

1 South Atlantic Interbasin Exchanges of Mass, Heat, Salt and 2 Anthropogenic Carbon

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19 **Abstract**

20 The exchange of mass, heat, salt and anthropogenic carbon (C^{ant}) between the South
21 Atlantic, south of $24^{\circ}S$, and adjacent ocean basins is estimated from hydrographic data
22 obtained during 2008-2009 using an inverse method. Transports of anthropogenic carbon
23 are calculated across the western (Drake Passage), eastern ($30^{\circ}E$) and northern ($24^{\circ}S$)
24 boundaries. The freshwater overturning transport of 0.09 Sv is southward, consistent with
25 an overturning circulation that exports freshwater from the North Atlantic, and consistent
26 with a bistable Meridional Overturning Circulation (MOC), under conditions of excess
27 freshwater perturbation. At $30^{\circ}E$, net eastward Antarctic Circumpolar Current (ACC)
28 transport, south of the Subtropical Front, is compensated by a 15.9 ± 2.3 Sv westward flow
29 along the Antarctic boundary. The region as a whole is a substantial sink for atmospheric
30 anthropogenic carbon of 0.51 ± 0.37 PgC yr⁻¹, of which 0.18 ± 0.12 PgC yr⁻¹ accumulates

31 and is stored within the water column. At 24°S, a 20.2 Sv meridional overturning is
32 associated with a 0.11 PgC yr⁻¹ C^{ant} overturning. The remainder is transported into the
33 Atlantic Ocean north of 24°S (0.28±0.16 PgC yr⁻¹) and Indian sector of Southern Ocean
34 (1.12±0.43 PgC yr⁻¹), having been enhanced by inflow through Drake Passage (1.07±0.44
35 PgC yr⁻¹). This underlines the importance of the South Atlantic as a crucial element of the
36 anthropogenic carbon sink in the global oceans.

37

38 **1 Introduction**

39 At the confluence of the southward-flowing deep water from the northern North Atlantic
40 Ocean and the eastward-flowing Antarctic Circumpolar Current (ACC), the South
41 Atlantic sector of the Southern Ocean is a key component of the global meridional
42 overturning circulation (MOC; Marshall and Speer 2012). The critical role of the South
43 Atlantic was recognised by Rintoul (1991), who quantified the basic heat and freshwater
44 exchange associated with balancing deep-water formation in the North Atlantic with
45 Intermediate Water and Bottom Water formation in the Southern Ocean. Overturning
46 within the South Atlantic is critical for the ventilation of older water masses facilitating
47 uptake and storage of anthropogenic carbon (C^{ant}) (Iudicone et al., 2011; Sallée et al.,
48 2012). Bottom Water formation; in particular, provides a mechanism for injection of C^{ant}
49 into the deep ocean (Brown et al., 2015; Vázquez-Rodríguez et al., 2009).

50

51 This paper focuses on the South Atlantic sector of the Southern Ocean south of 24°S
52 from Drake Passage to 30°E. The ACC crosses this region and, together with the Agulhas
53 Current, links the Pacific and Indian Ocean sectors of this region. The ACC transport is
54 concentrated into fronts (Subantarctic Front, SAF; Polar Front, PF; Southern ACC Front,
55 SACCF), which preferentially carry different water classes and properties across the
56 region (see Figure 1). Drake Passage is the narrow entry point for the ACC into the
57 Atlantic sector, after which, the Subantarctic Front protrudes northwards into the
58 Argentine Basin. This widens the meridional extent of the ACC, and separates the
59 warmer subtropical waters to the north from colder, Antarctic and Subantarctic water to
60 the south (Belkin and Gordon, 1996; Orsi et al., 1995).

61

62 North of the ACC in the Atlantic sector, the poleward-flowing Brazil Current (BC) lies
63 within the upper 300-600 dbar (Bryden et al., 2011; Peterson and Stramma, 1991). Fully
64 formed north of the Vitoria-Trinidad Seamounts at $\sim 20^{\circ}\text{S}$ (marked in Figure 1), it
65 intensifies southwards on the order of 5% per 100 km (Gordon and Greengrove, 1986)
66 with transport estimates at 24°S ranging between 4.1 Sv and 13.2 Sv (Bryden et al., 2011;
67 Evans et al., 1983; Evans and Signorini, 1985; Garfield, 1990; Signorini, 1978; Stramma,
68 1989; Zemba, 1991). At the eastern South Atlantic boundary within the Cape Basin, the
69 South Atlantic Current (SAC) feeds the northward flowing Benguela Current. Previous
70 transport estimates are of 6 Sv for the South Atlantic Current and 28 Sv for the Benguela
71 Current, respectively (Garzoli and Gordon, 1996; Mercier et al., 2003; Smythe-Wright et
72 al., 1998; Stramma and Peterson, 1990). The Benguela Current is also fed by the residual
73 westward flow into the South Atlantic from the Agulhas system, commonly termed
74 Agulhas leakage. The majority of the Agulhas Current flows along the East African
75 continent, and is retroflected at $16\text{-}20^{\circ}\text{E}$ (Lutjeharms and Van Ballegooyen, 1988) as the
76 eastward flowing Agulhas Return Current, closing the subtropical gyre of the South
77 Indian Ocean (Dencausse et al., 2010; Lutjeharms and Van Ballegooyen, 1988; Matano et
78 al., 1998).

79

80 South of the ACC in the Atlantic sector, previous studies (e.g. Meredith, 2013) have
81 suggested that the Weddell Sea contributes to about 40% of the global formation of
82 Antarctic Bottom Water (AABW). Westward inflow along the Antarctic shelf into the
83 Weddell Sea is partially comprised of recently formed Cape Darnley Bottom Water
84 (CDBW; Ohshima et al. 2013) and older AABW varieties from farther east. CDBW
85 contributes $\sim 13\text{-}30\%$ to global AABW production (Ohshima et al., 2013). Within the
86 Weddell Sea, local ventilation and interaction with the Filchner-Ronne (Whitworth et al.,
87 1998) and Larsen (Fahrbach et al., 1995; Weppernig et al., 1996) ice shelves contributes
88 to further AABW formation, carrying C^{ant} into the deep ocean (Huhn et al., 2013; van
89 Heuven et al., 2011). Some of this AABW recirculates within the eastward flowing
90 northern limb of the Weddell Gyre, whilst the remainder escapes either into the western
91 South Atlantic basin through narrow deep water pathways (e.g. Gordon et al., 2010,
92 Jullion et al., 2014), by South Scotia Ridge overflow (Jullion et al., 2014; Locarnini et al.,

93 1993; Naveira Garabato et al., 2002a), or into the eastern South Atlantic basin with 8 ± 2
94 Sv of AABW in total exported from the Weddell Gyre (Jullion et al., 2014). At the
95 Argentine Basin to Brazil Basin transition, northward AABW flow is restricted to key
96 topographical features (Figure 1): Vema Channel (25-50 km wide, sill depth ~ 4600 m;
97 Johnson and Biscaye (1976)) and Hunter Channel (200 km wide, sill depth ~ 4200 m;
98 Speer et al. (1992); Zenk et al. (1999)). Bottom water warming between the Weddell Sea
99 and 24°S alters the typical bottom water definition from $\theta \leq 0^\circ\text{C}$ to $\theta \leq 2^\circ\text{C}$. Bottom water
100 transports for $\theta < 2^\circ\text{C}$ are 4.0 ± 1.2 Sv at Vema Channel (Hogg et al., 1999), and 2.92 ± 1.24
101 Sv at Hunter Channel (Zenk et al., 1999).

102

103 The MOC, ACC, Agulhas system and Weddell Gyre are all major contributors to the
104 global large-scale ocean circulation, and therefore an understanding of their contribution
105 to interbasin fluxes is key for interpreting large-scale changes in volume, heat or
106 freshwater transports, and identifying linkages to broader changes in the Earth's climate.
107 Similarly interbasin fluxes of anthropogenic carbon (C^{ant}) provide an opportunity to
108 assess the South Atlantic's capacity to uptake and store anthropogenic carbon on decadal-
109 centennial timescales, in order to improve understanding of its responses to future
110 atmospheric CO_2 changes. Here, C^{ant} is estimated using the ΔC^* method following
111 Gruber et al. (1996), as described in section 2.1 and in further detail in Evans (2013). .

112

113 This paper uses a set of recent WOCE sections at the boundary of the South Atlantic
114 Ocean to update interbasin flux estimates of mass, heat and salt in comparison to earlier
115 studies (e.g. Rintoul, 1991), and to provide estimates of the interbasin flux of
116 anthropogenic carbon (C^{ant}). This paper is structured as follows: Section 2 describes the
117 data used. Section 3 outlines the inverse box methodology, as applied in this study. The
118 solution of the inverse box model is discussed in Section 4 in terms of geostrophic and
119 Ekman velocity fields, diapycnal mixing and air sea fluxes of heat and freshwater as well
120 as the transports of anthropogenic carbon at the South Atlantic boundary. The major
121 findings are described in Section 5.

122

123 **2 Data and Data Processing**

124 Hydrographic sections in Drake Passage (a repeat of World Ocean Circulation
125 Experiment (WOCE) section A21) in 2009, Africa to Antarctica along 30°E (repeat of
126 WOCE I6S) in 2008 and South America to Africa along 24°S in 2009 provide the data
127 for analysis. The Drake Passage and 24°S sections were occupied on board the research
128 vessel James Cook (King, 2010; McDonagh, 2009), with data stored within the British
129 Oceanographic Data Centre data archives, whilst the Africa to Antarctica occupation was
130 on board the Roger Revelle, with data stored by the CLIVAR (Climate Variability and
131 Predictability) and Carbon Hydrographic Data Office (CCHDO) (Speer and Dittmar,
132 2008; Wanninkhof et al., 2009).

133

134 Dissolved Inorganic Carbon (DIC) and Total Alkalinity were determined by coulometry
135 (Johnson et al. 1985, 1987, 1993; Johnson and Wallace 1992) and potentiometric titration
136 (Johnson et al., 1987; Dickson et al. 2003, 2007; Mintrop 2004), respectively. DIC and
137 Total Alkalinity were calibrated using Certified Reference Materials (CRM) (and gaseous
138 CO₂ loops for DIC along 30°E) to yield measurements with an accuracy of $\sim\pm 2-3 \mu\text{mol}$
139 kg^{-1} (Speer and Dittmar, 2008; McDonagh, 2009; King, 2010; Schuster et al. 2013,
140 2014). Oxygen was measured using Winkler titration (Culberson et al., 1991; Culberson
141 and Huang, 1987), whilst nitrate, phosphate and silicate measurements follow the
142 processes described in Gordon et al. (1993) and Kirkwood (1996). Estimated accuracies
143 according to CARINA methodology are oxygen (1%) and nutrients (2%) (Key et al.,
144 2010). All salinities used are on the PSS-78 scale (Fofonoff and Millard, 1983).

145

146 Hydrographic properties were recorded using a conductivity-temperature-depth (CTD)
147 profiler in 2 dbar intervals, to enable geostrophic transport estimates. Geostrophic
148 velocity within the 'bottom triangle' is set by nearest neighbour extrapolation to the
149 deepest common level for each station pair. DIC, nutrient and alkalinity measurements
150 are recorded for a maximum of 24, or 36 discrete depths per station for Drake Passage
151 and 24°S, and 30°E, respectively. Potential temperature (θ), salinity and oxygen are
152 linearly interpolated onto a 20 dbar vertical grid along the sections. Correction factors are
153 applied, as recommended by the GLODAP (Global Ocean Data Analysis Project) and
154 CARINA (Carbon in Atlantic Ocean) projects, listed in Table 1, to eliminate systematic

155 measurement biases (see Gouretski and Jancke, 2000; Hoppema et al., 2009; Key et al.,
156 2010, 2004; Lauvset et al., 2016; Olsen et al., 2016; Tanhua et al., 2010; Wanninkhof et
157 al., 2003).

158

159 The geographical locations of the sections are displayed in Figure 1. ACC fronts along
160 Drake Passage are determined as a distinct transition in θ -S space between θ -S
161 hydrographic station profiles. Each transition represents an ACC front separating each
162 frontal zone, and follows the Cunningham et al. (2003) analysis. Across 30°E,
163 thermohaline frontal definitions from Orsi et al. (1995) and Belkin and Gordon (1996) are
164 applied.

165

166 **2.1 Anthropogenic Carbon Calculation**

167 Anthropogenic carbon is estimated here using the ΔC^* method, whereby biological
168 effects, a pre-industrial background signal (based on ocean-atmospheric equilibrium
169 (C^{eqm}) (Brewer, 1978; Chen and Millero, 1979) and an estimate of ocean-atmosphere
170 disequilibrium (C^{diseq}) (Gruber et al., 1996) are removed from the modern inorganic
171 carbon signal. C^{eqm} is calculated based on pre-industrial fugacity ($fCO_2 = 280 \mu atm$), and
172 present day potential temperature, salinity, silicate and phosphate using “CO2SYS.m”
173 (Lewis and Wallace, 1998). C^{diseq} is represented using the linearised parameterisations for
174 specified potential temperature intervals from Pardo et al. (2011) (Indian/Pacific Ocean)
175 and Vázquez-Rodríguez et al. (2012) (Atlantic Ocean) and an Optimum Multiparameter
176 Analysis (OMP) technique (Karstensen and Tomczak, 1998; Sabine et al., 2002) below
177 the 5°C isotherm. C^{eqm} and C^{diseq} utilise the potential total alkalinity parameterisation
178 from Vázquez-Rodríguez et al. (2012), and the conversion from potential total alkalinity
179 to total preformed alkalinity following Brewer et al. (1975) and Fraga and Álvarez-
180 Salgado (2005). The uncertainty of C^{ant} estimates calculated using this method is up to ~6
181 $\mu mol kg^{-1}$ (Sabine et al., 1999). For visualisation and comparison in Section 4.3.1, a two-
182 dimensional distribution in neutral density:geopotential height or neutral
183 density:longitude space of C^{ant} is generated by least squares fitting using a ± 2 station and
184 $\pm 0.04 \gamma^n$ grid box centred at each CTD grid point. The geopotential height (ϕ) field is
185 calculated from the geopotential height anomaly at 500 dbar relative to 1500 dbar for

186 neighbouring stations. For transport calculations in Section 4.3.2, a two-dimensional
187 distribution in pressure:latitude or pressure:longitude space of C^{ant} is generated by least
188 squares fitting using a ± 2 station and ± 80 dbar grid box centred at each CTD grid point.
189 Full details of the calculation of C^{ant} are found in Evans (2013).

190

191 **2.1.1 Anthropogenic carbon storage**

192 Two independent methods are used to calculate the rate of accumulation of anthropogenic
193 carbon within the South Atlantic water column. The first is based on the assumption of a
194 transient steady state relationship between surface carbon changes and at depth, following
195 the methodology of Holfort et al. (1998) and Álvarez et al. (2003). This approach has
196 been indicated to be broadly consistent with Green's Function and inverse approaches
197 (Khatiwala et al., 2013). Secondly we use results of the Time Series Residual (TSR)
198 approach (van Heuven et al., 2011; van Heuven, 2013), where a residual DIC is
199 calculated from the difference between measured DIC and synthetic DIC values
200 constructed from a multivariate linear regression of all available data points. The time
201 trend of that residual DIC is interpreted as equivalent to the time trend of C^{ant} . The TSR-
202 based C^{ant} storage estimate (provided by van Heuven, S. (2016), manuscript in
203 preparation) uses all historical carbon data from 1972-2012 from the GLODAPv2 data
204 product (Olsen et al., 2016) as well as the climatologies produced therewith (Lauvset et
205 al., 2016).

