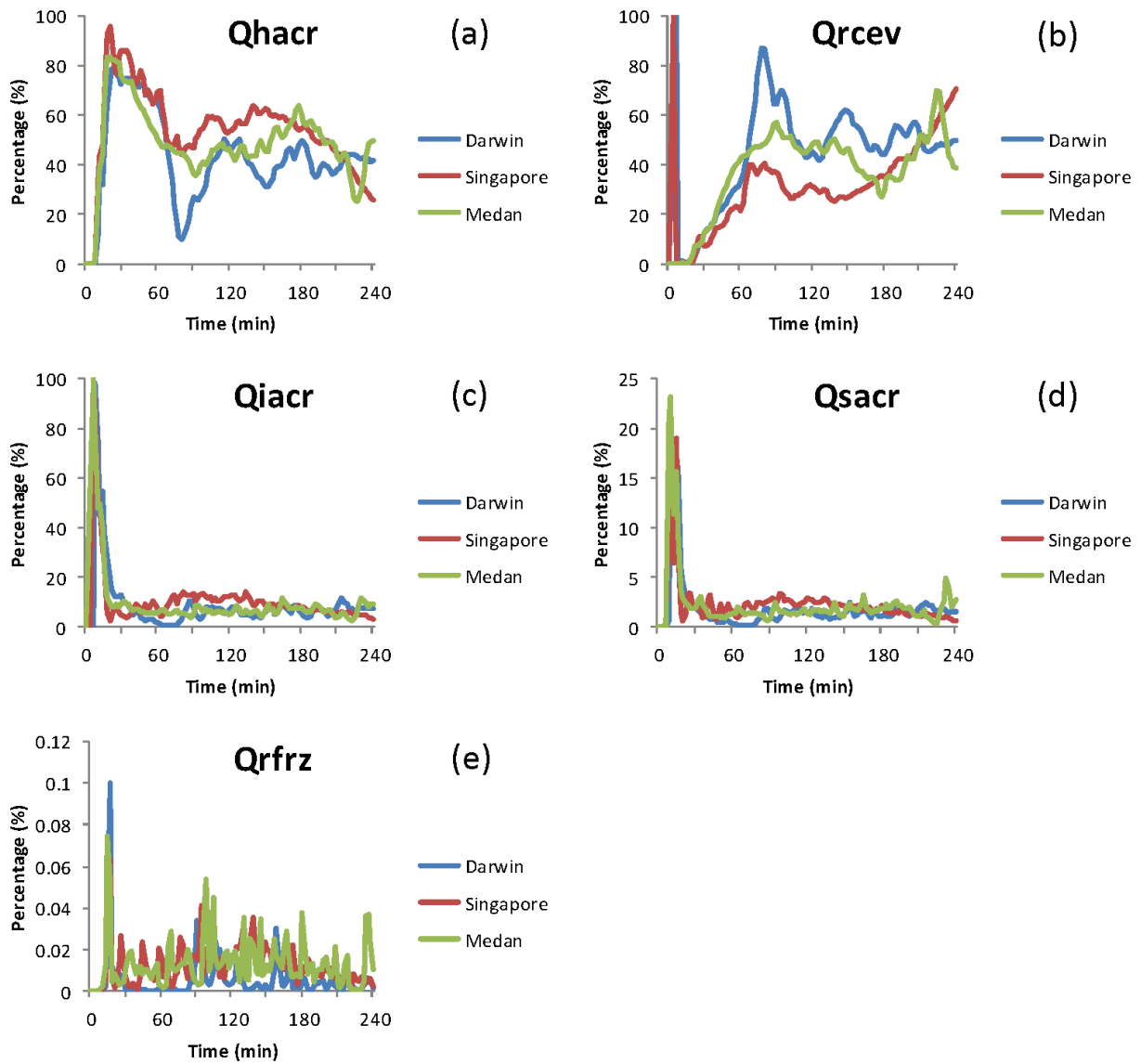
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670

671 **Figure 9.** Superimposed time series in all three storm simulations for the rain sources: (a) qhshr,
 672 (b) qhmlr, (c) qracw. See the appendix for the microphysical meanings of the acronyms.

