



## RESEARCH ARTICLE

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## Special Section:

Advances in scaling and  
modeling of land-atmosphere  
interactions

## Key Points:

- We describe a new tool for investigating canopy-chemistry-turbulence interactions at isoprene-dominated deciduous forests
- Spatial segregation of air masses in the canopy can bias reaction rates assuming well-mixed conditions by up to -48% to +23%
- There is stronger segregation for many reactions when soil NO emissions reflect the upper bound of observed values for temperate forests

## Supporting Information:

Supporting Information may be found in  
the online version of this article.

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## Large Eddy Simulation for Investigating Coupled Forest Canopy and Turbulence Influences on Atmospheric Chemistry

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**Abstract** Exchanges of reactive gases between forests and the atmosphere influence tropospheric chemistry, climate, and ecosystems. Chemistry inside canopies can alter these exchanges yet understanding of the chemistry and processes influencing the chemistry is incomplete. Previous work prioritizes complex chemical mechanisms over micrometeorology in the multilayer models of forests used to understand in-canopy chemistry and its impacts. Instead, here we use a new simplified chemical mechanism and resolve turbulence in a multilayer canopy model. Specifically, we describe a new version of the National Center for Atmospheric Research large eddy simulation (LES) coupled to both a multilayer canopy model and a chemical mechanism. The mechanism has 41 reactions and 19 gases for ozone, NO<sub>x</sub> (=NO + NO<sub>2</sub>), HO<sub>x</sub> (=OH + HO<sub>2</sub>), and isoprene chemistry. We evaluate the mechanism against two more complex mechanisms using box modeling. We configure the LES for a temperate deciduous forest and summer midday weak-wind buoyancy-forced conditions. The multilayer canopy model necessitates estimates of reactant sources and sinks at each canopy level, and thus we describe multilayer parameterizations for dry deposition and biogenic emissions. For the first time with an LES coupled to both a multilayer canopy and chemistry, we demonstrate variability in chemical reactions due to turbulence-induced segregation of reactants inside and just above the canopy. For the cases examined here, in-canopy segregation alters reaction rates that only consider well-mixed conditions by up to -48% to +23%. For many reactions, segregation is stronger when soil NO emissions reflect the upper bound of observed values for temperate forests.

**Plain Language Summary** The atmospheric boundary layer (ABL) (the part of the atmosphere closest to the earth) organizes itself into structured air motions that vertically transport momentum, energy, and reactive gases toward and away from the Earth's surface. This is important because the Earth's surface provides sources and sinks of momentum, energy, and reactive gases. The presence of forest canopies can alter the structured air motions in the ABL, and the structured air motions influence canopy sources and sinks of momentum, energy, and reactive gases, implying a coupling between the canopy and the structured air motions that may influence atmospheric chemistry. We describe a computer model with very high spatial resolution that can be used to examine the impact of the relationship between these organized motions, the canopy, and chemistry. Other computer models that have been used to examine atmospheric chemistry in canopies only consider the average effect of the organized motions. We show that the organized motions can influence chemical reaction rates substantially in the canopy, rendering the model described here an important tool for interpreting observations and advancing understanding of the impact on atmospheric chemistry important for climate and air pollution.

## 1. Introduction

Land-atmosphere exchanges of reactive gases play an important role in our earth system, influencing tropospheric chemistry, climate, and ecosystems (Arneth et al., 2011; Barth et al., 2005; Farmer & Riches, 2020; Fowler et al., 2009; He et al., 2021; Steiner, 2020). For example, emissions of biogenic volatile organic compounds (BVOCs) from vegetation contribute to the formation of tropospheric ozone and secondary organic aerosol, both air pollutants and short-lived climate forcers (Fehsenfeld et al., 1992; Fiore et al., 2012; Hallquist et al., 2009; Henze & Seinfeld, 2006; Trainer et al., 1987). Air pollution and the atmospheric oxidizing capacity also depend on NO<sub>x</sub> (=NO + NO<sub>2</sub>) (Crutzen, 1974, 1979; Levy, 1971), which has a soil source and a vegetation sink (Delaria

et al., 2018; Jacob & Bakwin, 1991; Wang et al., 2020; Yienger & Levy, 1995). A key but overlooked component of forest-atmosphere chemical exchanges is turbulence.

Generally, organized eddy structures comprise turbulence in the atmospheric boundary layer (ABL), with their size and character changing with driving forces (e.g., Moeng & Sullivan, 1994; Weckwerth et al., 1997). Inside and just above forest canopies, turbulence is unique. Near canopy top, the mean wind profile has a hydrodynamically unstable inflection point due to vegetation absorbing momentum through its depth. ABL-scale turbulence brings high-momentum air to the canopy top, producing shear-driven canopy-scale eddies (i.e., sweeps and ejections) distinct from the background flow (Finnigan et al., 2009; Gao et al., 1989, 1992; Raupach et al., 1996). In unstable conditions, buoyant plumes emanating from the ground and canopy also influence exchange between canopy layers and aloft (e.g., Brunet, 2020; Dupont & Patton, 2012; Patton et al., 2016).

Turbulence influences forest-atmosphere chemical exchanges in several ways. For example, turbulence transports reactive gases out of canopies, allowing compounds emitted by canopy elements or produced within canopy air spaces to influence tropospheric chemistry at local-to-global scales. Turbulence also transports gases into the canopy, allowing gases in the ABL to meet their fate via in-canopy ambient chemical loss and dry deposition to canopy surfaces. Organized structures in turbulence can lead to spatial variations in reactant abundances and thus the physical separation of reactants (“segregation”). Because organization in turbulence can happen on timescales important for atmospheric chemistry, segregation can cause reaction rates to speed up or slow down relative to rates that assume well-mixed conditions (Kim et al., 2016; Krol et al., 2000; Patton et al., 2001; Schumann, 1989; Sykes et al., 1992, 1994). Segregation may be particularly important in the canopy where there are distinct sources and sinks of reactants that vary in space and time (Edburg et al., 2012; Patton et al., 2001).