206

207 **2.1.1.1 Mean penetration depth (MPD)**

208 In general terms, the build-up of carbon within any ocean basin is given by the difference
209 between box boundary transports and atmosphere – sea-surface exchange, whilst
210 assuming a negligible effect for a number of compensatory processes (following Álvarez
211 et al. (2003)): riverine input (Holfort and Siedler, 2001; Jacobson et al., 2007), meltwater
212 input (Rignot et al., 2008), net organic carbon production and sediment burial (Rosón et
213 al., 2003; Sarmiento et al., 1995), and calcium carbonate dissolution and burial (Stoll et
214 al., 1996). Other transient terms related to seasonal or biological variability at the box
215 boundary are assumed as negligible for the basin-wide C^{ant} storage estimate. As diapycnal
216 and air-sea induced diapycnal transfer do not add or remove C^{ant} from the full depth C^{ant}

217 budget, the inclusion of only geostrophic and Ekman effects therefore create the
 218 following equation:

$$C^{\text{ant}} \text{ Storage} = F_{\text{air-sea}} + T_N + T_W + T_E \quad (1)$$

219 with the average air-sea C^{ant} flux being $F_{\text{air-sea}}$ and T_N , T_W and T_E being C^{ant} transports
 220 across the northern, western and eastern boundaries of the South Atlantic Ocean sector
 221 (Figure 1). For C^{ant} , storage is represented by the temporal increase of C^{ant} throughout the
 222 water column, or mathematically by:

$$C^{\text{ant}} \text{ Storage rate} = \frac{d \int C_z^{\text{ant}} dz}{dt} \quad (2)$$

223 where t is time and $\int C_z^{\text{ant}} dz$ is the accumulation of anthropogenic CO_2 at each depth level
 224 z yielding a storage rate with units of $\text{mol m}^{-2} \text{ yr}^{-1}$ (Álvarez et al., 2003). An
 225 approximation for the magnitude of the anthropogenic CO_2 storage is calculated from the
 226 mean penetration depth (MPD) from Broecker et al., (1979):

$$\text{MPD} = \frac{\int C_z^{\text{ant}} dz}{C_{\text{ml}}^{\text{ant}}} \quad (3)$$

227 where C_z^{ant} and $C_{\text{ml}}^{\text{ant}}$ are anthropogenic CO_2 estimates at depth, z , and within the mixed
 228 layer, respectively. MPD is therefore the C^{ant} column inventory divided by C^{ant} from the
 229 mixed layer, and always yields a depth which is shallower than the actual depth to which
 230 the tracer penetrates (Peacock, 2004). Combining equation 2 and 3 gives an estimate of
 231 the anthropogenic CO_2 storage rate:

$$C^{\text{ant}} \text{ Storage rate} = \frac{d \int C_z^{\text{ant}} dz}{dt} = \text{MPD} \times \frac{dC_{\text{ml}}^{\text{ant}}}{dt} \quad (4)$$

232 This assumes that the vertical profile of C^{ant} is constant in shape and scale depth with
 233 time following the transient steady-state assumption of Gammon et al. (1982). This
 234 prescribes that a conservative tracer propagating into an ocean with steady circulation,
 235 but forced by an exponentially-increasing atmospheric boundary source function, will
 236 reach a transient steady state with constant shape. Mixed layer increases in C^{ant} are then
 237 assumed to increase proportionally with tracer concentrations at all depths (Tanhua et al.,
 238 2007). C^{ant} is thought to have passed into transient steady-state, given the length of its
 239 atmospheric history (>200 years). The MPD assumptions are most problematic in regions
 240 of significant deep water ventilation, where the assumption of a constant vertical C^{ant}
 241 profile to the ocean bottom may be false. In this study, recently ventilated deep waters, in

242 the form of AABW along Drake Passage and 30°E sections, still maintain low C^{ant} . This
243 helps to validate the usage of an MPD-based C^{ant} storage rate estimate in this instance,
244 however, this methodology contributes to a relatively large uncertainty in the result.

245

246 The $\Delta C_{\text{ml}}^{\text{ant}}$ is calculated by determining the rate of change in mean C^{ant} within the mixed
247 layer between occupations. In this study, $\Delta C_{\text{ml}}^{\text{ant}}$ is computed using historical hydrographic
248 occupations of Drake Passage (Meteor: 1990) and 30°E (Marion Dufresne: 1996) with
249 further details in Evans (2013). Along the 24°S transect, $\Delta C_{\text{ml}}^{\text{ant}}$ is calculated based on
250 overlapping stations from meridional hydrographic occupations A13, A14 (Mercier and
251 Arhan, 1995), A15 (Smethie and Weatherly, 1994), A16 (Talley et al., 1989) and A17
252 (Mémery, 1994) within the South Atlantic. This constitutes all historical data for the
253 region available within GLODAPv2 (see Appendix B for details). The small sample of
254 repeat DIC measurements at the northern boundary increases storage uncertainty. Storage
255 rate is re-written as:

$$\text{Storage rate} = \text{MPD} \times \Delta C_{\text{ml}}^{\text{ant}} \times \rho_{\text{ml}} \quad (5)$$

256 where ρ_{ml} is the in-situ density within the mixed layer yielding storage rate with units of
257 $\text{mol m}^{-2} \text{yr}^{-1}$.

258

259 **2.1.1.2 Time Series Residual (TSR)**

260 TSR-based C^{ant} storage estimates rely upon assumptions that (i) the relationship between
261 DIC and the independent variables in the regression is linear, that (ii) bias and noise
262 within the sampling is considered negligible (or average out for the large dataset
263 employed) and that (iii) real changes in one or more independent variable is associated
264 with changes in one or more of the other independent variables (van Heuven, 2013). The
265 time trend in C^{ant} is expected to depend upon the ventilation age of the water mass, with
266 AOU used as a proxy for ventilation age. For a particular water mass, i , the time trend of
267 C^{ant} is represented by the linear regression of:

$$\frac{dC_i^{\text{ant}}}{dt} = a_i + \Delta \text{AOU} \cdot b_i \quad (6)$$

268 Where ΔAOU is the difference between the AOU of the sample and the mean AOU in
269 the water mass core (van Heuven, 2013). The contribution of a water mass to a given

270 sample is determined using Optimum Multiparameter analysis (OMP) (Karstensen and
 271 Tomczak, 1998; Sabine et al., 2002; van Heuven, 2013). The C^{ant} storage for the South
 272 Atlantic basin is thus estimated by the inclusion of a gridbox mass following:

$$\frac{d^{\text{INV}}C_i^{\text{ant}}}{dt} = \sum_{i=1}^j x_i \cdot (a_i + \Delta\text{AOU} \cdot b_i) \cdot \text{GBM} \quad (7)$$

273 Where x_i is the fractional contribution of water mass i to the inventory, and GBM
 274 represents the mass of a grid box surrounding each grid point, as described in van
 275 Heuven, (2013). The resulting inventory ($d^{\text{INV}}C_i^{\text{ant}}/dt$) can be expressed in units of PgC
 276 yr^{-1} .

277

278 **3 Box Inverse analysis**

279 **3.1 Setup**

280 The box inverse framework combines initial estimates of the circulation on each of the
 281 three hydrographic sections (Section 3.2) with constraints on the large-scale circulation,
 282 convergence of properties in the box, mixing and air-sea fluxes (Section 3.3). This
 283 generates an estimate of the circulation, the solution that is consistent across all three
 284 sections and the enclosed region (Section 3.4). The hydrographic sections used here
 285 (Figure 1) were made in February and early April, however, in either 2008 or 2009. The
 286 lack of synopticity of the data increases the uncertainty; however, this is partially
 287 accounted for by the choice of constraints to avoid a synoptic bias. This solution for this
 288 inverse box model is therefore most representative of South Atlantic circulation during
 289 austral summer.

290

291 The setup and method used is summarised here and detailed in Wunsch (1996). The
 292 inverse box model with the additional inclusion of noise vector ε to account for errors
 293 (Evans, 2013), is represented by:

$$\mathbf{E}\mathbf{x} + \varepsilon = \mathbf{y} \quad (8)$$

294 \mathbf{E} is an $m \times n$ matrix, \mathbf{x} is an $m \times 1$ vector of unknowns and \mathbf{y} is an $m \times 1$ vector of the
 295 imbalance between the initial field and the constraints. The coefficients in \mathbf{E} represent the
 296 geometry of the section. Each row of \mathbf{E} represents a constraint on the system. Each

297 column of \mathbf{E} represents an unknown. In this study, the system has 340 unknowns and 73
298 constraints. The unknowns are those elements of the system that can be adjusted in order
299 to satisfy the constraints. The inverse model solves for 217 depth-independent \mathbf{x}
300 velocities, one from each pair of adjacent hydrographic stations on each section. In
301 addition, a single unknown represents the correction to the Ekman transport on the 24°S
302 section, whilst 60 unknowns represent the mixing of volume, temperature and salinity
303 between density layers within the box and another 62 unknowns represent the
304 transformation between layers driven by air-sea interaction.

305

306 **3.2 Initial Field**

307 Flow across the sections is assumed to be geostrophic with an additional surface Ekman
308 transport across 24°S. An initial reference level and geostrophic field is constructed for
309 each section (Table 2) based on historical analysis. The basic premise of the box inverse
310 is to adjust the strength of the reference velocity at each station pair so that constraints are
311 satisfied within a given uncertainty (Section 3.2). In addition in this study, the box
312 inverse allows for a correction to initial estimates of the mixing between neutral density
313 layers, air sea fluxes and an Ekman transport. All diapycnal fluxes associated with
314 interior mixing or air-sea induced transformation are initialised to zero (McDonagh and
315 King, 2005). As the solution that is estimated is dependent upon the initial field, it is
316 important that the initial field is as representative as possible.

317

318 At Drake Passage, the reference level choice (Table 2) of the deepest common level
319 between the station pairs is based on the analysis of the mean volume transport of
320 multiple repeat stations across Drake Passage of $136.7 \pm 6.9 \text{ Sv}$ (Cunningham et al 2003,
321 Meredith et al., 2011), Lowered Acoustic Doppler Current Profiler transport estimates
322 (Meredith et al., 2011) and the scale of interannual variability (King and Jullion, in
323 prep.). At 24°S, the 1300 dbar reference level approximates the upper water/NADW
324 interface. At 30°E, Bryden et al. (2005) and Arhan et al. (2003) are used as a guide for
325 the vertical transition between the Agulhas Current and Agulhas Return Current, and the
326 Agulhas Undercurrent at depth (Beal and Bryden, 1999). On all sections the geostrophic

327 velocity within the ‘bottom triangle’ is set by nearest neighbour extrapolation to the
328 deepest common level for each station pair.

329

330 For the 24°S section, Ekman transport from NCEP (National Centers for Environmental
331 Prediction) wind stresses, an annual average calculated between 1980-2010 in Bryden et
332 al. (2011), of 3.3 Sv southward is applied as a single velocity above the 80 dbar Ekman
333 depth (D_{EK}). The Ekman component is included at 24°S as part of the initial field.

334

335 **3.3 Constraints**

336 **3.3.1 Constraints to circulation and property transports on sections**

337 The constraints across hydrographic sections, based on historical analyses and listed in
338 Table 3, are applied to better constrain the initial field, and later used to constrain the box
339 inverse model. Further details regarding the constraints in Table 3 are described below.

340

341 Across Drake Passage, full-depth volume transport is constrained to 136.7 Sv
342 (Cunningham et al., 2003; Meredith et al., 2011).

343

344 Bottom Water (BW) across 24°S has been defined to be below the 2 °C isotherm (Hogg
345 et al., 1999; McDonagh et al., 2002), shallower than the typical AABW neutral density
346 class definition (neutral density: $\gamma^n > 28.27$) in the Southern Ocean, and partly includes the
347 lower layers of the LCDW neutral density class within the Vema Channel and Hunter
348 Channel. Northward BW flow is constrained following Hogg et al. (1999), Zenk et al.
349 (1999) and McDonagh et al. (2002), as 6.9 Sv below the 2 °C isotherm. Within the
350 northern Cape Basin, east of Walvis Ridge (6°E), a zero mass transport constraint is
351 applied below the 2 °C isotherm (Arhan et al., 2003; McDonagh and King, 2005). For the
352 sectionwide upper 80 dbar, a southward, wind-driven estimate for the Ekman transport of
353 3.3 Sv is included following Bryden et al. (2011). For the upper 300 dbar, west of 35°W,
354 the Brazil Current is constrained to 4.9 Sv southward (Bryden et al., 2011). Finally, full
355 depth salinity transport across 24°S is constrained to be equal to the Bering Strait salinity
356 transport of 26.0 Sv psu, assuming salinity conservation (Coachman and Aagaard, 1988).

357

358 For the 30°E section, north of the Subtropical Front (42.9°S), the residual westward flow
359 of warm, salty Indian Ocean water into the Atlantic Ocean or ‘Agulhas leakage’ is
360 estimated based on McDonagh et al. (1999) as 9 Sv above the 3.5 °C isotherm. Finally, a
361 box-wide constraint for zero net salinity divergence is applied by summing together
362 salinity transport through the Agulhas regime, Drake Passage and across 24°S. Total
363 salinity transport outflow across the 30°E ACC regime is adjusted to match the inflow
364 across Agulhas regime, Drake Passage and 24°S (Table 3). The residual mass transport is
365 interpreted as the freshwater flux of the initial field.

366

367 **3.3.2 Property constraints in the box**

368 Each transect is split into 21 neutral density (γ^n) layers (Table 4; Jackett and McDougall
369 (1997)). Neutral density class interfaces, appropriate for the Southern Ocean, are
370 extracted from Heywood and King (2002), Naveira Garabato et al. (2009, 2002a, 2002b)
371 and Orsi et al. (1999, 1995). The layers are grouped into six neutral density classes. Each
372 γ^n layer represents an equation to be solved for, with an additional row for the full depth
373 water column. Conservation of mass, heat and salt (approximated as volume, potential
374 temperature anomalies and salinity anomalies) for each layer plus full depth conservation
375 gives 66 equations or constraints for the analysis. Additionally, full-depth silicate
376 conservation plus 6 constraints from previous knowledge of the circulation (Table 3)
377 gives a total of 73 constraints. Salinity and θ within each γ^n layer are conserved in the
378 form of a property anomaly, calculated by subtracting each property value by the
379 boundary-wide average, calculated using the whole domain boundary. The use of
380 property anomalies improves the matrix conditioning (Ganachaud, 2003; McIntosh and
381 Rintoul, 1997). For silicate, as argued by Ganachaud (1999), property anomalies are not
382 calculated given the large concentration range between surface and deep waters. Loss of
383 silicate through opal deposition is assumed negligible, given large uncertainties in the
384 silicate budget (Tréguer and De La Rocha, 2013) with this assumption encouraging
385 conservation within the silicate-rich bottom waters.

386

387 **3.4 Solution**

388 **3.4.1 Unknown velocities**

389 In this study, the columns of \mathbf{E} are constructed to solve for unknowns; geostrophic,
390 diapycnal, air-sea fluxes and Ekman transports, and each row in \mathbf{E} represents an equation
391 or constraint. In order to better condition the pre-inversion matrix for solving for the
392 unknown velocities, each row and each column of the $m \times n$ coefficient matrix \mathbf{E} is
393 weighted based on estimates of the previously known, ‘a priori’ uncertainties within each
394 component (see Appendix A). Solution weightings are applied as stated in Appendix A
395 following the method of McDonagh and King (2005) and Tsubouchi et al. (2012).

396

397 The geostrophic component of each cross-sectional station pair is applied with an a priori
398 uncertainty of $1 \times 10^{-2} \text{ m s}^{-1}$, as in Naveira Garabato et al. (2003), McDonagh and King
399 (2005) and Jullion et al. (2010). The a priori uncertainty is uniform for all station pairs
400 across all transects.

401

402 For the inverse model, the Ekman transport adjustment is initialised as a single unknown.
403 The coefficient matrix \mathbf{E} , initialised for a single unknown representative of the Ekman
404 transport adjustment, is initialised by the area above D_{Ek} , the property mean of the Ekman
405 layer, and the proportional contribution of the Ekman transport to each γ^n layer above
406 D_{Ek} . As the climatological data contains uncertainties, which are difficult to quantify, an
407 a priori uncertainty of 50% of the initial estimate of the Ekman transport adjustment is
408 assigned.