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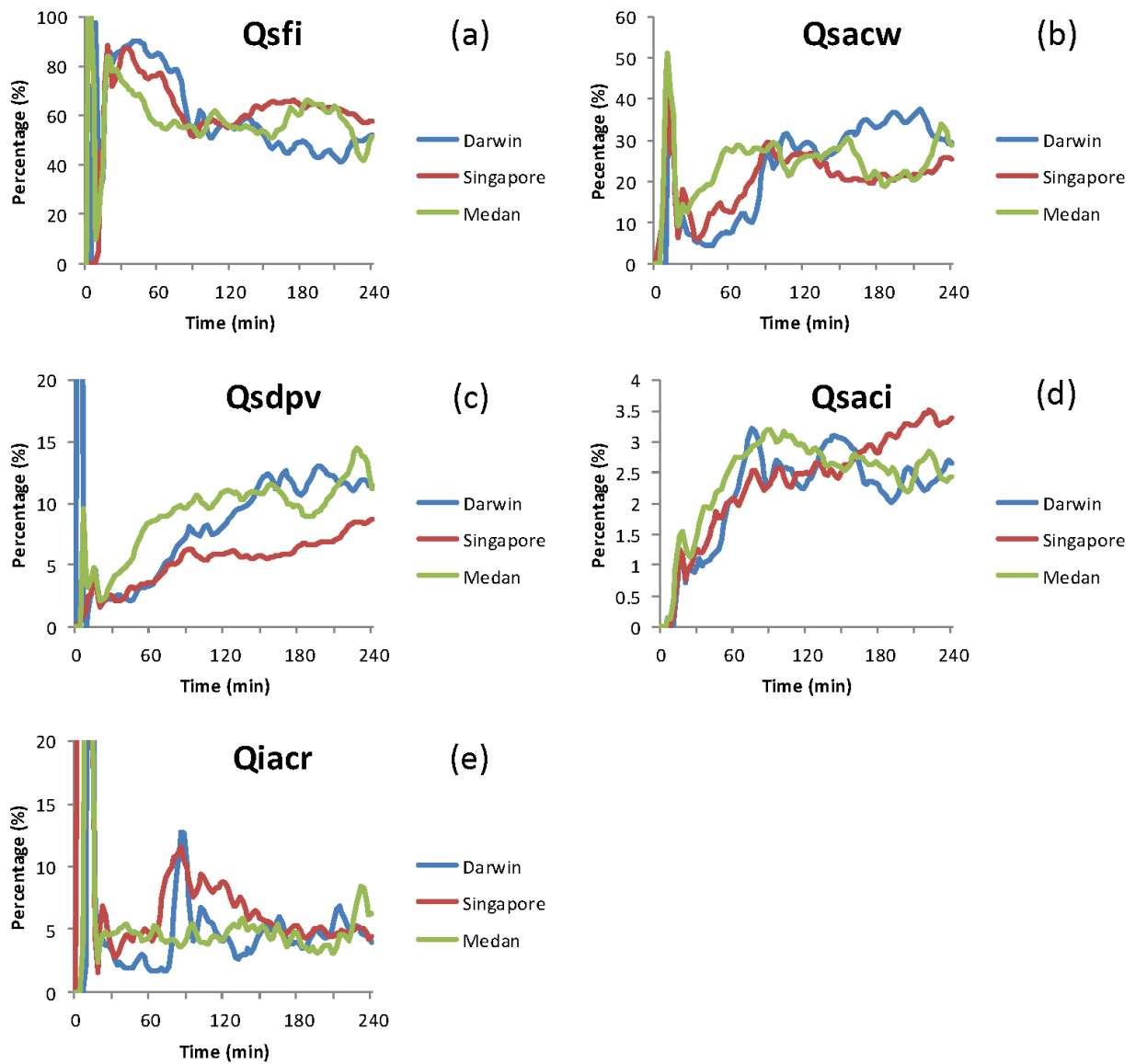


674

675 **Figure 10.** Superimposed time series in all three storm simulations for the rain sinks: (a) qhacr,
 676 (b) qrcev, (c) qiacr, (d) qsacr, (e) qrfrz. See the appendix for the microphysical meanings of the
 677 acronyms.