Despite the key role of turbulence in forest-atmosphere chemical exchanges, models almost always only consider the influence of turbulence in a bulk sense. Thus, turbulent transport is often not accurately represented, and the impact of segregation cannot be quantified. While multilayer models of vegetation canopies (hereinafter, MLMs) have been coupled to complex chemical mechanisms to examine exchanges of reactive gases and the related in-canopy chemistry for many years (e.g., Ashworth et al., 2015; Bryan et al., 2012; Forkel et al., 2006; Ganzeveld et al., 2002; Wei et al., 2021; Wolfe & Thornton, 2011), chemistry MLMs have yet to be coupled to turbulence-resolving large eddy simulation (LES) with one exception (Fuentes et al., 2022). LES is a powerful tool for investigating turbulence including within and above canopies (e.g., Kanda & Hino, 1994; Ma & Liu, 2019; Shaw & Schumann, 1992; Su et al., 1998) yet can be computationally expensive.

The National Center for Atmospheric Research LES is coupled to an MLM (hereinafter, NCAR-LES-MLM) such that vertical changes in light attenuation, canopy and soil sources and sinks of moisture and heat, the canopy-imposed sink of momentum, and canopy drag all influence turbulent flow (Patton et al., 2016). We document a new version of the NCAR-LES-MLM coupled to a simplified chemical mechanism (hereinafter, NCAR-LES-MLM-CHEM) that represents ozone,  $\text{NO}_x$ ,  $\text{HO}_x$  ( $=\text{OH} + \text{HO}_2$ ), and isoprene chemistry. Our model prioritizes resolution of turbulence over complex chemistry, in contrast to the current suite of tools used to investigate chemical exchanges that prioritize complex chemistry. To our knowledge, reactant segregation in and above the canopy has never been demonstrated using an LES coupled to a chemistry MLM through canopy drag and interactive sources and sinks of heat, water vapor, and reactants. Specifically, unique to the NCAR-LES-MLM-CHEM is dynamic and spatially varying leaf-level emission and dry deposition of reactants, as well as water vapor and heat sources. In Fuentes et al. (2022), which also does not investigate segregation, sources and sinks of heat and water vapor are forced with observed fluxes, and emissions of isoprene are imposed with an offline vertical distribution.

We configure the NCAR-LES-MLM-CHEM for midday summer weak-wind buoyancy-forced conditions at a temperate deciduous forest. Deciduous forests are key sources of isoprene, globally the largest emitted BVOC and non-methane volatile organic compound (Guenther et al., 1995, 2012). Our simulated forest, with its relatively high ambient humidity and soil moisture, dense isoprene-emitting deciduous trees, idealizes a southeast US forest. The southeast US has been a focus in atmospheric chemistry research due to field campaigns in summer 2013 that aimed to understand climate and pollution relevant low- $\text{NO}_x$  isoprene chemistry (e.g., Carlton et al., 2018; Mao et al., 2018). Previous work used airborne observations over the southeast US during summer 2013 to constrain segregation between OH and isoprene in the ABL (Kaser et al., 2015). Here we leverage the NCAR-LES-MLM-CHEM to probe how segregation impacts first-generation isoprene chemistry. We focus on inside and just above the canopy where segregation may be particularly important due to unique turbulence and

**Table 1**  
Simulated Conditions by the NCAR-LES-MLM-CHEM

Case	$z_i$ (m)	$u_*$ (m s <sup>-1</sup> )	$\overline{w'\theta'_v}$ (m K s <sup>-1</sup> )	$w_*$ (m s <sup>-1</sup> )	$w_m$ (m s <sup>-1</sup> )
Low soil NO	402	0.22	0.14	1.22	1.23
High soil NO	402	0.22	0.14	1.22	1.23

Note. All variables are horizontally and temporally averaged. All variables except the ABL height ( $z_i$ ) are from canopy top.  $u_*$  is friction velocity.  $\overline{w'\theta'_v}$  is the buoyancy flux.  $w_*$  is the Deardorff convective velocity scale.  $w_m$  is a scaling velocity combining shear and buoyancy forcing.

interactive reactant sources and sinks. We also investigate segregation in the context of uncertainty in soil NO emissions. For example, southeast U.S. soil NO emissions are relatively unconstrained but expected to be an important contribution to regional NO<sub>x</sub> and air pollution (Romer et al., 2018; Travis et al., 2016), and the few observations of soil NO emissions for temperate deciduous forests vary by three orders of magnitude (Thornton et al., 1997; Valente & Thornton, 1993; Williams & Fehsenfeld, 1991).

## 2. Model Description and Evaluation

### 2.1. NCAR Multilayer Canopy Large Eddy Simulation

The NCAR-LES (e.g., Moeng, 1984; Moeng & Wyngaard, 1988; Sullivan et al., 1996; Sullivan & Patton, 2011) solves the three-dimensional time-dependent spatially filtered fluid motion equations under the Boussinesq approximation. A discrete Poisson equation for pressure enforces incompressibility. In Section 2.1, we describe the NCAR-LES-MLM, and in Section 2.2 we introduce the NCAR-LES-MLM-CHEM. Note that the NCAR-LES has been coupled to chemistry before but without a vertically resolved and coupled forest canopy (hereinafter, NCAR-LES-CHEM).

In the NCAR-LES-MLM (Patton et al., 2016), terms representing pressure and viscous drag as well as scalar sources and sinks are produced by spatially filtering the conservation equations in the presence of canopy elements (Patton & Finnigan, 2013). Boundary conditions in the horizontal directions are periodic. A rough-wall boundary condition is imposed at the surface with Monin Obukhov Similarity Theory (MOST) determining transfer between the lowest grid point and the ground. Sub-filter scale (SFS) momentum and scalar fluxes are parameterized with down-gradient diffusion assumptions and estimated SFS turbulent kinetic energy (TKE) (Deardorff, 1980; Shaw & Patton, 2003). We assume that all scalars mix similarly to heat at SFSs.