409

410 **3.4.1.1 Interior diapycnal velocities**

411 A separate diapycnal velocity is resolved for each property (McIntosh and Rintoul, 1997)
412 and for each layer interface. The interface mean for each property (S , θ) is generated
413 using the WOCE Global Hydrographic Climatology (WGHC) by Gouretski and
414 Koltermann (2004). The WGHC data is on a 0.5° grid, and averaged along isopycnal
415 surfaces, such that the properties are broadly in agreement with the properties along the
416 sections. The layer interface area for each of the neutral density interfaces in this study is
417 constructed from the initial 45 levels from WGHC for each mapped property field. For
418 the diapycnal mixing, a priori uncertainties are dependent on the pre-existing estimates of
419 diapycnal velocities (ω) and assigned as 10^{-5} m s^{-1} , following Orsi et al. (1999) and

420 Naveira Garabato et al. (2003), for an estimate of an upper value for deep ocean
421 diapycnal velocities.

422

423 **3.4.1.2 Diapycnal transfers induced through Air-Sea interactions**

424 Heating and cooling of neutral density classes, as the isopycnals outcrop at the ocean
425 surface provides a mechanism for across isopycnal transformation (Speer and Tziperman,
426 1992; Tziperman and Speer, 1994). Within the Southern Ocean, all layers are assumed to
427 outcrop given the upwelling of deep neutral density classes. Following Jullion et al.
428 (2010a), net air-sea fluxes of mass (freshwater) M_v and heat M_θ are calculated for each
429 layer, whilst the diapycnal volume flux induced by air-sea interaction F_v is included for
430 each layer interface. The area of outcrop for each neutral density layer is estimated from
431 monthly averaged sea surface temperature and salinity fields from World Ocean Atlas
432 (WOA) on a 1° grid (Antonov et al., 2010; Boyer et al., 2005; Locarini et al., 2010). To
433 ensure an area of outcrop for the densest γ^n layers, the area of outcrop for all LCDW and
434 AABW layers was averaged, and this value was assigned to all LCDW and AABW
435 layers.

436

437 Heat flux terms are supplied by monthly-averaged estimates from the National
438 Oceanography Centre (NOC v2.0) climatology (Berry and Kent, 2011, 2009). Net heat
439 flux Q_{net} is the sum of contributions from latent (Q_H) and sensible heat flux (Q_E),
440 longwave flux (Q_{LW}) and shortwave flux (Q_{SW}) (Grist and Josey, 2003). The mean heat
441 flux for the January-February-March (JFM) period is 65 W m^{-2} .

442

443 Freshwater flux is based on the climatologies recommended by Schanze et al. (2010):
444 Global Precipitation Climatology Project (GPCP) for precipitation (Adler et al., 2003),
445 and Objectively Analysed Ocean-Atmosphere Flux (OAFlux) for evaporation (Yu et al.,
446 2008; Yu and Weller, 2007). Evaporation is subtracted by precipitation (E-P) at each grid
447 point using the 2008 and 2009 estimates, before finding the inverse box model mean. A
448 priori uncertainties are estimated to be 50% of the initial estimates. Uncertainties arise
449 from the uncertainty of the climatologies as described in Lumpkin and Speer (2007), as
450 well as from not considering the contribution of sea-ice near the Antarctic continent.

451

452 **3.4.2 Choice of preferred solution**

453 The solution rank of 60 out of 73 is chosen after application of SVD. Truncation to the
454 solution rank occurs at the point at which the noise added by including additional rows
455 negates the information gained. The co-dependency between ocean layers gives reason
456 for selection of a solution rank below the full rank. Ranks $\sim >50$ are suitable solutions
457 with a full depth volume transport $\sim <1$ Sv, equivalent to the freshwater divergence.
458 Reference velocities for the geostrophic component are generally within ± 0.5 cm/s with
459 all adjustments off continental shelves within ± 0.7 cm/s.

460

461 **3.5 Model Diagnostics**

462 **3.5.1 Overturning freshwater and heat transport**

463 The overturning component of the salinity transport at 24°S is calculated for comparison
464 to the outputs of Bryden et al. (2011) using the M_{ov} salt transport, in addition to the
465 azonal component M_{az} . Additionally the heat transport associated with the ‘overturning’
466 and ‘gyre’ components is separated following the methods of Bryden and Imawaki
467 (2001) and Bryden et al. (2011).

468

469 For freshwater, values for M_{ov} and M_{az} are calculated following Bryden and Imawaki
470 (2001), Dijkstra (2007), Huisman et al. (2010) and Bryden et al. (2011):

$$M_{ov} = -1/\langle S \rangle \int \langle v \rangle (\langle S \rangle - \overline{\langle S \rangle}) L(z) dz \quad (9)$$

$$M_{az} = -1/\langle S \rangle \iint (v - \langle v \rangle) (S - \langle S \rangle) dx dz \quad (10)$$

471 where v is the northward velocity, S is salinity, L is zonal section width and z is depth.
472 Triangular brackets indicate a zonal average and an overline represents a vertical average.
473 The M_{ov} and M_{az} transports are effectively the freshwater transports associated with the
474 overturning and gyre circulation components, respectively. Cimadoribus et al. (2012)
475 suggest that an increase in the zonal salinity contrast across the South Atlantic increases
476 M_{az} and that this is compensated by a decrease in M_{ov} . Changes in M_{az} could therefore
477 dictate potential MOC shutdown (Cimadoribus et al., 2012).

478 The volume transports and overturning freshwater transports associated with the MOC
479 are detailed in section 4.1. Geostrophic and non-geostrophic results are described and
480 circulation features examined in section 4.2. For section 4.3, C^{ant} transports are calculated
481 for each layer, whilst C^{ant} air-sea flux is considered in section 4.3.3.

482

483 **4 Inverse Model Solution**

484 **4.1 Volume and overturning freshwater transports**

485 **4.1.1 Geostrophic solution**

486 The geostrophic velocities of the final solution are shown in Figure 2. The overall
487 velocity pattern is for strong flow into the box through Drake Passage and an outflowing
488 velocity along 30°E , south of the Subtropical Front. North of the Subtropical Front,
489 positive and negative velocities reflect the Agulhas Current inflow and Agulhas Return
490 Current outflow. The box-wide salinity transport conservation results in a net volume
491 imbalance of -0.47 Sv, interpreted as a loss of freshwater, balanced by excess
492 precipitation over the box.

493

494 The net transport (Figure 3, right) indicates convergence (positive numbers) or
495 divergence (negative numbers) of a neutral density class within the box. Convergence can
496 be interpreted as destruction of that neutral density class within the box and divergence
497 reflects production of that neutral density class. Basinwide UCDW layer convergence is
498 caused by upwelling of the MOC southern limb (see Section 4.1.2.1), resulting in
499 northward flowing surface and mode water and AABW layer formation to the south.
500 LCDW layer divergence corresponds with greater outflow across 30°E (44.4 Sv)
501 compared to Drake Passage inflow (28.1 Sv), caused by mixing the NADW and AABW
502 layers with the LCDW layer.

503

504 **4.1.1.1 Drake Passage**

505 The final solution decreases the Drake Passage initial field full-depth volume transport of
506 136.7 ± 10 Sv to 128.4 ± 8.3 Sv (Table 3). This is within the uncertainty of the volume
507 transport, estimated as 126.3 - 147.1 Sv (King and Jullion, in prep., and Meredith et al.
508 (2011) (their Figure 11)). Transport of UCDW layers constitutes almost half of the Drake

509 Passage full depth volume transport (58.1 Sv out of 128.4 Sv ; Figure 3), in agreement
510 with the 62.3 Sv estimate of Cunningham et al. (2003), relative to the deepest common
511 level. Within the SACCF, the transport is equally split between UCDW and LCDW
512 layers. The contribution of SAMW and AAIW layer transport increases progressively to
513 the north along the section.

514

515 **4.1.1.2 24°S**

516 For the Brazil Current, the final solution of $5.8 \pm 0.1 \text{ Sv}$ falls within the historical range as
517 described in Bryden et al. (2011) with the salty Brazil Current being important for the
518 total salinity transport across 24°S . Bottom water exchange from the northern Cape Basin
519 into the eastern South Atlantic basin is limited by Walvis Ridge. The final solution shows
520 $0.2 \pm 0.1 \text{ Sv}$ of southward AABW layer transport, and is similar to McDonagh and King
521 (2005)'s estimate of $0.1 \pm 0.5 \text{ Sv}$.

522

523 The southward basin-wide full-depth salinity transport at 24°S ($25.8 \pm 0.2 \text{ Sv}$ psu, Table
524 3) closely matches observations from the Bering Strait throughflow (Coachman and
525 Aagaard, 1988; Woodgate and Aagaard, 2005) and is similar to Holfort and Siedler
526 (2001)'s $26.75 \pm 0.77 \text{ Sv}$ southward salinity transport for the quasi-zonal A10 WOCE
527 section across 30°S . Historical meridional freshwater, heat and salt transports across
528 24°S , 30°S and 32°S are included for comparison with the results from our box inverse
529 (Table 5). Focussing firstly on net freshwater transport, the difference between 0.8 Sv
530 Bering Strait volume transport and the southward 0.7 Sv volume transport at 24°S
531 provides an indirect 0.1 Sv estimate for freshwater divergence between Bering Strait and
532 24°S . Figure 4, adapted from Piecuch and Ponte (2012), compares hydrographic
533 estimates of meridional heat transport, following Hall and Bryden (1982), within the
534 Atlantic Ocean. The estimate from this study is added (marked with a red point),
535 calculated by adjusting the inverse model solution to yield zero net mass transport along
536 24°S by adding an additional barotropic velocity. The estimate of $0.40 \pm 0.08 \text{ PW}$ out of
537 the box is within the range of the anticipated heat transport across 24°S .

538

539 In order to assess the overturning circulation, each of the 21 γ^n layers (Table 4) is
540 grouped, depending on flow direction. The circulation consists of 0.8 ± 4 Sv of southward
541 flowing surface water (layers 1-2), as a result of the Ekman transport, 15.8 ± 3 Sv
542 northward flow of upper ocean water (layer 3-12), 20.2 ± 2 Sv southward flow of deep
543 water (layers 13-18) and 4.6 ± 1 Sv northward flow of lower LCDW and AABW (layers
544 19-21). The MOC strength is estimated as the 20.2 Sv southward flow of deep water,
545 comparable with the previous estimates in Table 5.

546

547 **4.1.1.3 30°E**

548 On the 30°E section north of 34°S, strong westward flow of warm, salty Indian Ocean
549 water close to the continental slope results in a total westward transport of 65.7 Sv
550 (Figure 3), similar to findings by Casal et al. (2009). Between ~ 34 - 35° S, westward
551 transport is interrupted by eastward flow. The maximum westward flow is 84.5 ± 2.0 Sv
552 for the Agulhas Current. The Agulhas Return Current is attributed to the net eastward
553 flow south of $\sim 36.25^\circ$ S, occupying a broader meridional extent compared to the Agulhas
554 Current. The Agulhas Return Current transport is estimated as 82.2 ± 2.0 Sv, extending
555 between 36.25° S and the Subtropical Front (42.9° S). Above 3.5° C, a 10.7 ± 1.3 Sv
556 Agulhas leakage is detected, comparable with an estimate of 15 Sv from observations
557 using subsurface floats and surface drifters (Richardson, 2007).

558

559 South of the Subtropical Front (STF), the net eastward transport of 131.7 Sv is dominated
560 by the ACC. This estimate is lower than the previous estimates of 160 Sv (full 30°E
561 section, Park et al. (2001)), 147 ± 10 Sv (STF to SACCF between 0° E and 30° E, Legeais
562 et al. (2005)), 136 Sv to 153 Sv for baroclinic and total transport (north of 54.75° S
563 between 0° E and 20° E, (Gladyshev et al., 2008)) and 141.6 ± 2.9 Sv along 30° E (Naveira
564 Garabato et al., 2014). The estimate is closer to the Drake Passage volume transport, as a
565 consequence of constraining the salinity transport around the box boundary. Significant
566 westward flow of AABW is predominately associated with the westward-flowing
567 southern limb of the Weddell Gyre, as previously observed by Schröder and Fahrbach
568 (1999), Park et al. (2001) and Jullion et al., (2014) along the Antarctic continent at 0° E
569 and 30° E.

570

571 **4.1.2 Non-geostrophic terms**

572 **4.1.2.1 Diapycnal transfer of volume, freshwater and heat in the ocean**

573 **interior**

574 A positive diapycnal volume flux represents an upward diapycnal transfer from a denser
575 neutral density class to a lighter neutral density class. In this study, the net diapycnal
576 velocities and volume fluxes (Figure 6a-b) indicate that diapycnal transfer is primarily
577 within the denser layers with nearly zero diapycnal volume flux for layer 10 and above.
578 The vertical structure becomes more significant within the UCDW layer with a tendency
579 for positive fluxes of up to 1 Sv suggesting diapycnal upwelling, including for NADW
580 defined as at the UCDW/LCDW interface ($27.90 < \sigma_\theta < 28.10$), equal to layers 16 and 17
581 (Table 4). The lighter LCDW layer also upwells (4.9 Sv), whilst the most significant
582 downwelling signal of 2 Sv contributes to the production of the densest LCDW layer.
583 The production within this layer is furthered by significant upwelling of 6.3 Sv of AABW
584 layer to LCDW layer with a diapycnal velocity of $\sim 1.5 \times 10^{-5} \text{ m s}^{-1}$. The rough
585 topography of the Scotia Sea (Heywood et al., 2002; Naveira Garabato et al., 2004), and
586 deep passages, such as Vema Channel (Morris et al., 2001), potentially contribute to the
587 significant upwelling and mixing of the AABW and LCDW layers. The absence of large
588 scale diapycnal flux of NADW to lighter neutral density classes supports the findings of
589 Sloyan and Rintoul (2001) for deep to intermediate water conversion in the Southern
590 Ocean to occur along isopycnals, rather than by uniform interior upwelling as suggested
591 in historical conceptual models (e.g Munk, 1966; Gordon, 1986).

592

593 Upward diapycnal salinity flux (Figure 6d) from the SAMW layer towards the fresher
594 surface water and downward diapycnal salinity flux towards the AAIW layer implies a
595 divergence of salinity from the SAMW layer. The SAMW layer is relatively salty in
596 comparison to the waters above and below. This salty SAMW signature is consistent with
597 SAMW sourced from the inflow of salty Indian Ocean water south of Africa, as opposed
598 to fresher SAMW through Drake Passage, in agreement with Sloyan and Rintoul (2000).
599 A similar, if smaller divergence of the salinity flux is observed for the NADW layer at
600 the UCDW/LCDW boundary. Upwelling of salinity to lighter UCDW layers, and

601 downwelling to denser LCDW layers, contributes to the erosion of the NADW salinity
602 maximum.

603

604 Diapycnal temperature velocities (Figure 6f) greater than 0.1 m s^{-1} are only found within
605 the LCDW and AABW layers. For the temperature fluxes, the contribution from
606 diffusion results in the upwelling of temperature flux from denser to lighter LCDW
607 layers, and the downward mixing of LCDW temperature flux to the AABW layer.

608

609 **4.1.2.2 Diapycnal transfer of volume, freshwater and heat by Air-Sea**

610 **interaction**

611 Air-sea interaction contributes to the formation of 14.6 Sv of SAMW through the AAIW
612 to SAMW flux in Figure 7a. This matches (despite the difference in area) the 14 Sv
613 estimate of Sloyan and Rintoul (2001) for their South Atlantic box, nominally bounded
614 by transects at Drake Passage, 0°E and $12\text{-}19^\circ\text{S}$. This process is hypothesised to dominate
615 within the southwest Atlantic region, in the vicinity of the energetic Brazil-Malvinas
616 Confluence (BMC) (Jullion et al., 2010a). Convergence of dense surface water/SAMW is
617 approximately compensated by the divergence of deep neutral density classes: UCDW
618 and upper LCDW ($\sim 15.2 \text{ Sv}$; c.f. 8 Sv (Sloyan and Rintoul, 2001b)). Upwelling of lighter
619 deep neutral density classes, primarily UCDW, and transformation to SAMW/AAIW via
620 exposure to wind, heat and freshwater fluxes contributes towards the MOC southern
621 limb.

622

623 Net freshwater flux contributes to volume flux induced by air-sea interaction, and reflects
624 adjustments to the initial freshwater flux estimate, with extra evaporation required from
625 the surface water layer. As described in Jullion et al. (2010a), freshwater flux is difficult
626 to estimate accurately given uncertainties in upper ocean baroclinic variability and
627 therefore the a priori uncertainties applied to the inverse box model (Ganachaud, 2003;
628 Naveira Garabato et al., 2003).