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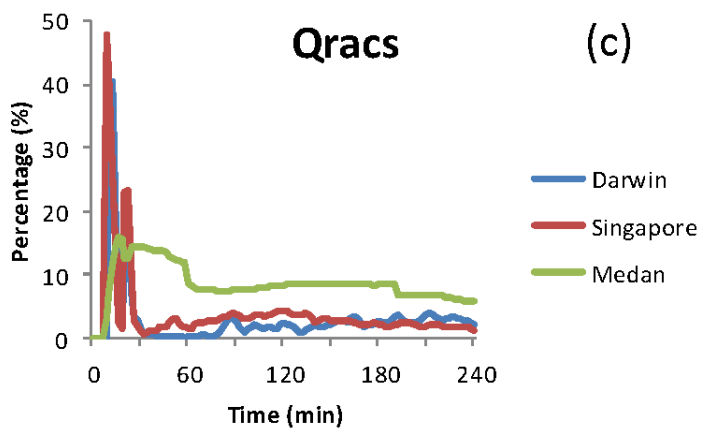
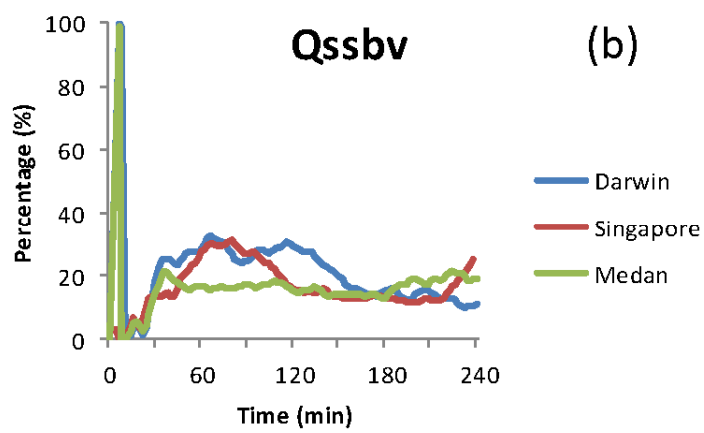
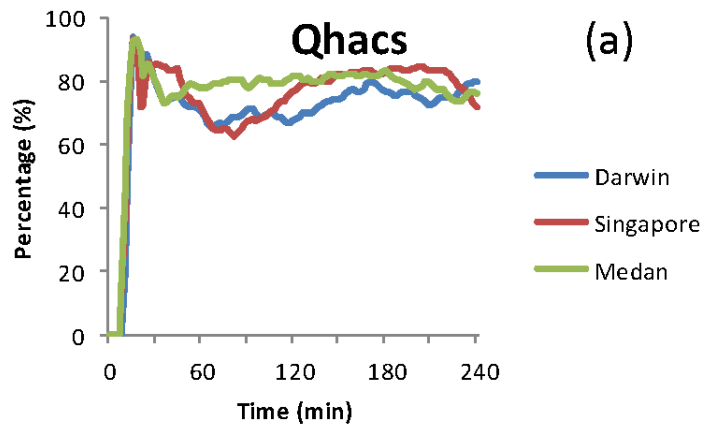
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681 **Figure 11.** Superimposed time series in all three storm simulations for the snow sources: (a) qsf_i,
 682 (b) qsac_w, (c) qsd_{pv}, (d) qsac_i, (e) qiac_r. See the appendix for the microphysical meanings of the
 683 acronyms.