The physical model configuration follows Clifton and Patton (2021). In brief, the simulation domain is 2,048 m in both horizontal directions and 1,024 m in the vertical direction. The spatial resolution is 2 m × 2 m × 2 m. The horizontal dimensions are at least fivefold larger than the ABL height ( $z_i$ ; Table 1) to minimize the influence of periodic boundary conditions on ABL-scale turbulent structures.

The NCAR-LES-MLM simulates flow interacting with a forest with horizontally homogeneous properties. The specified vertical leaf area density (LAD) profile has a maximum in the upper canopy and a relatively sparse understory, typical for a deciduous tree (Figure 1). The canopy height ( $h$ ) is 20 m. The NCAR-LES-MLM represents multiple canopy layers in each 2 m × 2 m vertical column in the 2,048 m × 2,048 m horizontal domain. In each grid cell inside the canopy, the MLM simulates processes like leaf energy balance and photosynthesis

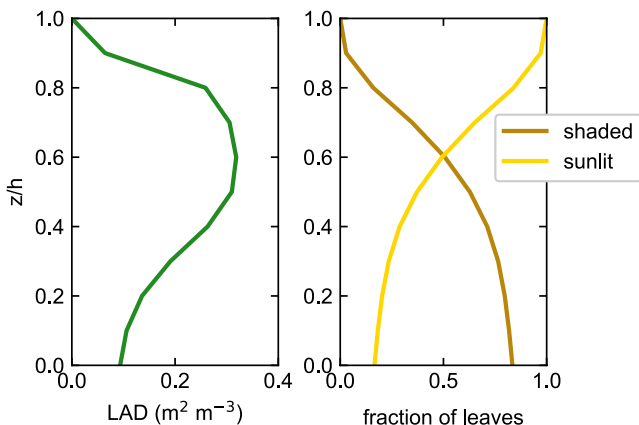
for sunlit and shaded leaves in response to the LES forcing and provides leaf-level water and heat sources back to the LES at each time step. Canopy drag occurs primarily through pressure forces and varies spatially with LAD and instantaneous local wind speed (Patton et al., 2016).

We prescribe the vertical distribution of sunlit versus shaded leaves (Figure 1) using Beer's Law for which incoming direct radiation differs based on LAD and solar zenith angle ( $\theta$ ) (e.g., Monteith & Unsworth, 1990; Moon et al., 2020). At a given height ( $z$ ), the sunlit fraction ( $f_{\text{sun}}$ ) is:

$$f_{\text{sun}} = e^{-\Omega \frac{G}{\cos \theta} \int_0^z \text{LAD}(z) dz} \quad (1)$$

The variable  $\theta$  is the solar zenith angle in radians (0.317), the parameter  $G$  relates to leaf orientation (0.5), and the parameter  $\Omega$  is the clumping index (0.85).

Each vertical column has a 1-m deep soil layer resolved at four levels, and there is a set of roots accessing the soil at every horizontal grid point. Initial soil moisture and temperature conditions are from a 2-year High-Resolution



**Figure 1.** Vertical profiles of leaf area density (LAD) and the fraction of sunlit versus shaded leaves (and canopy air space) from the ground to canopy height ( $h$ ).

Land Data Assimilation System simulation (Chen et al., 2007; Patton et al., 2016) to ensure that the initial soil conditions are in equilibrium with the atmosphere.

We prescribe initial vertical profiles of specific humidity, virtual potential temperature, and streamwise wind. The streamwise geostrophic wind is  $1 \text{ m s}^{-1}$  and the initial vertical profile of the streamwise wind is set to the streamwise geostrophic wind. Initial virtual potential temperature has a  $0.08 \text{ K m}^{-1}$  gradient spanning 100 m centered on the initial ABL height (400 m) and a  $0.003 \text{ K m}^{-1}$  gradient in the free troposphere (FT; Sullivan & Patton, 2011). Incoming solar radiation is  $847 \text{ W m}^{-2}$ . The initial specific humidity is  $18 \text{ g kg}^{-1}$  in the ABL and  $1 \text{ g kg}^{-1}$  in the FT.

To initiate turbulence, divergence-free fluctuations are added to the initial horizontal velocity field over the five vertical grid points spanning canopy top (Patton et al., 2016). We discard model results before  $w_m$ , a velocity scale incorporating shear and buoyancy forces (Moeng & Sullivan, 1994), becomes constant at canopy top. We then average results over five turnover times ( $=z_i/w_m$ ; Table 1) corresponding to the largest and slowest eddies in the domain. Our simulations, excluding spin up, are each 27 min. The time step varies but averages 0.32 s.

Unless stated otherwise, all quantities reported can be assumed to be averaged horizontally and temporally. We only use overbars in fluxes or covariances, or when the emphasis is needed to understand a quantity; in these cases, the overbars represent horizontal averages, and the primes represent deviations from the horizontal average. Time averaging happens after horizontal averaging. Because all variables are averaged in time unless stated otherwise, we do not include notation denoting this averaging.

Although our model domain extends up into the FT, the focus of this paper is the vicinity of the canopy, strictly speaking from the ground to  $3h$ .