629

630 Air-sea heat fluxes are dominated by the higher temperature surface ocean within the
631 western South Atlantic basin (Figure 7c). Air-sea heat flux adjustments reach -0.53 PW

632 for radiative heat loss from the warm uppermost surface layer, as it moves northwards
633 towards the North Atlantic Ocean. However, over the water column, the total net
634 adjustment is -0.07 PW for the net air-sea heat flux input estimate of 2.15 PW (65 W m^{-2}
635 over the South Atlantic area) as denser surface layers are heated by the atmosphere.
636 Therefore whilst the whole column adjustment is insignificant, alterations for individual
637 layers show greater significance. Small overall adjustments suggest good agreement
638 between the NOC (v2.0) climatology and observations, despite variability between NOC
639 (v2.0) climatology and alternative heat flux climatology products, particularly in the
640 Southern Ocean (Liu et al., 2011).

641

642 **4.1.2.3 Ekman**

643 The model diagnoses Ekman transport adjustments, assumed meridionally uniform across
644 24°S , in addition to the initial field Ekman transport. Total volume transport adjustment
645 is 0.5 Sv contributing towards the $0.7 \pm 0.3 \text{ Sv}$ freshwater flux. Given uncertainty within
646 the NCEP wind stress (Brunke et al., 2011) used to derive the initial field Ekman
647 transport, the additional transport associated with the Ekman adjustments is only
648 significant within the context of ensuring a net salinity transport of about 26 Sv psu
649 across 24°S .

650

651 **4.2 South Atlantic circulation**

652 Schematic circulation of geostrophic flow within the South Atlantic is shown in Figure 8
653 for the upper and deep ocean neutral density classes. Conversion of the AAIW layer to
654 surface water and SAMW layers occurs between Drake Passage and the 30°E ACC
655 regime. Accumulation within the LCDW layer between Drake Passage and 30°E is offset
656 by AABW layer inflow, as part of the Weddell Gyre southern limb. These results also
657 suggest that the entrainment of the AABW layer into the Circumpolar Deep Water layer
658 is more significant than the intermediate to deep water conversion based on the
659 convergence of the AABW layer at the box boundary.

660

661 Within the subtropics, surface water and SAMW entering the South Atlantic through the
662 Agulhas regime is entrained at the South Atlantic Current/Benguela Current transition,

663 and joins the northward pathway for Agulhas-sourced upper ocean water across 24°S.
664 Given a southward flow of 18.1 Sv of deep water (UCDW and LCDW) across 24°S, the
665 eastward flow of 5.9 Sv of deep water across the Agulhas regime proportionally accounts
666 for approximately one-third of the deep water exiting the South Atlantic that entered the
667 South Atlantic across 24°S. The remainder of the deep water flows into the Southern
668 Ocean and contributes to both Circumpolar Deep Water, and the MOC southern limb.

669

670 **4.2.1 North Atlantic Deep Water layer circulation**

671 For the NADW layer (Figure 9), defined as $27.90 < \gamma^{\rho} < 28.10$, the box-wide circulation is
672 as follows. A net excess inflow from the sum of the box boundary transports requires the
673 divergence of 7.5 Sv from the NADW layer, predominately by upwelling to lighter
674 neutral density classes. This broadly matches the estimate of diapycnal fluxes induced by
675 air-sea interaction of 7.3 Sv from Figure 7a for the NADW layer (layers 16, 17).

676

677 **4.2.2 Antarctic Bottom Water sources and recirculation**

678 A significant source of AABW formation at the Cape Darnley polynya (65°E - 69°E)
679 (Meijers et al., 2010, Ohshima et al., 2013) contributes to full depth cumulative transport
680 of 15.9 ± 2.3 Sv (Figure 3) for the westward flowing, Weddell Gyre southern limb, south
681 of 64.25°S (Naveira Garabato et al., 2014, 2002a). This is largely comprised of LCDW
682 (6.3 ± 1 Sv) and AABW (8.8 ± 0.5 Sv), and comparable to the 24 ± 4 Sv flow associated with
683 the Antarctic Slope Front by Jullion et al. (2014) or 9.6 ± 2.3 Sv Antarctic Slope Front
684 estimate by Dong et al. (2016). Within the Weddell Sea, LCDW and AABW are
685 modified and subsequently exported northward, with wind-forcing thought to dominate
686 this process (Gordon et al., 2010; Jullion et al., 2010b; Wang et al., 2012).
687 Comparatively, the recirculating northern limb of the Weddell Gyre shows a much
688 weaker eastward AABW layer flow across 30°E (Figure 3). The difference between the
689 8.8 ± 0.5 Sv inflow of the AABW layer, as part of the Weddell Gyre southern limb, and
690 the smaller AABW layer outflow across 24°S of 2.6 ± 0.5 Sv is, at least, partially offset by
691 6.3 ± 1.0 Sv of diapycnal upwelling to the densest LCDW layers. This contributes to a
692 6.7 ± 2.2 Sv northward flow, below the 2 °C isotherm, west of the Mid-Atlantic Ridge

693 (10°W), whilst the remaining AABW layer is hypothesised to recirculate within the
694 South Atlantic box.

695

696 **4.2.3 Overturning and gyre circulation for heat and freshwater**

697 The overturning component (Table 6a) is particularly sensitive to the Ekman transport
698 (assumed uniform across the section initially) and initial constraints on the Brazil Current
699 transport. Both components of the total heat transport were similar to those estimated by
700 Bryden et al. (2011).

701

702 The M_{ov} estimates (Table 6b, Figure 5a) are similar to Bryden et al. (2011) and indicate a
703 net southward freshwater transport. Positive M_{az} in this study and Bryden et al. (2011),
704 corresponds with the gyre and the flow near the boundaries transporting freshwater out of
705 the South Atlantic box (Figure 5b).

706

707 **4.3 Anthropogenic Carbon**

708 **4.3.1 Distributions**

709 The Drake Passage C^{ant} distributions in Figure 10 are calculated using the ΔC^* method,
710 with the C^{ant} transports in section 4.3.2 all calculated using the 2009 transect. This
711 transect indicates C^{ant} concentrations markedly shallow from north to south, partly
712 following the general trend of the neutral density isopycnals. The transect maximum of
713 $>30 \mu\text{mol kg}^{-1}$ is primarily within surface, SAMW and AAIW neutral density classes
714 with negligible C^{ant} for the AABW neutral density class. Across 30°E (Figure 11), higher
715 concentrations ($>25 \mu\text{mol kg}^{-1}$) are either found within the Agulhas regime down to 1000
716 dbar or within the upper 200 dbar, south of the Agulhas regime. C^{ant} transports in this
717 study, all make use of the 2008 transect across 30°E. Across 24°S (Figure 12), lower
718 concentrations ($<10 \mu\text{mol kg}^{-1}$) are predominately below 1000 dbar.

719

720 **4.3.2 Transports**

721 Total C^{ant} fluxes of individual neutral density classes are controlled by the underlying
722 volume transport. Net imports of C^{ant} into the South Atlantic box occur only through
723 Drake Passage (Table 7, Figure 13). Across 24°S, although total net DIC flow is

724 southward (Gruber et al., 2009), the large surface-to-deep C^{ant} gradient causes a net
725 northward transport, in line with previous estimates (Holfort et al., 1998). A C^{ant}
726 overturning estimate of 0.11 PgC yr^{-1} , associated with the 20.2 Sv overturning, is
727 calculated based on the southward transport of C^{ant} -poor deep water (layers 13-18
728 following Section 4.1.1.2). A net eastward C^{ant} transport within the Agulhas regime is
729 caused by ventilation within the highly energetic South Atlantic sector of the Agulhas
730 regime and C^{ant} increase in the upper ocean, prior to the eastward return flow.

731

732 The mean transport-weighted (TW) C^{ant} is calculated for each neutral density class at the
733 box boundaries by dividing the total C^{ant} transport by the total volume transport (Table 8).
734 Transport-weighted values are most heavily weighted at the location of the transport
735 maximum, and hence are directly influenced by changes in the transport profile (Georgi
736 and Toole, 1982; Tillinger and Gordon, 2010). Neutral density classes with the largest
737 volume transports along both Drake Passage and the 30°E ACC regime, particularly
738 UCDW and LCDW layers (Figure 8), therefore contribute significantly to the observed
739 C^{ant} divergence (Figure 13). Small systematic biases within these low C^{ant} waters, below
740 the level of the adjustments calculated as part of GLODAPv2, could contribute towards
741 the significant C^{ant} divergence. The C^{ant} divergence shown by larger eastward-flowing
742 TW C^{ant} at 30°E, compared to either eastward-flowing TW C^{ant} at Drake Passage or
743 southward-flowing C^{ant} at 24°S is suggestive of an air-sea C^{ant} input requirement.

744

745 Storage rate is calculated using MPD estimates from Drake Passage, 24°S and 30°E
746 multiplied by the mean rate of C^{ant} increase in the mixed layer (Table 9). As listed in
747 Table 9, MPD from Drake Passage and 30°E are notably shallower than 770m for the
748 region between 10°S and 30°S (Holfort et al., 1998), and 790 m at 24.5°N (Rosón et al.,
749 2003). As described in Álvarez et al. (2003), areas with higher stratification yield
750 shallower MPD, with comparatively lower penetration of C^{ant} below the upper 2000 dbar
751 at Drake Passage, compared to 30°E, resulting in the shallower MPD. Increased
752 convection, therefore leads to increased uncertainty in the time variability of the MPD
753 (Khatiwala et al., 2013; Pérez et al., 2008).

754

755 Storage rates of $0.22\pm 0.29 \text{ mol m}^{-2} \text{ yr}^{-1}$ along Drake Passage, $0.81\pm 0.53 \text{ mol m}^{-2} \text{ yr}^{-1}$
756 along 24°S and $0.29\pm 0.18 \text{ mol m}^{-2} \text{ yr}^{-1}$ along 30°E extend the range of previous South
757 Atlantic storage rate estimates from repeat hydrography (Table 10). The values show
758 similarities with the time-averaged Green's Function Inversion in Khatiwala et al. (2013)
759 (their Figure 7). The Drake Passage estimate reflects its shallower MPD of 259.8 m than
760 for other parts of the Southern Ocean, given that less C^{ant} has penetrated into deeper
761 neutral density classes based on the lower TW C^{ant} estimates for UCDW, LCDW and
762 AABW (Table 8). Along 30°E , the C^{ant} values are normalised by temperature to remove
763 biases caused by cooler temperatures within the mixed layer in the 2008 occupation
764 compared to the 1996 occupation. The temperature normalisation reduced the initial high
765 $\Delta C_{\text{ml}}^{\text{ant}}$ estimate of $1.52 \mu\text{mol kg}^{-1} \text{ yr}^{-1}$ along 30°E to $0.45 \mu\text{mol kg}^{-1} \text{ yr}^{-1}$. The $0.45 \mu\text{mol}$
766 $\text{kg}^{-1} \text{ yr}^{-1}$ estimate is at the lower range of previous South Atlantic estimates of CO_2 uptake
767 ($0.6\text{-}1.0 \mu\text{mol kg}^{-1} \text{ yr}^{-1}$) (Murata et al., 2008; Peng and Wanninkhof, 2010; van Heuven,
768 2013). The 24°S estimate is similar to Holfort et al. (1998)'s estimate of $0.59\pm 0.12 \mu\text{mol}$
769 $\text{kg}^{-1} \text{ yr}^{-1}$ for the 10°S and 30°S region and within their 20% uncertainty estimate.

770

771 For the South Atlantic box, the mean storage rate for Drake Passage, 24°S and 30°E
772 (Table 10), calculated from the mean MPD, mean $\Delta C_{\text{ml}}^{\text{ant}}$ and mean ρ_{ml} (Table 9) and
773 integrated over the ocean surface area (estimated as $3.3 \times 10^{13} \text{ m}^2$ assuming a
774 parallelepiped ocean) yields a basin-wide C^{ant} storage of $0.18\pm 0.12 \text{ Pg C yr}^{-1}$. Application
775 of the TSR-based C^{ant} storage estimation method, which makes use of additional
776 historical hydrographic cruise data from the interior of the South Atlantic Ocean sector,
777 generates a storage term of $0.21\pm 0.06 \text{ Pg C yr}^{-1}$ (van Heuven, S. (2016), manuscript in
778 preparation). The two estimates compare well despite substantially different
779 methodologies. The smaller TSR uncertainty represents its greater robustness as a
780 calculation approach, due to the additional data and lack of structural assumptions
781 compared with the MPD method (transient steady state, parallelepiped ocean). Historical
782 storage estimates for the South Atlantic regions show slightly higher values: 0.30 Pg C
783 yr^{-1} between $2^{\circ}\text{S}\text{-}58^{\circ}\text{S}$ based on decadal hydrographic observations (Peng and
784 Wanninkhof, 2010) and $0.29 \text{ Pg C yr}^{-1}$ between $0^{\circ}\text{S}\text{-}58^{\circ}\text{S}$ from multiple global ocean
785 inversions based on hydrographic section data (Mikaloff Fletcher et al., 2006). Based on

786 this study, usage of MPD calculations appear to have some value in providing a
787 reasonable estimate for C^{ant} storage in the absence of full basin-scale historical data.
788 However, greater uncertainty will be assigned to estimates if the sampling pattern of the
789 hydrographic cruises chosen does not fully capture the north-south variability within the
790 Southern Ocean of the column inventory of ΔC^{ant} (see Figure 7.13 from van Heuven,
791 (2013)). Similarly, MPD calculations are also dependent upon the shape of the C^{ant}
792 profile, such that the presence of increasing amounts of C^{ant} within bottom water layers
793 (due to proximity to bottom water ventilation locations) may compromise the MPD
794 assumption (Khaliwala et al., 2013; Pérez et al., 2008). However, the sections used here
795 are not thought to suffer from this at this stage, with negligible bottom-water C^{ant} change
796 identified (Evans, 2013).

797

798 **4.3.3 Anthropogenic CO₂ air-sea flux**

799 The C^{ant} budget for the South Atlantic box - comprising storage and divergent flux terms
800 at the box boundaries (Figure 13) - is balanced by a $0.51 \pm 0.37 \text{ Pg C yr}^{-1}$ air-sea flux term.
801 This compares to a global anthropogenic CO₂ uptake of 2.2 to $2.6 \pm 0.3 \text{ Pg C yr}^{-1}$
802 estimated from ocean inverse and biogeochemical models (DeVries, 2014; Gruber et al.,
803 2009), or more generally 2 Pg C yr^{-1} from a range of oceanic and atmospheric
804 observations (Wanninkhof et al., 2013). The Southern Ocean is the largest annual sink
805 region of total (natural and anthropogenic) CO₂ of more than $0.42 \text{ Pg C yr}^{-1}$ south of 44°S
806 (Lenton et al., 2013). Regional observations and model outputs for its Atlantic sector
807 combined within the South Atlantic from $18\text{-}58^\circ\text{S}$, broadly similar to our South Atlantic
808 box but excluding the small sea-air CO₂ flux south 58°S (Lenton et al., 2013; van
809 Heuven, 2013), suggest a net annual mean total (natural and anthropogenic) CO₂ flux of
810 $0.19 - 0.38 \text{ Pg C yr}^{-1}$ (Lenton et al., 2013; Schuster et al., 2013). This is smaller than the
811 air-sea uptake estimate derived here that only quantifies the anthropogenic component.
812 However, large outgassing of natural carbon identified in the Southern Ocean (Mikaloff
813 Fletcher et al., 2007) suggests that any estimates of regional CO₂ uptake here will be
814 disproportionately of anthropogenic origin. A global ocean circulation inverse model
815 assimilating potential temperature, salinity, CFC-11 and radiocarbon observations
816 (DeVries, 2014) supports the distinction between natural and anthropogenic CO₂ uptake,

817 with an estimated total (natural and anthropogenic) CO₂ uptake for the South Atlantic
818 box of 0.43 Pg C yr⁻¹ of which 0.38 Pg C yr⁻¹ is anthropogenic CO₂. Although the air-sea
819 C^{ant} uptake estimate here is larger than other observational and model estimates this is not
820 entirely unexpected, as a seasonal bias may exist in the input C^{ant} estimates due to the
821 austral summer-based cruise timings: increased stratification and intense biological
822 production draw down surface carbon levels and increase the air-sea ΔpCO₂ difference.
823 Combined with a temperature-related increase in the Revelle factor (Sabine et al., 2004)
824 that enables greater anthropogenic carbon loadings, the associated uptake reaches its
825 maximum during the summer months and is a likely major contributor to the large budget
826 residual.