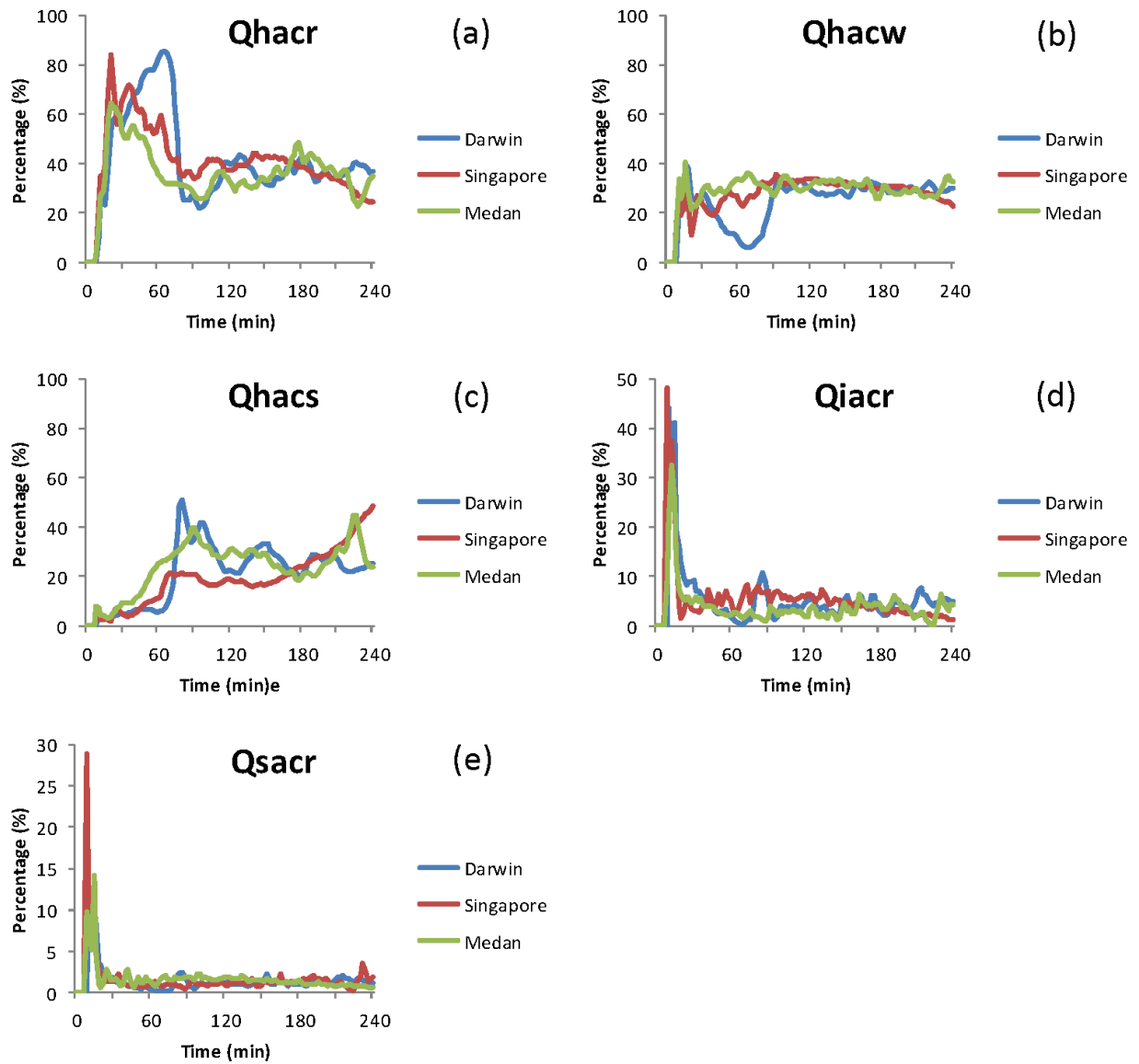
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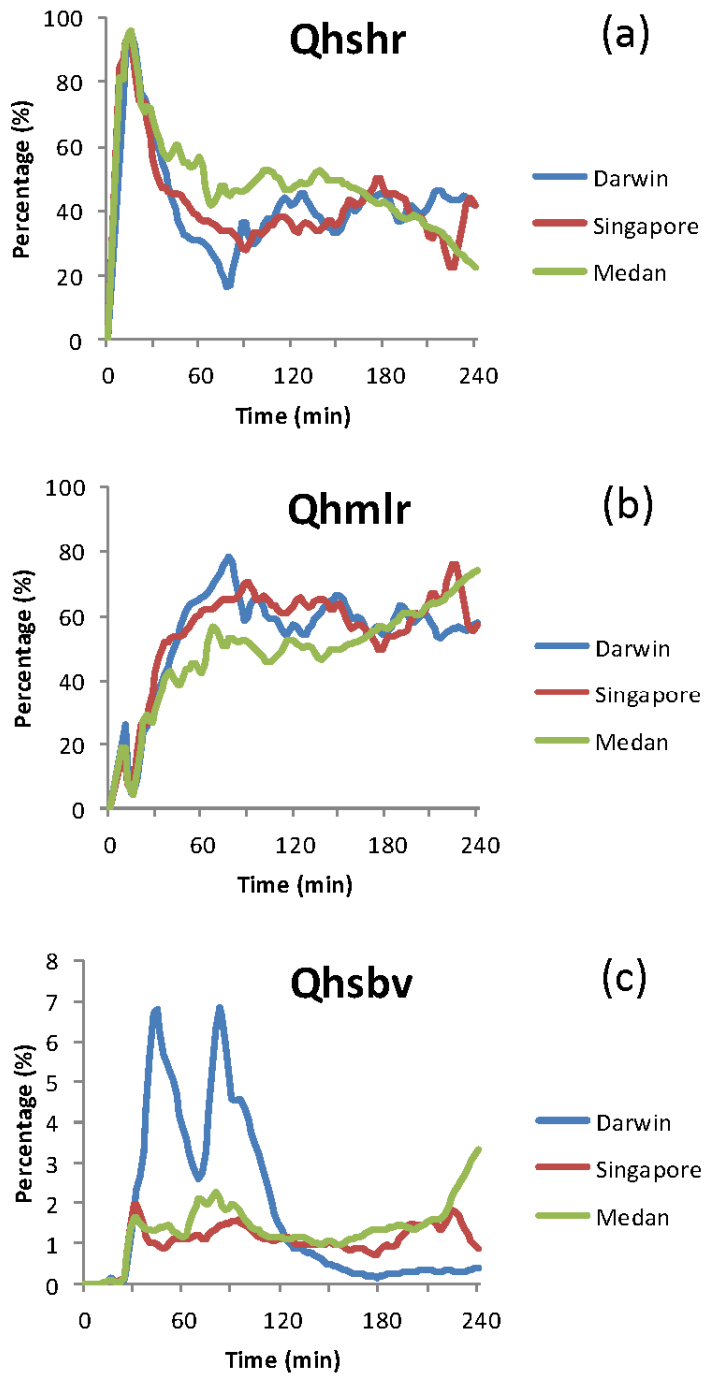
686 **Figure 12.** Superimposed time series in all three storm simulations for the snow sinks: (a) qhacs,
687 (b) qssbv, (c) qracs. See the appendix for the microphysical meanings of the acronyms.