## 2.2. Chemistry

### 2.2.1. New Simplified Chemical Mechanism

Due to the computational expense of the NCAR-LES-MLM-CHEM, we employ a simplified gas phase chemical mechanism that simulates the features of tropospheric chemistry involving ozone, isoprene,  $\text{NO}_x$  and  $\text{HO}_x$  over deciduous forests. The mechanism has 19 gases (Table 2) and 41 reactions (Table 3) and is designed to represent first-generation isoprene chemistry and similar to Wang et al. (2021). The mechanism is based on Su et al. (2016). Reaction kinetics are updated to reflect the Jet Propulsion Laboratory (JPL; Burkholder et al., 2015) or the International Union of Pure and Applied Chemistry (IUPAC; Atkinson et al., 2004, 2006) databases. The isoprene oxidation is updated based on MOZART T1 (Emmons et al., 2020). We include a simplified representation for the isomerization of the organic peroxy radical ( $\text{RO}_2$ ) from OH reaction with isoprene to form an isoprene-derived hydroperoxy aldehyde (HPALD), which is lumped with isoprene hydroxy aldehydes (HYDRALD). Isomerization of  $\text{RO}_2$  is important for  $\text{HO}_x$  recycling (Crounse et al., 2011). Due to mechanism design, later-generation isoprene chemistry may not be as well represented.

The NCAR-LES-MLM-CHEM uses the Kinetic PreProcessor (KPP; Damian et al., 2002; Sandu & Sander, 2006) to connect the specified chemical mechanism to a Rosenbrock solver for solving the stiff set of ordinary differential equations. Though the version of the NCAR-LES-CHEM with KPP was reported in Lenschow et al. (2016) with a three-reaction mechanism, we emphasize that using KPP and the Rosenbrock solver is an update from other NCAR-LES-CHEM studies (e.g., Kim et al., 2012; Li, Barth, et al., 2016; Li, Su, et al., 2016; Li et al., 2017; Vilà-Guerau de Arellano et al., 2005), allowing the user to easily change chemical mechanisms and employ a highly accurate conservative solver.

We initialize the NCAR-LES-MLM-CHEM with ABL and FT mixing ratios (Table 2) from an NCAR-LES-CHEM simulation for the Southeast Atmosphere Study (SAS) campaign without the resolved canopy at 12:00 Central Standard Time (CST) that started at 6:00 CST (the chemistry was reinitialized to observed values at 08:00 CST). Our NCAR-LES-MLM-CHEM spin up is 1,108 s, reflecting the time it takes the turbulence to develop and reach equilibrium with the forcing. During this transition, ABL-averaged reactant mixing ratios reveal unphysical fluctuations (e.g., see the first hundreds of s in Figure S1 in Supporting Information S1). However, there is likely no effect of the fluctuations on the spatial distribution of the reactants during the production part of the simulation. For example, because time averaging begins 358 s ( $\sim 6$  min) after the final fluctuation in the ABL-averaged reactant mixing ratios, and the turnover time for ABL-size eddies is  $\sim 325$  s during production, any effect of

**Table 2**

*The Gaseous Species Used in the Chemical Mechanism, the Properties of Each Species Needed for Dry Deposition Estimates, and Their Initial Mixing Ratios in the Atmospheric Boundary Layer (ABL) and Free Troposphere (FT)*

Chemical species	Molecular weights (g mol <sup>-1</sup> )	Reactivity $f_0$ (unitless)	Solubility $H^{\text{eff}}$ (M atm <sup>-1</sup> )	Initial mixing ratios (ppb)
OH	17.01	1.0	$3.90 \times 10^1$	$0.584 \times 10^{-4}$ (ABL) $0.193 \times 10^{-3}$ (FT)
HO <sub>2</sub>	33.01	0.1	$6.90 \times 10^2$	0.026 (ABL) 0.021 (FT)
O <sub>3</sub>	48.00	1.0	$1.03 \times 10^{-2}$	33.66 (ABL) 37.84 (FT)
H <sub>2</sub> O <sub>2</sub>	34.01	1.0	$8.70 \times 10^4$	2.49 (ABL) 1.23 (FT)
CO	28.01	0.0	$9.81 \times 10^{-4}$	105.53 (ABL) 90.27 (FT)
NO	30.01	0.0	$1.92 \times 10^{-3}$	0.031 (ABL) 0.012 (FT)
NO <sub>2</sub>	46.01	0.1	$1.20 \times 10^{-2}$	0.067 (ABL) 0.018 (FT)
NO <sub>3</sub>	62.00	1.0	$3.80 \times 10^{-2}$	$0.491 \times 10^{-5}$ (ABL) $0.129 \times 10^{-5}$ (FT)
N <sub>2</sub> O <sub>5</sub>	108.01	0.1	2.14	$0.272 \times 10^{-3}$ (ABL) $0.146 \times 10^{-4}$ (FT)
CH <sub>4</sub>	16.04	0.0	$1.41 \times 10^{-3}$	1,831.36 (ABL) 1,803.33 (FT)
CH <sub>3</sub> O <sub>2</sub>	47.03	1.0	2.38	0.009 (ABL) 0.014 (FT)
HCHO	30.03	1.0	$3.23 \times 10^3$	2.40 (ABL) 0.491 (FT)
ISOP	68.11	0.0	$3.45 \times 10^{-2}$	2.82 (ABL) $0.109 \times 10^{-3}$ (FT)
ISOPAO <sub>2</sub>	117.12	1.0	2.38	0.020 (ABL) $0.436 \times 10^{-5}$ (FT)
ISOPBO <sub>2</sub>	117.12	1.0	2.38	0.012 (ABL) $0.277 \times 10^{-5}$ (FT)
ISOPOOH	118.13	1.0	$3.50 \times 10^6$	0.952 (ABL) 0.007 (FT)
[HYDRALD + HPALD]	108.11	1.0	$1.70 \times 10^5$	0.550 (ABL) 0.002 (FT)
XO <sub>2</sub>	149.12	1.0	$1.00 \times 10^{11}$	0.003 (ABL) $0.957 \times 10^{-4}$ (FT)
[MVK + MACR]	70.09	1.0	23.75	1.550 (ABL) 0.194 (FT)

*Note.* Values in columns (2)–(3) are from Schwantes et al. (2020) or CESM2.2 except values for ISOPAO<sub>2</sub> and ISOPBO<sub>2</sub>, for which we use CH<sub>3</sub>O<sub>2</sub> values, and XO<sub>2</sub> for which we assume very high solubility and reactivity. Brackets around species added together denote species lumped as one in our model. XO<sub>2</sub> is an RO<sub>2</sub> formed oxidation of [HYDRALD + HPALD] and ISOPOOH (Table 3; Emmons et al., 2010).