827

828 Differences from alternative estimates may also be partially methodological in nature.
829 Given the large volume transports associated with the UCDW and LCDW neutral density
830 classes in this study, systematic biases within these deep waters could potentially
831 contribute to large differences in C^{ant} between Drake Passage and 30°E, which are
832 inferred as being balanced by the air-sea flux. The differences between volume transport-
833 weighted C^{ant} estimates at Drake Passage and 30°E (Table 8) also imply that these deeper
834 neutral density classes must be gaining C^{ant} within the South Atlantic. Khatiwala et al.,
835 (2013) describe a key difference between the ‘ocean inversion’ method, where
836 hydrographic section estimates of C^{ant} are combined with Ocean General Circulation
837 Models (OGCMs), first applied in Gloor et al. (2003) and later in further depth in
838 Mikaloff Fletcher et al. (2006, 2007), Gruber et al. (2009) and Khatiwala et al., (2013),
839 and C^{ant} flux estimates from ship transects. Hydrographic occupations are accurate for a
840 single point in time and thus subject to sampling biases, whilst the ocean inversion
841 method represents a transport integrated in time since the industrial revolution, and
842 typically scaled to any selected year (e.g. 1995 in Mikaloff Fletcher et al. (2006)).
843 Additionally, seasonal variability affects hydrographic fluxes (Wilkin et al., 1995) with
844 Lachkar et al. (2009) suggesting that subtropical South Atlantic seasonal variability
845 corresponds to up to 20% of the annual mean transport of C^{ant}. The inverse model in the
846 current study is designed to create a 2008-2009 ocean mean such that the calculated
847 divergence within the South Atlantic Ocean is representative of that time period.

848

849 **5 Conclusions**

850 An inverse box model was used to examine net exchange between the South Atlantic
851 Ocean and surrounding basins, inspired by the work of Rintoul (1991). We revisit this
852 study with newer data and the inclusion of C^{ant} . The key findings include:

- 853 • The 15.9 Sv of westward Weddell Gyre return flow at 30°E contains 8.8 ± 0.5 Sv of
854 the AABW layer, contributing to a net 13.8 ± 1.0 Sv inflow of the AABW layer to the
855 box across all sections. Diapycnal upwelling of 6.3 ± 1.0 Sv from the AABW layer to
856 the LCDW layer within the box, leads to a net AABW recirculation within the South
857 Atlantic of 7.5 ± 1.4 Sv.
- 858 • A Meridional Overturning Circulation of 20.2 Sv with a net mass transport of 0.7 ± 0.3
859 Sv southward and a freshwater transport associated with the overturning component
860 M_{ov} of 0.09 Sv southward across 24°S. This southward overturning freshwater flux of
861 0.09 Sv supports the notion of MOC bistability.
- 862 • Agulhas leakage, defined as westward flow above the 3.5 °C isotherm, is 10.7 ± 1.7
863 Sv. Total eastward transport of Circumpolar Deep Water is 5.9 ± 2.2 Sv beneath the
864 Agulhas Current system, north of the Subtropical Front. Agulhas leakage contributes
865 towards the northward flowing upper ocean water across 24°S, whilst up to one-third
866 of southward-flowing deep water across 24°S, exits the South Atlantic underneath the
867 net westward-flowing Agulhas leakage.
- 868 • The C^{ant} divergence from the South Atlantic box of 0.33 ± 0.31 Pg C yr⁻¹ and
869 0.18 ± 0.12 Pg C yr⁻¹ of C^{ant} storage correspond to a C^{ant} air-sea uptake of 0.51 ± 0.37 Pg
870 C yr⁻¹. While 0.18 ± 0.12 Pg C yr⁻¹ of anthropogenic carbon is stored within the box,
871 89% of C^{ant} input to the South Atlantic box is exported from the South Atlantic. C^{ant}
872 export from the South Atlantic occurs across both the 24°S section (0.28 ± 0.16 Pg C
873 yr⁻¹), and across 30°E, associated with the 1.04 ± 0.42 Pg C yr⁻¹ ACC and the
874 0.08 ± 0.07 Pg C yr⁻¹ Agulhas Current and its return flow.
- 875 • Significant C^{ant} divergence within the South Atlantic box is only sustainable with
876 significant C^{ant} uptake from the atmosphere. C^{ant} uptake of 0.51 ± 0.37 Pg C yr⁻¹
877 equivalent to approximately 25% of previous estimates of global C^{ant} uptake may be

878 caused through the upwelling of C^{ant} -poor NADW as part of the MOC, which
879 subsequently absorbs atmospheric CO_2 into the ocean surface layers.

880 In conclusion, the South Atlantic circulation diagnosed in this study is characterised by
881 inflow through Drake Passage, overturning south of 24°S consistent with southward-
882 flowing UCDW and LCDW and conversion to lighter neutral density classes through
883 diapycnal processes. Northward flows of surface water, SAMW and AAIW layers merge
884 with a net westward Agulhas leakage from the Agulhas system to complete the MOC
885 upper cell. AAIW, UCDW and LCDW flow eastward below the Agulhas system, whilst
886 further south, eastward transport in the ACC dominates. Near the Antarctic continental
887 margin, a westward flow supplies AABW to the Weddell Sea.

888

889 Ventilation and transformation within the Weddell Sea precedes the northward flow of
890 the renewed AABW layer out of the Weddell Sea, whereupon significant diapycnal
891 processes convert the AABW layer to the LCDW layer, limiting the volume of AABW
892 exiting the South Atlantic. There is net SAMW production, LCDW layer creation and
893 AABW layer destruction in the South Atlantic. For C^{ant} , an imbalance between the
894 transport-weighted inflow and outflow for each neutral density class indicates significant
895 uptake of CO_2 from the atmosphere within the South Atlantic, subsequently supplying the
896 Atlantic Ocean north of 24°S and the Indian sector of the Southern Ocean with C^{ant} . Inter-
897 basin exchange within the South Atlantic therefore ventilates CDW, receives, modifies
898 and then consumes AABW, and supplies C^{ant} to the rest of the global ocean.

899

900

901 **Appendix A**

902 **Constraint weighting**

903 Each constraint has an associated uncertainty. As each constraint is represented by a row
904 in **E**, each row is weighted according to the constraint's uncertainty. For the layer volume
905 constraints, larger a priori uncertainties (ϵ_j) are applied to the upper ocean than the deep
906 ocean following Ganachaud (2003) for the neutral density classes: Surface ($\pm 4\text{Sv}$),
907 SAMW ($\pm 4\text{Sv}$), AAIW ($\pm 3\text{Sv}$), UCDW ($\pm 2\text{Sv}$), LCDW ($\pm 1\text{Sv}$) and AABW ($\pm 0.5\text{Sv}$).
908 For volume transport constraints, the reciprocal of the a priori uncertainty is applied as
909 the row weighting whilst for property transports, the reciprocal of the a priori uncertainty
910 multiplied by 2, and multiplied by the larger of either the property standard deviation or
911 property mean is applied for each layer/row. Typically a property standard deviation is
912 applied, however, the property mean is included to cope with excessively small standard
913 deviation values, and to better weight higher temperature anomalies within the surface
914 waters across the box. For full depth salinity anomaly transport around the box boundary,
915 a small a priori uncertainty (0.2 Sv psu) is applied to better constrain the system, making
916 use of well-constrained values for Drake Passage (Cunningham et al., 2003; Meredith et
917 al., 2011) and 24°S (Coachman and Aagaard, 1988; Woodgate and Aagaard, 2005),
918 following the constraint for full depth boundary salinity transport applied to the initial
919 field. The small uncertainty improves the zero salinity convergence constraint for the
920 inverse box, rather than reflecting actual uncertainty. Only small full-depth residual
921 imbalances for volume of -0.47 Sv and salinity anomaly of -1.08 Sv psu remain after
922 applying the inverse box model.

923 **Weightings for unknown velocities**

924 The accuracy of the depth-independent velocities is affected by the inclusion of a priori
925 uncertainties for weighting each column in **E**, and designed to optimally weight the
926 different components of the solution. Column weighting takes the general form of the a
927 priori uncertainty divided by the appropriate area and subsequently square rooted.

928

929 **Appendix B**

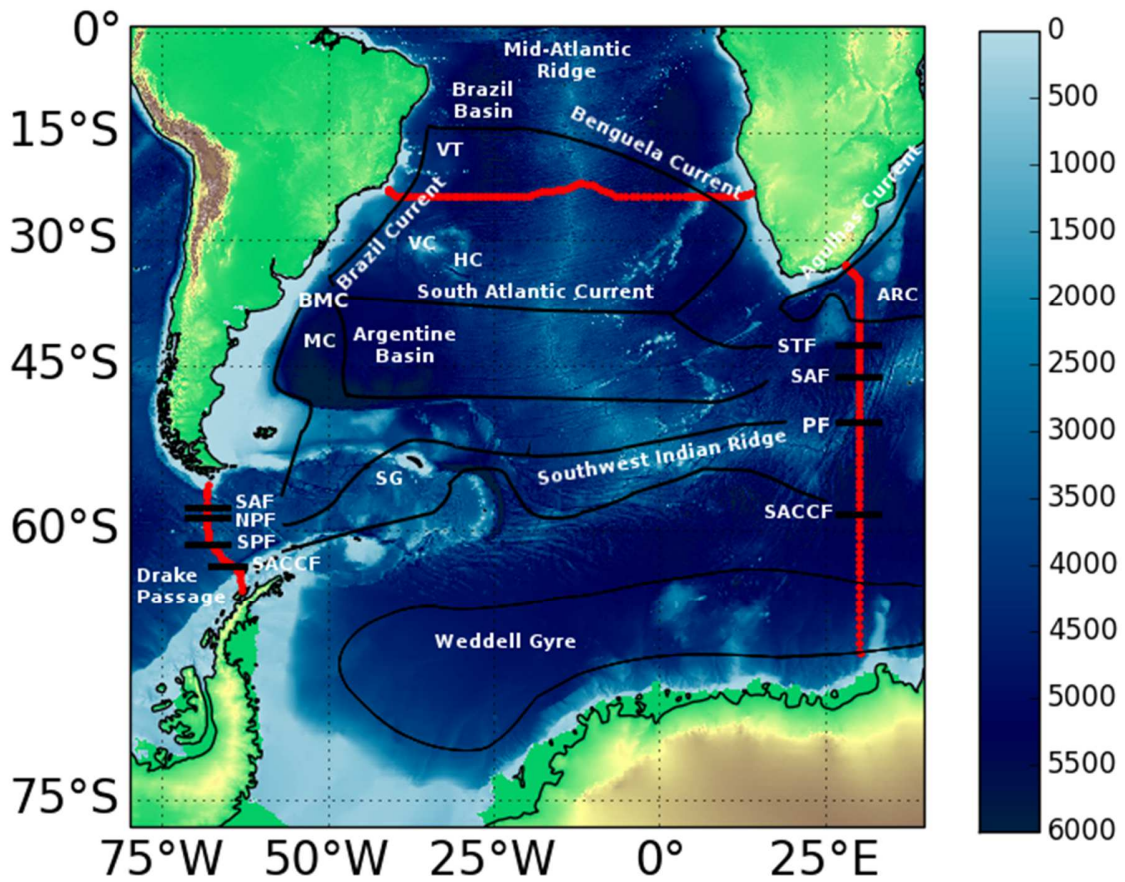
930 Historical surface data from five meridional cruises that intersect the 2009 24°S section
931 across its full extent have been used to generate estimates of the change of anthropogenic
932 carbon within the mixed layer (ΔC_{ml}^{ant}) and thus C^{ant} storage rates across 24°S, as detailed
933 in Sections 2.1 and 4.3. Each meridional cruise provides a single intersection for
934 comparison to the 24°S zonal transect. C^{ant} was calculated in an identical manner to the
935 other box sections. The C^{ant} profile of the nearest station, in terms of latitude and
936 longitudes coordinates, along each of the meridional sections is matched to the nearest
937 station along the 24°S zonal transect to help determine ΔC_{ml}^{ant} . Historical cruises used
938 were as follows: A14 (35A3CITHER3_1) occupying a longitude of 9°W at 24°S between
939 January-February 1995 (Mercier and Arhan, 1995); A13 (35A3CITHER3_2) crossing
940 through 24°S at 8°E between February-April 1995 (Mercier and Arhan, 1995);
941 A15/AR15 (316N142_3) crossing 24°S at 19°W in May 1994 (Smethie and Weatherly,
942 1994); A16 (318HYDROS4) crossing 24°S at 25°W in March 1989 (Talley et al., 1989);
943 and A17 (3230CITHER2_1-2) intersecting 24°S at 33°W in February 1994 (Mémery,
944 1994). Data from each of these cruises is accessible from the Carbon Hydrographic Data
945 Office (CCHDO).

946

947 **Acknowledgments**

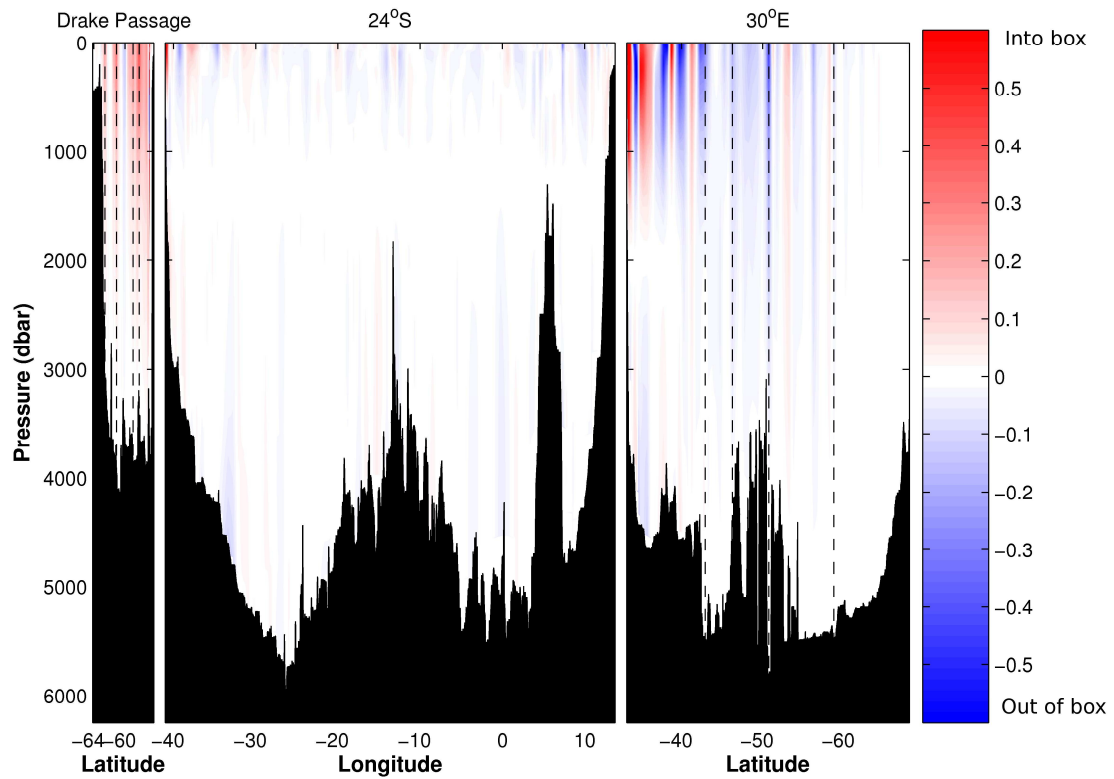
948 GRE thanks the Marine Physics and Ocean Climate group at the National Oceanography
949 Centre, and the Natural Environment Research Council (NERC) for Ph.D. funding. ELM,
950 BAK and the 24°S and Drake Passage Cruises were funded by NERC through
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956 2008 I6S cruise was funded as part of the US repeat hydrography program, with principal
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961 Peltola, Rik Wanninkhof and Mark Stinchcombe requires particular acknowledgment.