688



689

690 **Figure 13.** Superimposed time series in all three storm simulations for the hail/graupel sources:
 691 (a) qhacr, (b) qhacw, (c) qhacs, (d) qiacr, (e) qsacr. See the appendix for the microphysical
 692 meanings of the acronyms.
 693



694

695 **Figure 14.** Superimposed time series in all three storm simulations for the hail/graupel sinks: (a)
 696 qhshr, (b) qhmlr, (c) qhsbv. See the appendix for the microphysical meanings of the acronyms.

697

698 **The appendix.** Definition of acronyms for microphysical processes

Acronym	Process
qhaci	Accretion of cloud ice by hail
qhacr	Accretion of rain by hail
qhacs	Accretion of snow aggregates by hail
qhacw	Accretion of cloud water by hail
qhcev	Condensation/Evaporation of vapor to/from wet hail
qhens	Autoconversion of snow to hail
qhdpv	Vapor deposition to hail
qhmlr	Melting of hail to rain
qhsbv	Vapor sublimation from hail
qhshr	Rain water shed from hail
qhspait	Secondary production I of cloud ice from hail
qhspbit	Secondary production II of cloud ice from hail
qiacr	Accretion of rain by ice to form snow or hail
qiacw	Accretion of cloud water by cloud ice
qiint	Nucleation of pristine cloud ice
qimlw	Melting of cloud ice to cloud water
qivds	Vapor Deposition/Sublimation to/from cloud ice
qraci	Accretion of cloud ice by rain to form snow or hail
qracs	Accretion of snow by rain to form snow or hail
qracw	Accretion of cloud water by rain
qrcev	Evaporation of rain
qrcnw	Autoconversion of cloud water to rain
qrfrz	Probabilistic freezing of rain to form snow or hail
qsaci	Accretion of cloud ice by snow
qsacr	Accretion of rain by snow to form snow or hail
qsacw	Accretion of cloud water by snow
qscev	Condensation/Evaporation of vapor to/from wet snow
qscni	Autoconversion of cloud ice to snow
qsdpv	Vapor deposition to snow
qsfi	Bergeron process, transfer of cloud ice to snow
qsfw	Bergeron process, transfer of cloud water to snow
qsmllr	Melting of snow to rain
qssbv	Vapor sublimation from snow
qsspait	Secondary production I of cloud ice from snow
qsspbit	Secondary production II of cloud ice from snow
qvend	Condensation/Deposition to/on cloud water/cloud ice
qves	Evaporation/Sublimation of cloud water/cloud ice
qwfzi	Homogeneous freezing of cloud water to cloud ice