**Table 3**  
The Reactants, Products, and Reaction Coefficients of the Reactions in the Chemical Mechanism

Reactants	Products	Rate coefficient (Molecules cm <sup>-3</sup> s <sup>-1</sup> )
O <sub>3</sub>	2 OH	$K_{\text{H}_2\text{O}} = 1.63 \times 10^{-10} [\text{H}_2\text{O}] e^{60/T}$ $K_{\text{N}_2} = 2.15 \times 10^{-11} 0.79 [\text{M}] e^{110/T}$ $K_{\text{O}_2} = 3.30 \times 10^{-11} 0.21 [\text{M}] e^{55/T}$ $j_{\text{O}(1\text{D})} = 4.551 \times 10^{-5}$ $K_{\text{H}_2\text{O}} \frac{j_{\text{O}(1\text{D})}}{k_{\text{H}_2\text{O}} + k_{\text{N}_2} + k_{\text{O}_2}}$ $j_{\text{NO}_2} = 1.144 \times 10^{-2}$ $j_{\text{HCHO}_r} = 3.669 \times 10^{-5}$ $j_{\text{HCHO}_m} = 5.352 \times 10^{-5}$ $4.8 \times 10^{-11} e^{250/T}$ $1.70 \times 10^{-12} e^{-940/T}$ $(2.3 \times 10^{-13} e^{600/T} + 1.7 \times 10^{-33} [\text{M}] e^{1,000/T}) (1.0 + 1.4 \times 10^{-21} [\text{H}_2\text{O}] e^{2,200/T})$ $2.03 \times 10^{-16} e^{\frac{693}{T}} \left( \frac{T}{300} \right)^{4.57}$ $1.4 \times 10^{-12} e^{-1,310/T}$ $\frac{1.80 \times 10^{-30} \left( \frac{300}{T} \right)^3 [\text{M}]}{1 + \frac{1.80 \times 10^{-30} \left( \frac{300}{T} \right)^3 [\text{M}]}{2.80 \times 10^{-11}}}$ $0.6^{1/1 + \log_{10}(1.80 \times 10^{-30} (300/T)^3 [\text{M}]/2.80 \times 10^{-11})^2}$
NO <sub>2</sub>	NO + O <sub>3</sub>	
HCHO	2 HO <sub>2</sub> + CO	
HCHO	CO	
OH + HO <sub>2</sub>	H <sub>2</sub> O	
OH + O <sub>3</sub>	HO <sub>2</sub>	
HO <sub>2</sub> + HO <sub>2</sub>	H <sub>2</sub> O <sub>2</sub>	
HO <sub>2</sub> + O <sub>3</sub>	OH	
NO + O <sub>3</sub>	NO <sub>2</sub>	
NO <sub>2</sub> + OH		
NO <sub>2</sub> + O <sub>3</sub>	NO <sub>3</sub>	
NO + HO <sub>2</sub>	NO <sub>2</sub> + OH	
NO <sub>2</sub> + NO <sub>3</sub>	N <sub>2</sub> O <sub>5</sub>	$\frac{2.40 \times 10^{-30} \left( \frac{300}{T} \right)^3 [\text{M}]}{1 + \frac{2.40 \times 10^{-30} \left( \frac{300}{T} \right)^3 [\text{M}]}{1.60 \times 10^{-12} \left( \frac{300}{T} \right)^{-0.7}}}$ $0.6^{1/1 + \log_{10}(2.40 \times 10^{-30} (300/T)^3 [\text{M}]/1.60 \times 10^{-12} (300/T)^{-0.7})^2}$
NO + NO <sub>3</sub>	NO <sub>2</sub> + NO <sub>2</sub>	$1.8 \times 10^{-11} e^{\frac{110}{T}}$ $4.0 \times 10^{-12}$
HO <sub>2</sub> + NO <sub>3</sub>	OH + NO <sub>2</sub>	
N <sub>2</sub> O <sub>5</sub> + M	NO <sub>2</sub> + NO <sub>3</sub>	$\frac{1.8 \times 10^{-30} \left( \frac{300}{T} \right)^3 [\text{M}]}{1 + \frac{1.8 \times 10^{-30} \left( \frac{300}{T} \right)^3 [\text{M}]}{2.8 \times 10^{-11}}}$ $0.6^{1/1 + \log_{10}(1.8 \times 10^{-30} (300/T)^3 [\text{M}]/2.8 \times 10^{-11})}$ $1.724 \times 10^{26} e^{-10840/T} / [\text{M}]$
N <sub>2</sub> O <sub>5</sub>		
NO <sub>3</sub>	0.11 NO + 0.89 NO <sub>2</sub> + 0.89 O <sub>3</sub>	$2.5 \times 10^{-22} [\text{H}_2\text{O}]$ $j_{\text{NO}_3, \text{O}} = 2.591 \times 10^{-2} + 2.051 \times 10^{-1}$
CH <sub>3</sub> O <sub>2</sub> + NO	HCHO + 0.99 NO <sub>2</sub> + HO <sub>2</sub>	$2.80 \times 10^{-12} e^{300/T}$ $4.10 \times 10^{-13} e^{250/T}$
CH <sub>3</sub> O <sub>2</sub> + HO <sub>2</sub>	HO <sub>2</sub>	$1.5 \times 10^{-13} (1 + 6 \times 10^{-7} k_b [\text{M}] T)$ $2.45 \times 10^{-12} e^{-1,775.0/T}$
CO + OH	CH <sub>3</sub> O <sub>2</sub>	
CH <sub>4</sub> + OH	HO <sub>2</sub> + CO	$5.50 \times 10^{-12} e^{125.0/T}$ $2.54 \times 10^{-11} e^{410.0/T}$
HCHO + OH		
ISOP + OH	0.6 ISOPAO <sub>2</sub> + 0.4 ISOPBO <sub>2</sub>	