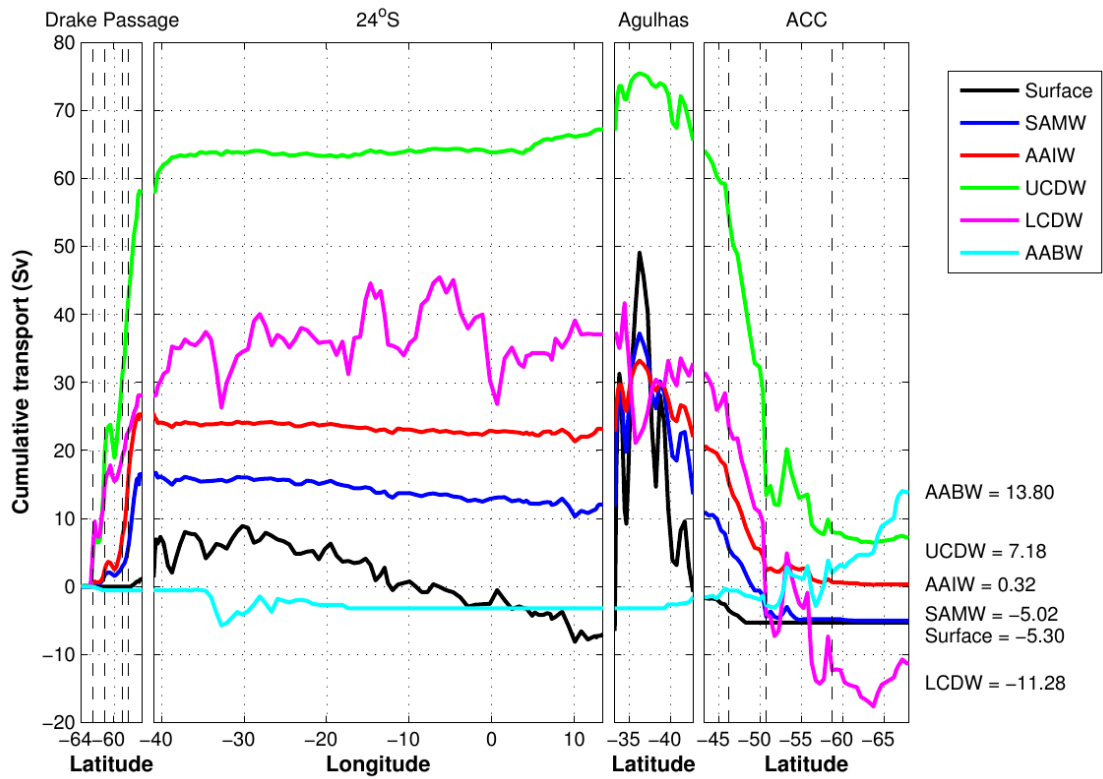


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Figure 1: Map of the hydrographic sections that form the boundaries to the South Atlantic inverse box model. Sections are A21 (Drake Passage), I6S (30°E) and 24°S. The Subtropical Front (STF), Subantarctic Front (SAF), North Polar Front (NPF), South Polar Front (SPF) and Southern Antarctic Circumpolar Current Front (SACCF) are indicated. Major topographical and circulation features are: Vitoria-Trinidad seamounts VT, Vema Channel VC, Hunter Channel HC, Brazil Malvinas Confluence BMC, Malvinas Current MC, South Georgia SG and the Agulhas Return Current ARC.



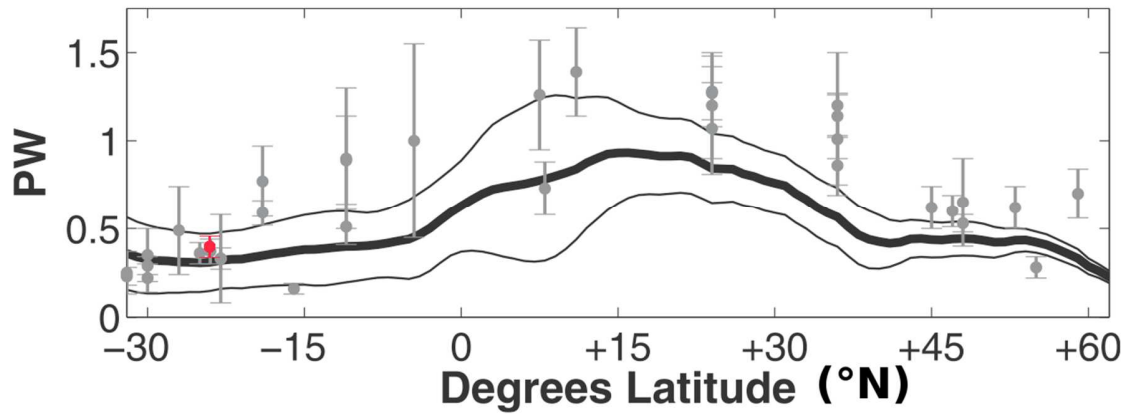
971
 972 Figure 2: Geostrophic velocities (barotropic plus baroclinic velocities from the final
 973 solution) on the box boundary in units of m s^{-1} . Into (out of) the box is shown by red
 974 (blue). The dashed lines indicate frontal positions along the Drake Passage section from
 975 south to north: SACCF, SPF, NPF and SAF, and along the 30°E section from north to
 976 south: STF, SAF, PF, SACCF.
 977



978

979 Figure 3: Cumulative transport along the box boundary for the final solution for each
 980 neutral density class in units of Sv. The total cumulative transport for each neutral density
 981 class is shown. Positive transports refer to a net gain by the box, whilst negative
 982 transports refer to a net loss. Vertical dashed lines indicate fronts.

983



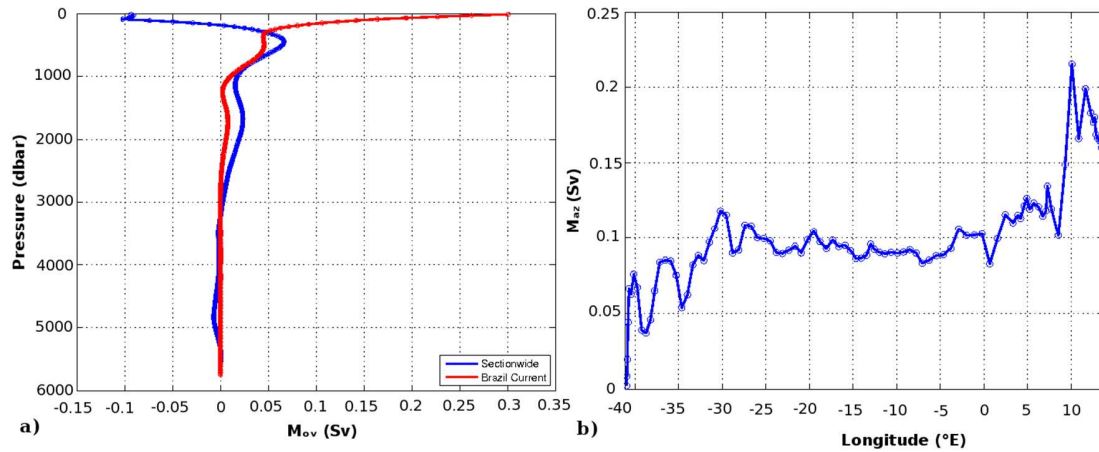
984

985 Figure 4: Heat transport (red; petawatts (PW)) for zero net mass transport across 24°S.
 986 Additional hydrographic estimates and errors (grey bars) are shown together with
 987 meridional heat transport (from Piecuch and Ponte (2012)) with an average time-mean
 988 ECCO (Estimating the Circulation and Climate of the Ocean; black solid thick line)
 989 estimate from model-observation syntheses. The uncertainty interval is given as the
 990 standard deviation of the heat transport time series (black thin lines).

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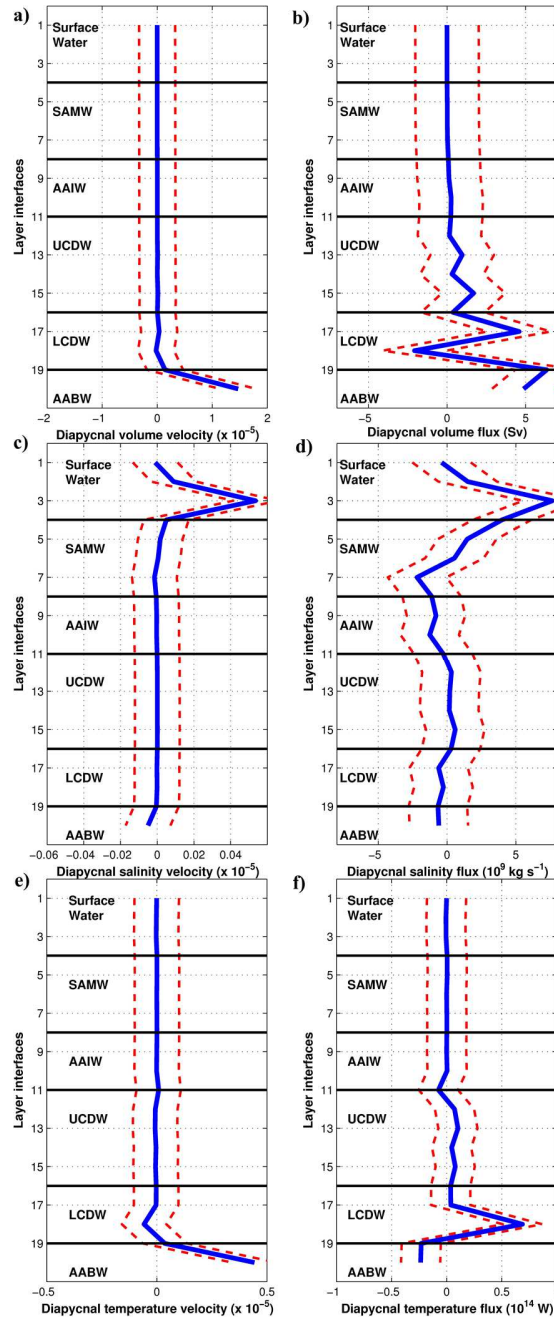
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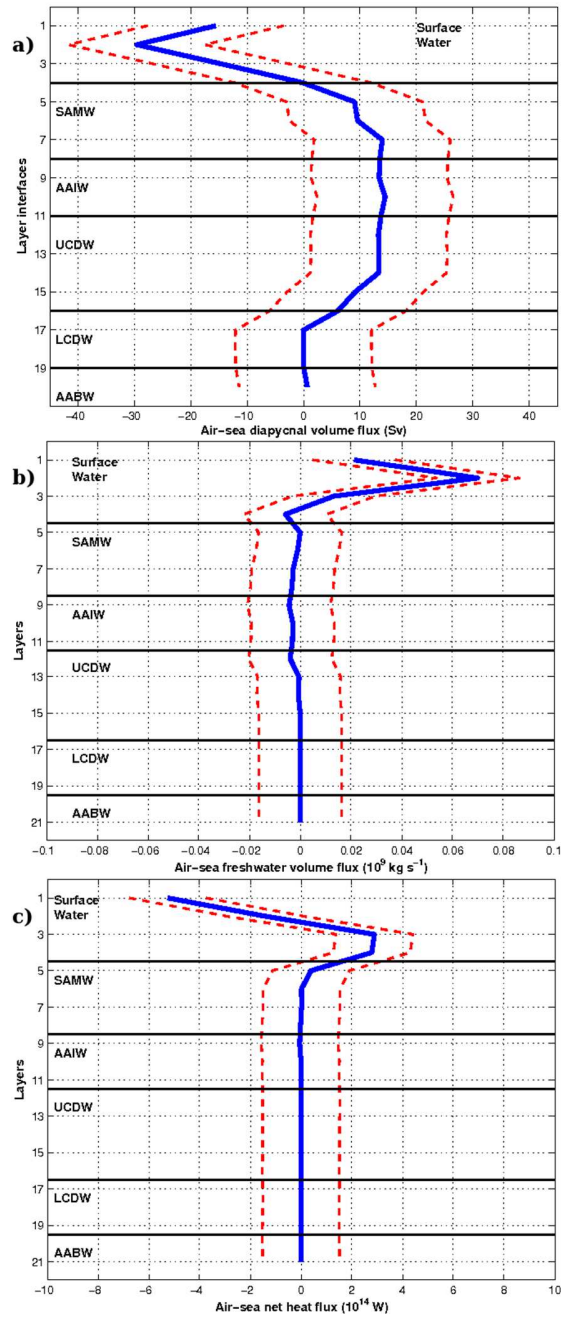
995 Figure 5: a) Cumulative M_{ov} as a function of pressure for the 24 $^{\circ}$ S section (blue), and for
 996 west of 35 $^{\circ}$ W, inclusive of the Brazil Current (red). Positive (negative) M_{ov} is northward
 997 (southward). Units of Sv. b) Cumulative M_{az} as a function of longitude. Positive
 998 (negative) M_{az} is northward (southward). Units of Sv.

999



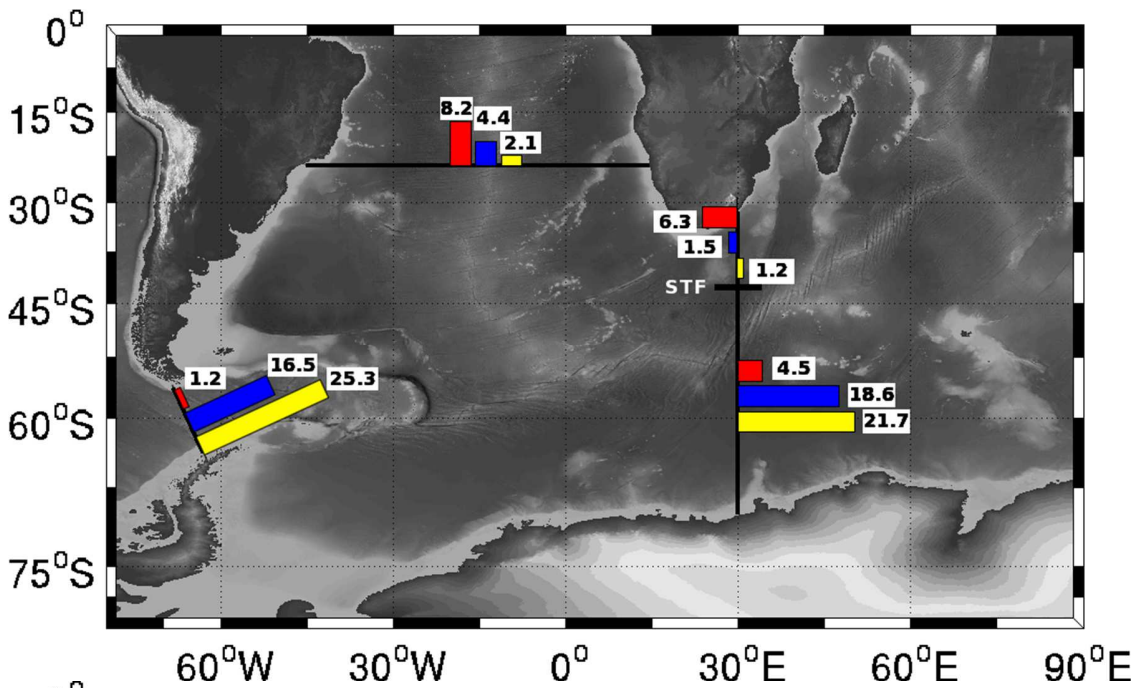
1000

1001 Figure 6: a) Diapycnal volume velocity (m s^{-1}) and b) volume flux (Sv), c) diapycnal
 1002 salinity velocity (m s^{-1}) and d) salinity flux (kg s^{-1}) and e) diapycnal temperature velocity
 1003 (m s^{-1}) and f) temperature flux (W) across each layer interface within the South Atlantic
 1004 box. A positive (negative) velocity or transport represents an upward (downward) flow.
 1005 The dashed lines represent one standard deviation. Neutral density class boundaries are
 1006 marked (solid black line), and neutral density classes labelled.