**Table 3**  
*Continued*

Reactants	Products	Rate coefficient (molecules cm <sup>-3</sup> ) <sup>-(n-1)</sup> s <sup>-1</sup>
ISOPAO <sub>2</sub> + CH <sub>3</sub> O <sub>2</sub>	HO <sub>2</sub> + 1.5 HCHO + 0.75 [MVK + MACR]	$5.0 \times 10^{-13} e^{4000/T}$
ISOPAO <sub>2</sub> + HO <sub>2</sub>	ISOPOOH	$8.0 \times 10^{-13} e^{7000/T}$
ISOPAO <sub>2</sub> + NO	0.9 NO <sub>2</sub> + 0.9 [MVK + MACR] + 0.9 HCHO + 0.9 HO <sub>2</sub> + 0.9 CO	$4.4 \times 10^{-12} e^{1800/T}$
ISOPBO <sub>2</sub> + CH <sub>3</sub> O <sub>2</sub>	HO <sub>2</sub> + 0.75 HCHO + 0.75 [HYDRALD + HPALD]	$5.0 \times 10^{-13} e^{4000/T}$
ISOPBO <sub>2</sub> + HO <sub>2</sub>	ISOPOOH	$8.0 \times 10^{-13} e^{7000/T}$
ISOPBO <sub>2</sub>	[HYDRALD + HPALD] + HO <sub>2</sub>	$1.6 \times 10^{-9} e^{-8300/T}$
ISOPBO <sub>2</sub> + NO	0.87 [HYDRALD + HPALD] + 0.9 NO <sub>2</sub> + 0.97 HO <sub>2</sub> + 0.15 CO + 0.1 HCHO	$4.4 \times 10^{-12} e^{1800/T}$
ISOPOOH + OH	0.4 XO <sub>2</sub> + 0.6 OH	$1.52 \times 10^{-11} e^{2000/T}$
ISOPOOH	0.69 [MVK + MACR] + 0.69 HCHO + HO <sub>2</sub>	$j_{\text{ISOPOOH}} = 6.457 \times 10^{-6}$
[HYDRALD + HPALD] + OH	XO <sub>2</sub>	$1.86 \times 10^{-11} e^{1750/T}$
[HYDRALD + HPALD]	OH + HO <sub>2</sub>	0.004 $j_{\text{NO}_2}$
XO <sub>2</sub> + HO <sub>2</sub>		$8.0 \times 10^{-13} e^{7000/T}$
XO <sub>2</sub> + NO	0.9 NO <sub>2</sub> + HO <sub>2</sub> + 0.25 CO + 0.25 HCHO	$2.7 \times 10^{-12} e^{3600/T}$
[MVK + MACR] + OH	HO <sub>2</sub> + HCHO + 1.5 CH <sub>3</sub> O <sub>2</sub>	$2.4 \times 10^{-11}$
ISOP + O <sub>3</sub>	0.785 HCHO + 0.29 H <sub>2</sub> O <sub>2</sub> + 0.363 CO + 0.125 HO <sub>2</sub> + 0.205 OH + 0.675 [MVK + MACR] + 0.181 CH <sub>3</sub> O <sub>2</sub>	$1.03 \times 10^{-14} e^{-19950/T}$
ISOP + NO <sub>3</sub>		$3.15 \times 10^{-12} e^{-4500/T}$
H <sub>2</sub> O <sub>2</sub>	2 OH	$j_{\text{H}_2\text{O}_2} = 8.499 \times 10^{-6}$

*Note.* The parameter  $n$  is the number of reactants in the chemical reaction;  $T$  is air temperature in K;  $j_X$  is photolysis rate for species  $X$  (values given are not scaled by fraction of sunlight within the canopy);  $[M]$  is the third-body concentration (concentration of air used in molecules cm<sup>-3</sup>);  $k_b$  is the Boltzmann constant ( $1.38 \times 10^{-16}$  kJ K<sup>-1</sup>). Brackets around species added together denote species lumped as one in our model.

**Table 4**  
*Box Model Conditions*

Condition	Value
Solar zenith angle	18°
Temperature	301 K
Relative humidity	80%
Pressure	101,325 Pa
CH <sub>4</sub>	1,831 ppb
CO	106 ppb
NO	0.1 ppb (clean), 2.5 ppb (polluted)
HCHO	2.0 ppb
H <sub>2</sub> O <sub>2</sub>	1.9 ppb
Isoprene	8 ppb
O <sub>3</sub>	30 ppb

*Note.* Reactant mixing ratios reflect initial conditions. Other reactants are initialized at 0 ppb.

spin up on the distributions has been mixed. Because kinks during spin up may affect mean mixing ratios for species with lifetimes longer than a few minutes, we report the ABL-averaged mixing ratios in Figure S1 in Supporting Information S1.

Corresponding to the solar radiation at 12:00 CST (847 W m<sup>-2</sup>), photolysis rates are estimated offline with the NCAR Tropospheric UV and Visible Radiation Model (TUV) 5.0 for clear sky conditions. We scale the estimated photolysis rates using the observed versus estimated photolysis frequency for NO<sub>2</sub> (Table 3). The scaled photolysis rates are fed into the NCAR-LES-MLM-CHEM; they do not vary with time for our short simulations. Aside from inside the canopy, the photolysis rates do not vary vertically. The photolysis rates applied within the canopy are scaled by the fraction of sunlit air space (Figure 1). Our assumption is that the fraction of sunlit canopy air space is the same as the fraction of sunlit leaves, common in chemistry MLMs (e.g., Gao et al., 1993; Moon et al., 2020).

### 2.2.2. Box Modeling

We evaluate the performance of the simplified chemical mechanism using a zero-dimensional photochemical box model (Wang & Pratt, 2017; Wang et al., 2019), which solves the ordinary differential equations using the Back-

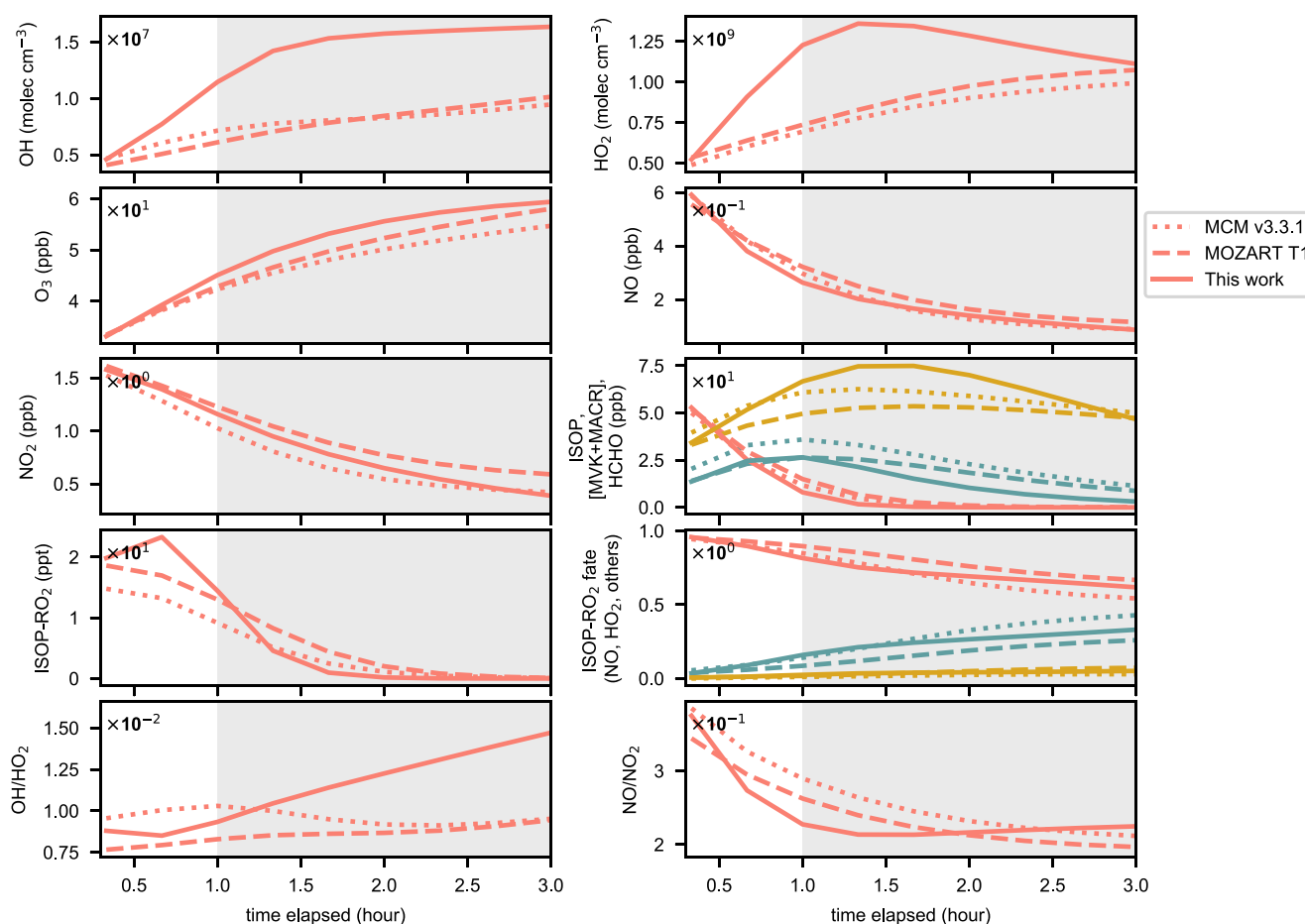
wards Differentiation Formula (also known as the Gear approach; Brown et al., 1989). We compare the time evolution of reactants in our simplified mechanism, the MOZART T1 mechanism (151 gases and 287 reactions; Emmons et al., 2020) and the isoprene mechanism in MCM version 3.3.1 (610 gases and 1,945 reactions; Jenkin et al., 2015). MCM 3.3.1 (hereinafter, MCM) is a near-explicit mechanism widely considered as a benchmark, whereas MOZART T1 (hereinafter, T1) is widely employed in chemical transport modeling.

We test the three mechanisms under “clean” and “polluted” scenarios (Table 4). The clean scenario reflects the lowest mean NO in the vicinity of the canopy across our NCAR-LES-MLM-CHEM simulations and the polluted scenario reflects the highest. In all box-model simulations, photolysis frequencies are calculated using the parameterization in MCM; photolysis frequencies calculated with the solar zenith angle value listed in Table 4 are used (the same solar zenith angle as in the NCAR-LES-MLM-CHEM). The temperature and relative humidity used reflect mean conditions in the vicinity of the canopy in the NCAR-LES-MLM-CHEM simulations.

Although the initial chemical conditions for our NCAR-LES-MLM-CHEM simulations arise from an NCAR-LES-CHEM simulation that has been running for a couple of hours, we integrate the chemical evolution forward in time in the NCAR-LES-MLM-CHEM for no longer than 1 hr. We thus examine the three mechanisms with the box modeling, respectively, over 3 hr (1 hour of which we consider as spin up). In general, the box-model results demonstrate that NO<sub>x</sub> and ozone as well as the first-generation isoprene chemistry are well represented by our mechanism.