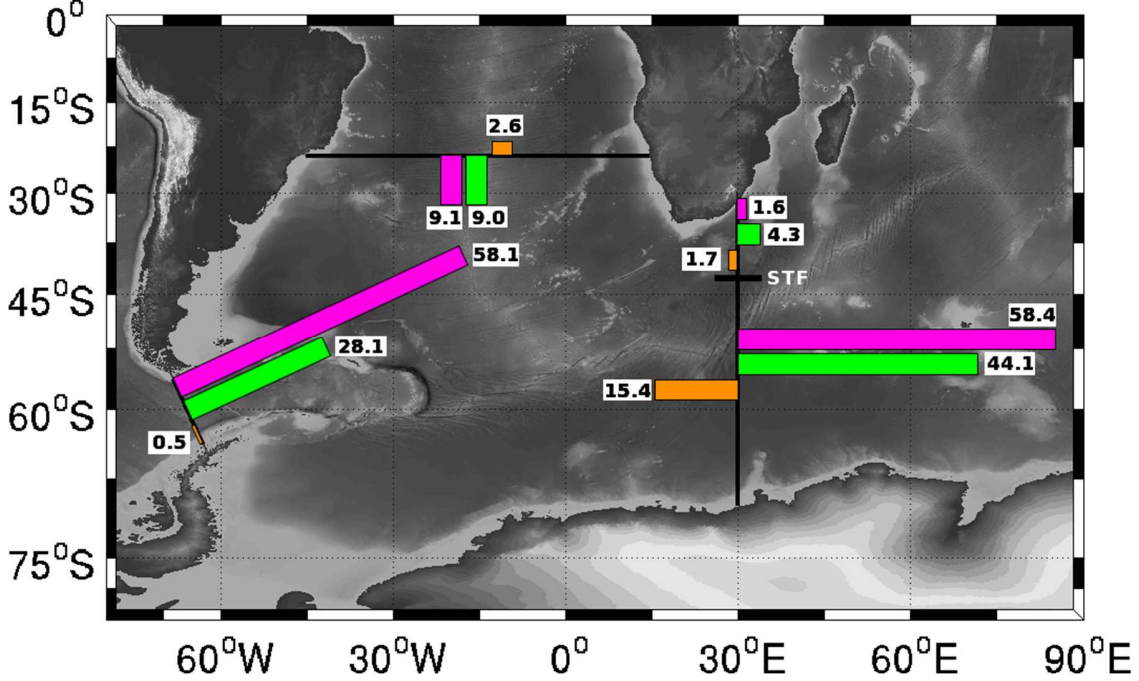


1007

1008 Figure 7: Air-sea interaction induced diapycnal a) volume flux, b) freshwater flux and c)
 1009 heat flux within the South Atlantic box. Diapycnal volume flux is estimated at the layer
 1010 interface, freshwater flux and heat flux induced by air-sea interaction is into each
 1011 individual layer. Positive (negative) values indicate a flux towards lighter (heavier)
 1012 neutral density classes. Neutral density class boundaries are marked (solid black line),
 1013 and neutral density classes labelled.



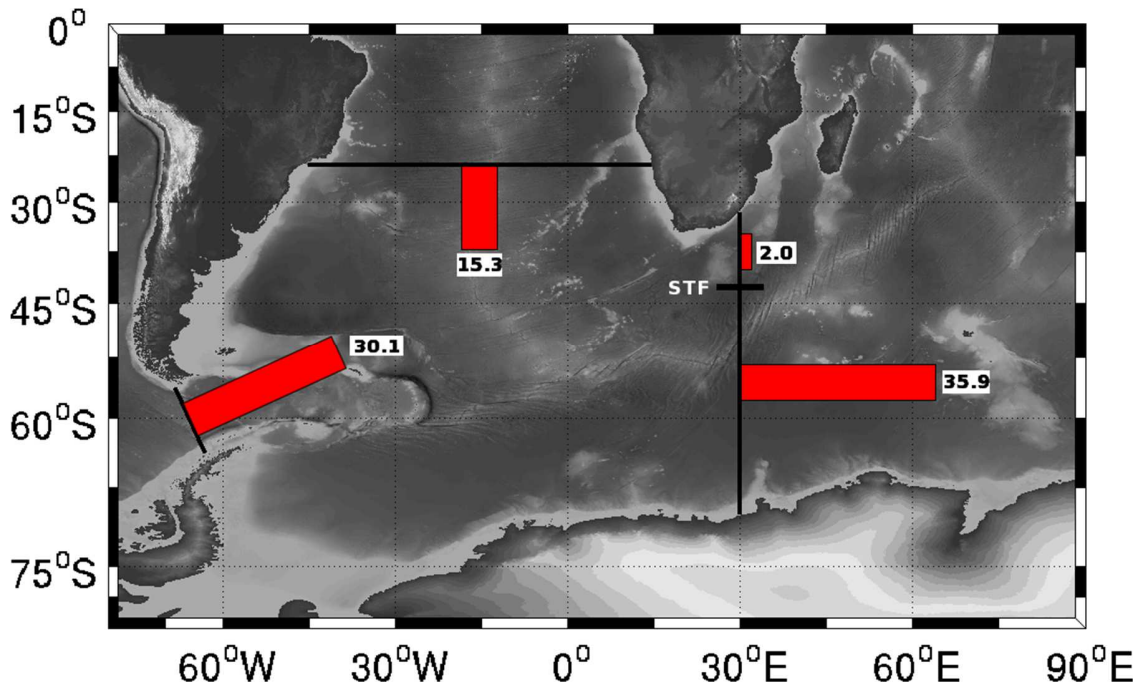
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1016 Figure 8: Schematic circulation for the inverse model solution. The length of each bar is
 1017 proportional to the net transport associated with each neutral density class. Neutral
 1018 density classes shown are a) surface water (red), SAMW (blue), and AAIW (yellow) and
 1019 b) UCDW (pink), LCDW (green) and AABW (orange). Numbers at the end of each bar
 1020 give transports in Sv. A priori uncertainties for transport in each neutral density class

1021 transport are: surface water 4 Sv, SAMW 4 Sv, AAIW 3 Sv, UCDW 2 Sv, LCDW 1 Sv
1022 and AABW 0.5 Sv.



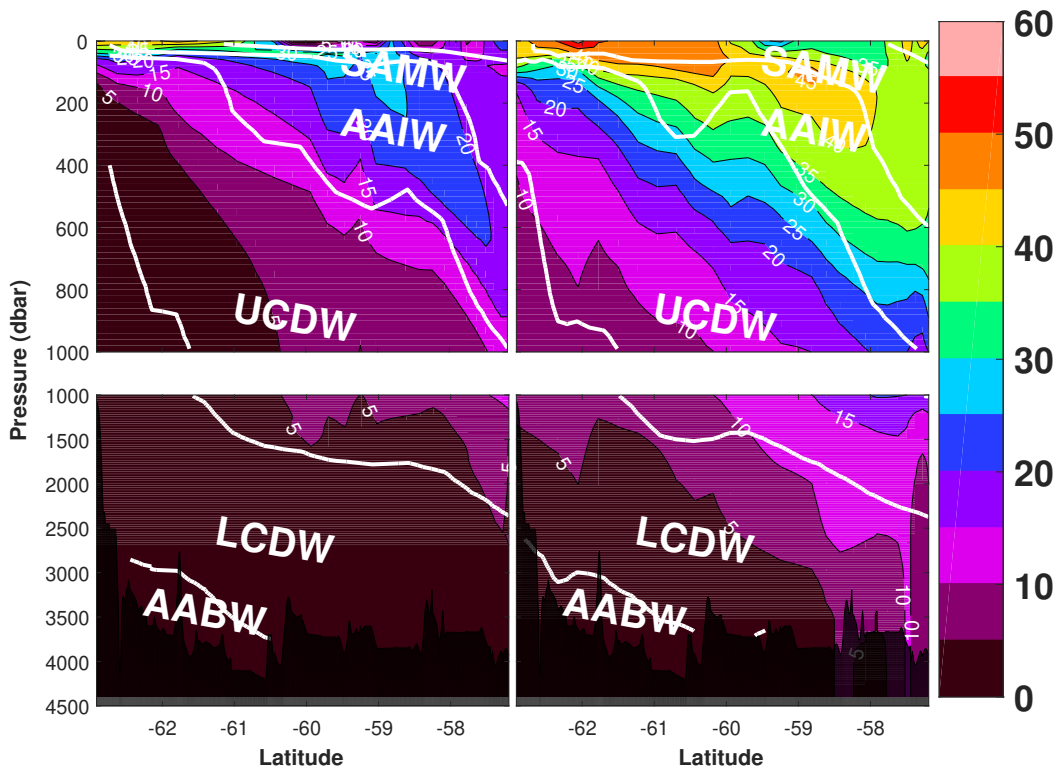
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1024 Figure 9: Schematic circulation for the NADW from the inverse model solution, defined
1025 as at the UCDW/LCDW interface ($27.90 < \gamma^n < 28.10$), equal to layers 16 and 17 (Table 4).
1026 The length of each bar is proportional to the net transport. Numbers at the end of each bar
1027 give transports in Sv with an uncertainty of 2 Sv.

1028

1029

1030



1031

Figure 10: ΔC^* -derived distribution across Drake Passage of C^{ant} for Left: 1990 and Right: 2009. The neutral density:geopotential height interpolation scheme mentioned in Section 2.1 uses a 0.02 geopotential height (φ) grid across Drake Passage. Neutral density classes are labelled following the neutral density interfaces in Table 4. Units of $\mu\text{mol kg}^{-1}$.

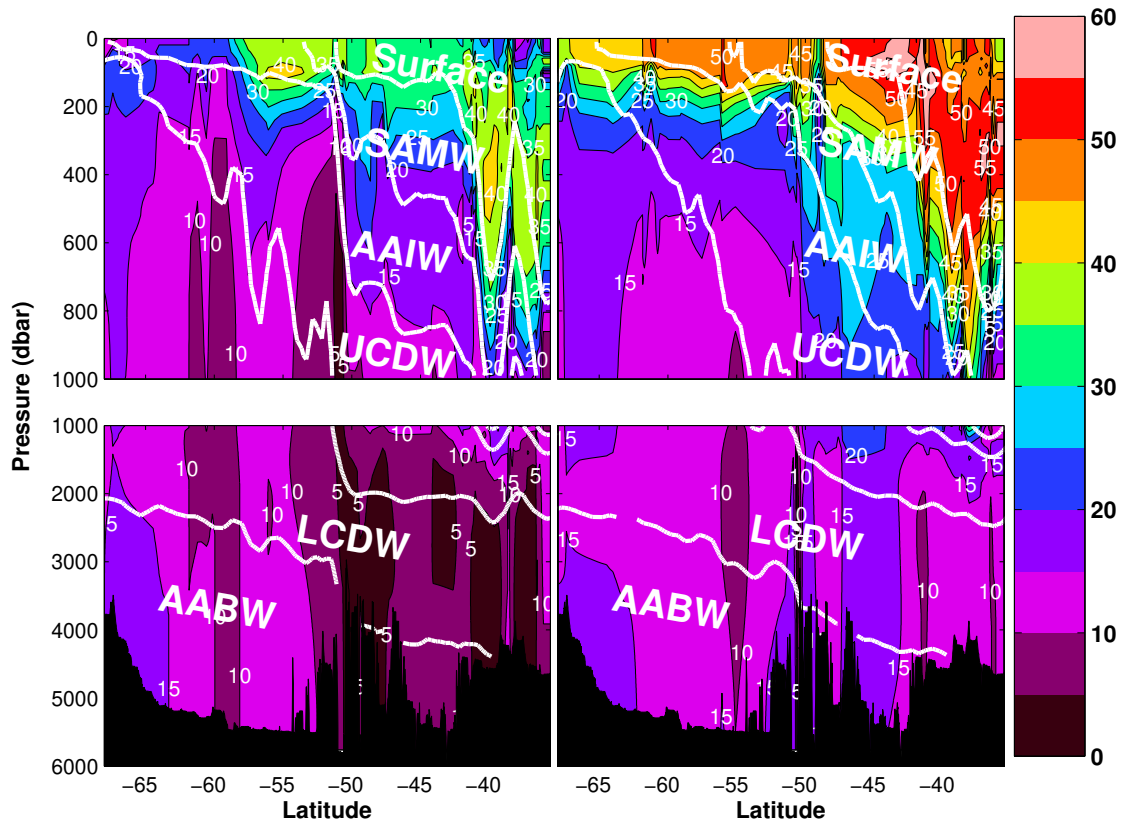


Figure 11: ΔC^* -derived distribution across 30°E of C^{ant} for Left: 1996 and Right: 2008. The neutral density:geopotential height interpolation scheme mentioned in Section 2.1 uses a 0.02 geopotential height (φ) grid across 30°E between 35°S and 58°S and a 0.002 φ grid south of 58°S . Neutral density classes are labelled following the neutral density interfaces in Table 4. Units of $\mu\text{mol kg}^{-1}$.

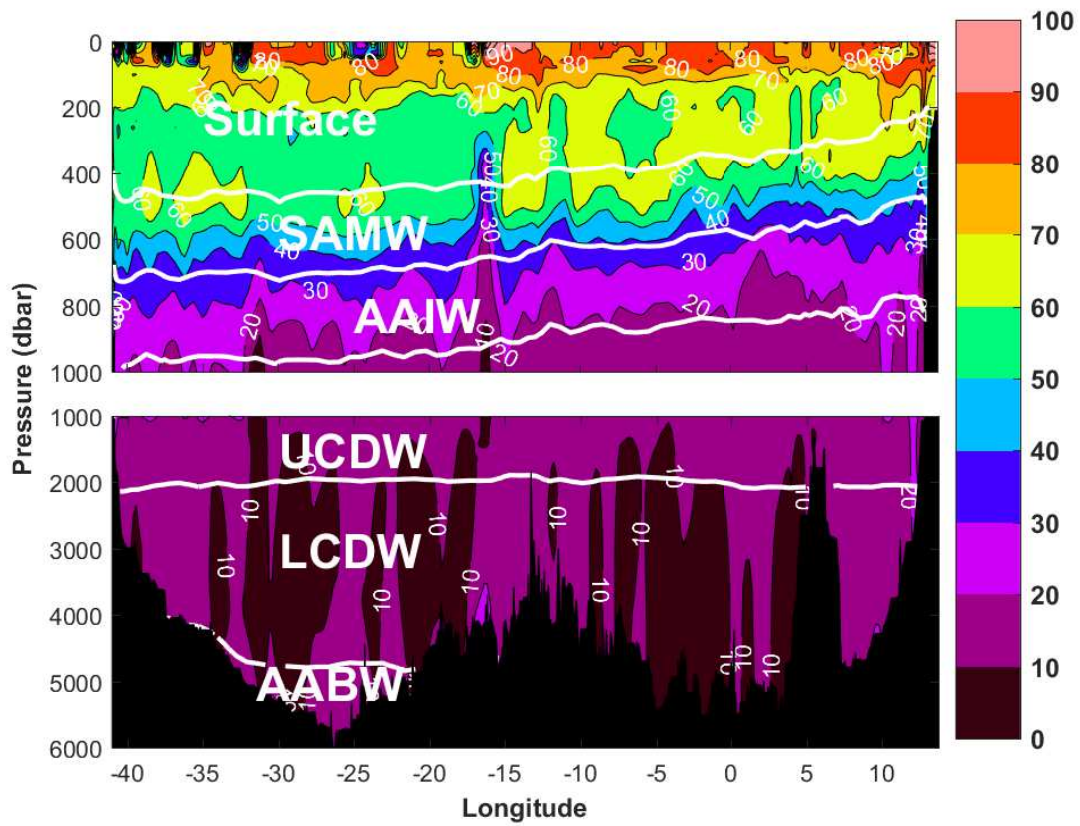
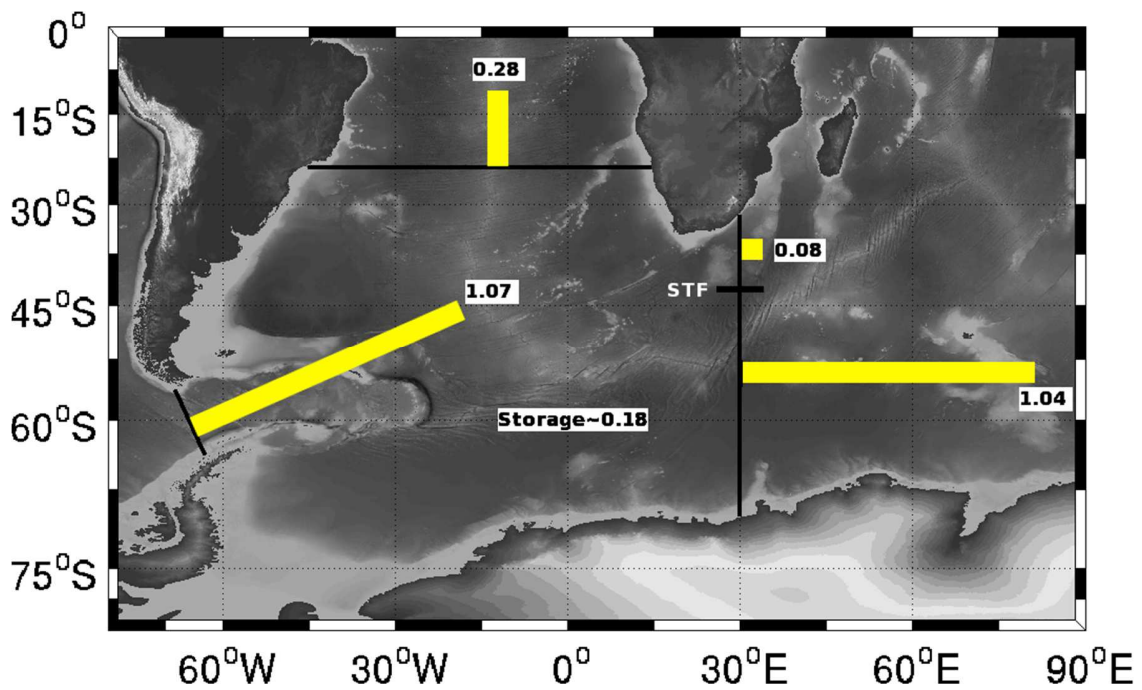


Figure 12: ΔC^* -derived distribution across 24°S of C^{ant} in 2009. Neutral density classes are labelled following the neutral density interfaces in Table 4. Units of $\mu\text{mol kg}^{-1}$.

1032



1033

1034 Figure 13: Schematic circulation for the each component of C^{ant} transport within the
 1035 inverse model solution (PgC yr^{-1}). The length of each bar is proportional to the net
 1036 transport. The implied net air-sea flux required to maintain the C^{ant} divergence is
 1037 $0.51 \pm 0.37 \text{ PgC yr}^{-1}$. Numbers at the end of each bar give transports in PgC yr^{-1} .
 1038 Uncertainties are presented in Table 7.

Table 1: GLODAP/CARINA correction factors as detailed in Gouretski and Jancke (2000), Wanninkhof et al. (2003), Key et al. (2004) and Hoppema et al. (2009). GLODAPv2 correction factors are detailed in Lauvset et al. (2016) and Olsen et al. (2016). Adjustments applied to hydrographic cruises along A13, A14, A15, A16 and A17 are required for section 2.1. Nitrate, phosphate, silicate and alkalinity are in units of $\mu\text{mol kg}^{-1}$. Salinity is listed as an addition in parts per million. Oxygen is listed in units of ml/l requiring multiplication by a factor of 43.55 to convert to $\mu\text{mol kg}^{-1}$ for all cruises apart from A21 (Drake Passage 2009) and 24°S 2008 where the multiplicative factors have already been optimised for $\mu\text{mol kg}^{-1}$.

	Salinity	Nitrate	Phosphate	Oxygen	Silicate	Alkalinity
A21 (Drake Passage 1990)	+1.1	+0.04	-0.06	+0.03	+4.9	×1.0
A21 (Drake Passage 2009)	×1.0	×0.975	×1.0	×1.035	×1.0	-6.0
I6S (30°E 1996)	×1.0	×0.96	×0.97	×1.0	×0.9	×1.0
24°S 2008	×1.0	×0.99	×1.0	×1.035	×0.95	×1.0
A13 (8°E 1995)	+2.8	-1.3	-0.153	+0.003	-3.0	×1.0
A14 (9°W 1995)	+2.3	-0.19	-0.033	+0.016	-1.9	×1.0
A15 (19°W 1994)	+0.3	-0.3	-0.023	-0.001	-1.5	×1.0
A16 (25°W 1989)	-0.5	-0.28	-0.029	+0.019	+0.3	×1.0
A17 (33°W 1994)	+1.8	+0.06	-0.024	+0.001	+1.6	×1.0

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1040 Table 2: Reference levels for each of the box boundaries. The 30°E section has been split
 1041 by the Subtropical Front (42.9°S) into an Agulhas and ACC regime.

Section	Reference Level	Reference
Drake Passage	Bottom	Meredith et al., (2011), (King & Jullion, in prep.).
24°S	1300 dbar	Bryden et al., (2011), Warren and Speer, (1991)
30°E Agulhas regime (North of 40°S)	2000 dbar	Arhan et al., (2003; Bryden et al., (2005)
30°E Agulhas regime (40°S – 42.9°S)	Bottom	Arhan et al., (2003)
30°E ACC regime (South of 42.9°S)	Bottom	Park et al., (2001)

1042 Table 3: Constraints applied to better construct the initial field for each of the sections
 1043 along the box boundary. Positive (negative) values indicate a transport into (out of) the
 1044 box. The boundary salinity transport refers to the net inflow of salinity transport across
 1045 Drake Passage, 24°S and the 30°E Agulhas regime combined to equal the net outflow of
 1046 salinity transport through the ACC regime at 30°E. All constraints are applied to better
 1047 constrain the initial field. The Ekman transport and the ACC regime boundary salinity
 1048 transport are not included as explicit constraints within the box inversion. Stated errors
 1049 are the residual noise terms from the conservation equations.

	Reference	Property	Vertical extent	Constraint	Solution
Drake Passage	Cunningham et al. (2003), Meredith et al. (2011)	Volume	Full depth	136.7±10 Sv	128.4±8.3 Sv
24°S:					
Full section	Coachman and Aagaard, (1988)	Salinity	Full depth	26±0.2 Sv psu	25.8±0.2 Sv psu

Vema and Hunter Channel	Hogg et al. (1999), McDonagh et al. (2002), Zenk et al. (1999)	Volume	$\theta < 2 \text{ }^\circ\text{C}$	$-6.9 \pm 2 \text{ Sv}$	$-6.7 \pm 1.9 \text{ Sv}$
Brazil Current	Bryden et al. (2011)	Volume	Above 300 dbar	$4.9 \pm 5 \text{ Sv}$	$5.8 \pm 0.1 \text{ Sv}$
Cape Basin (East of 6°E)	Arhan et al. (2003), McDonagh and King (2005)	Volume	$\theta < 2 \text{ }^\circ\text{C}$	$0 \pm 1 \text{ Sv}$	$0.2 \pm 0.1 \text{ Sv}$
Ekman transport 30°E:	Bryden et al. (2011)	Volume	Above 80 dbar	3.3 Sv	
Agulhas regime	McDonagh et al. (1999)	Volume	$\theta > 3.5 \text{ }^\circ\text{C}$	$9 \pm 3 \text{ Sv}$	$10.7 \pm 1.3 \text{ Sv}$
ACC regime	This study	Salinity	Full depth	Salinity transport inflow to box (4773.64 Sv psu)	

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1051 Table 4: Neutral density limits for each layer and corresponding water classes.
 1052 Definitions following Orsi et al. (1999), Naveira Garabato et al. (2002a), Heywood and
 1053 King (2002) and Naveira Garabato et al. (2009). A North Atlantic Deep Water (NADW)
 1054 neutral density class is labelled at $27.90 < \gamma^n < 28.10$ primarily for usage along the 24 °S
 1055 section where NADW is prevalent.

Layer	Lower limits	Water classes
1	$\gamma^n < 24$	
2	$24 < \gamma^n < 25$	Surface Water
3	$25 < \gamma^n < 26$	
4	$26 < \gamma^n < 26.80$	
5	$26.80 < \gamma^n < 26.90$	
6	$26.90 < \gamma^n < 27.00$	Subantarctic Mode Water (SAMW)
7	$27.00 < \gamma^n < 27.10$	
8	$27.10 < \gamma^n < 27.23$	
9	$27.23 < \gamma^n < 27.30$	
10	$27.30 < \gamma^n < 27.40$	Antarctic Intermediate Water (AAIW)
11	$27.40 < \gamma^n < 27.50$	
12	$27.50 < \gamma^n < 27.60$	
13	$27.60 < \gamma^n < 27.70$	Upper Circumpolar
14	$27.70 < \gamma^n < 27.80$	Deep Water
15	$27.80 < \gamma^n < 27.90$	(UCDW)
16	$27.90 < \gamma^n < 28.00$	North Atlantic Deep Water
17	$28.00 < \gamma^n < 28.10$	Lower Circumpolar (NADW)
18	$28.10 < \gamma^n < 28.20$	Deep Water
19	$28.20 < \gamma^n < 28.27$	(LCDW)
20	$28.27 < \gamma^n < 28.35$	
21	$28.35 < \gamma^n$	Antarctic Bottom Water (AABW)

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1058 Table 5: Meridional property transport from inverse studies and empirical analysis across
 1059 24°S, 30°S (WOCE A10) and nominally at 45°S (WOCE A11), adapted from McDonagh
 1060 and King (2005) and Williams (2007). The MOC strength in this study, is interpreted as
 1061 the southward flow of deep water, primarily NADW. The MOC strength estimate in
 1062 Dong et al. (2009) is an average of 17 hydrographic occupations. A northward net flux is
 1063 positive.

Source	Section	Freshwater (Sv)	Heat (PW)	Salt (Gg s⁻¹ or Sv psu)	MOC strength (Sv)
Ganachaud (1999)	A11	-	0.66±0.12	-	18±4
Holfort and Siedler (2001)	A11	-0.55±0.02	0.37±0.02	-26.37±0.73	21.7
McDonagh and King (2005)	A11	-0.7	0.43±0.08	-26	21.0±2
Naveira Garabato et al. (2014)	A11	-0.7±0.48	0.14±0.06	-29.2±17.2	15.8
Dong et al. (2009)	35°S	-	0.55±0.14	-	17.9
Rintoul (1991)	32°S	-	0.25	-	-
Lumpkin and Speer (2007)	32°S	-	0.60±0.08	-	-
Ganachaud (1999)	30°S	-	0.35±0.15	-	23±3
Holfort and Siedler (2001)	30°S	-0.51±0.02	0.29±0.05	-26.75±0.77	22.7
Ganachaud and Wunsch (2003)	30°S	-0.5±0.1	-	-26.7	-
McDonagh and King (2005)	30°S	-0.5±0.1	0.22±0.08	-26	19.9±2
Naveira Garabato et al. (2014)	30°S	-0.58±0.48	0.31±0.04	-13.8±17.1	13.7
Bryden et al. (2011)	24°S	-0.34/-0.29	0.7	-26	21.5 / 16.5
This study	24°S	-0.7±0.3	0.40±0.08	-25.8±0.2	20.2±2

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1066 Table 6: a) Net heat flux across 24°S separated into overturning and gyre components. b)
 1067 Overturning component of the salinity transport and associated M_{ov} and M_{az} transports.
 1068 Positive (negative) transport is defined as northwards (southwards) for compatibility with
 1069 Bryden et al. (2011).

a)	Overturning (PW)	Gyre (PW)	Total (PW)
This study	0.52	-0.12	0.40
Bryden et al. (2011) 2009 section	0.76	-0.07	0.68
Bryden et al. (2011) 1983 section	0.53	-0.14	0.38

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b)	Overturning (Sv psu)	M_{ov} (Sv)	M_{az} (Sv)
This study	3.3	-0.09	0.16
Bryden et al. (2011) 2009 section	4.6	-0.13	0.12
Bryden et al. (2011) 1983 section	3.3	-0.09	0.21

Table 7: C^{ant} transports at the box boundary, C^{ant} storage within the box and C^{ant} air-sea flux in Pg C yr^{-1} . Positive (negative) values indicate a transport into (out of) the box.

Section	C^{ant} Transport (Pg C yr^{-1})
Drake Passage	$+1.07 \pm 0.44$
24°S	-0.28 ± 0.16
30°E: Agulhas	-0.08 ± 0.07
30°E: ACC	-1.04 ± 0.42
Total	-0.33 ± 0.31
Storage	$+0.18 \pm 0.12$
Air-Sea flux	$+0.51 \pm 0.37$

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Table 8: Transport-weighted C^{ant} ($\mu\text{mol kg}^{-1}$) at each box boundary. For 24°S and 30°E (Agulhas), the transports are separated into north-south, or east-west components respectively, given the substantial flow in both directions. Uncertainties are the standard error of the mean with units of $\mu\text{mol kg}^{-1}$.

Neutral density	Drake Passage	24°S		30°E (Agulhas)		30°E
		North	South	West	East	
Surface	34.9±0.4	65.2±0.3	63.2±0.4	48.2±0.3	49.7±0.2	50.0±0.5
SAMW	39.2±0.3	51.9±0.4	50.8±0.4	29.2±0.6	35.6±0.7	40.9±0.6
AAIW	36.0±0.5	26.9±0.3	25.8±0.3	16.8±0.5	17.8±0.4	31.8±0.5
UCDW	16.3±0.2	14.3±0.1	13.9±0.1	16.5±0.1	16.3±0.1	16.6±0.2
LCDW	6.9±0.1	10.2±0.1	10.8±0.1	11.1±0.1	12.6±0.1	10.3±0.1
AABW	2.2±0.1	12.6±0.2	11.1±0.2	10.1±0.1	11.7±0.1	11.1±0.1
Total	20.9±0.2	30.1±0.2	25.6±0.2	33.8±0.3	35.1±0.3	18.8±0.1

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1077 Table 9: Mean Penetration Depth (MPD), mean ΔC_{ml}^{ant} ($\mu\text{mol kg}^{-1} \text{yr}^{-1}$) within mixed layer
 1078 and mean in-situ density ρ_{ml} (kg m^{-3}) within mixed layer for Drake Passage, 30°E, and
 1079 24°S and a mean of the hydrographic sections. Along 30°E, C^{ant} is normalised to a mean
 1080 temperature. MPD is estimated to have a 20% uncertainty, and a $\pm 0.5 \mu\text{mol kg}^{-1} \text{yr}^{-1}$
 1081 ΔC_{ml}^{ant} uncertainty.

	MPD (m)	ΔC_{ml}^{ant} ($\mu\text{mol kg}^{-1} \text{yr}^{-1}$)	ρ_{ml} (kg m^{-3})
Drake Passage	259.8	0.84	1027.0
24°S	933.2	0.85	1024.8
30°E	624.3	0.45	1026.2
Mean	605.8	0.71	1025.9

1082 Table 10: Comparison of C^{ant} storage rate ($\text{mol m}^{-2} \text{yr}^{-1}$) for the South Atlantic (south of
 1083 15°S), and South Atlantic sector of the Southern Ocean. For Peng and Wanninkhof
 1084 (2010), the two estimates derive from two different calculation methods.

Author	Region	Storage rate ($\text{mol m}^{-2} \text{yr}^{-1}$)
Holfort et al. (1998)	10°S -30°S	0.59±0.12
Murata et al. (2008)	Along 30°S	0.6±0.1
Peng and Wanninkhof (2010)	South of 15°S	0.56/0.35±0.3
Wanninkhof et al. (2010)	South of 15°S	0.76
Ríos et al. (2012)	10°N-55°S, western basin	0.92±0.13
This study	Drake Passage	0.22±0.29
This study	24°S	0.81±0.53
This study	30°E	0.29±0.18
This study	Mean	0.44±0.30

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