We start describing the box-model results for the polluted scenario because simplified mechanisms likely work best at higher NO<sub>x</sub> given that then peroxy radicals react exclusively with NO. For the polluted scenario, our mechanism overestimates OH and HO<sub>2</sub> relative to T1 and MCM (Figure 2). For HO<sub>2</sub>, the overestimate is strongest in the second hour and then declines, agreeing with T1 and MCM at the end of the third hour. OH/HO<sub>2</sub> is overestimated relative to T1 and MCM, implying too rapid HO<sub>x</sub> cycling. Isoprene rapidly decays to negligible levels. Our mechanism overestimates HCHO with the largest overestimate in the second hour and underestimates [MVK + MACR]. Our mechanism captures isoprene peroxy radicals (ISOP-RO<sub>2</sub>) simulated by MCM and that ISOP-RO<sub>2</sub> fate is dominated by reaction with NO. Our mechanism also agrees with T1 and MCM that loss via reaction with HO<sub>2</sub> is second in importance, with the importance increasing with time (and the importance of reaction with NO decreasing in time).

For the clean scenario, our mechanism underestimates OH relative to MCM but agrees with T1 (Figure 3). Our mechanism underestimates HO<sub>2</sub>, with a weaker underestimate relative to T1 than MCM. The bias in OH and HO<sub>2</sub> between MCM and our mechanism increases with time. Likely our mechanism underestimates HO<sub>x</sub> recycling from the RO<sub>2</sub> isomerization pathway (and thus important at low NO) (Jenkin et al., 2015; Wennberg et al., 2018)



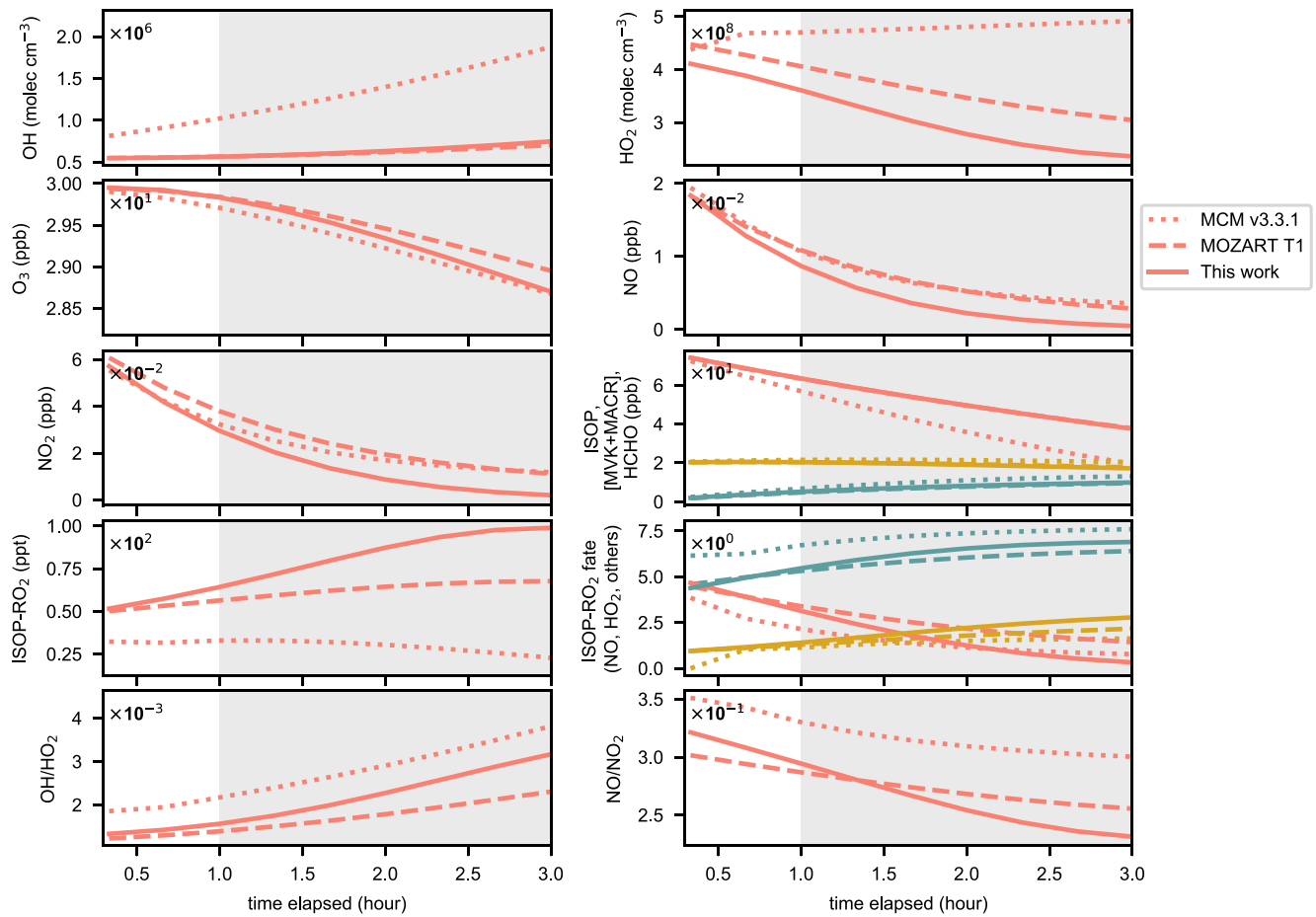
**Figure 2.** Chemical evolution from box-model simulations of the different chemical mechanisms as a function of time since the simulation start for polluted conditions. For plots with multiple quantities, the first quantity listed on the y-axis label is orange, the second is teal, and the third is gold. The first hour is dictated by initial conditions and a few reactions and does not represent the full behavior of the mechanisms so we focus on the final 2 hr, highlighted in gray.

included in MCM, but not in T1 or our mechanism. Nonetheless, even complex mechanisms, including MCM, struggle to capture  $\text{HO}_x$  recycling and disagree with respect to  $\text{HO}_x$  production under low- $\text{NO}_x$  high-isoprene conditions because the chemistry is complex and not well constrained (Bates & Jacob, 2019; Schwantes et al., 2020; Wennberg et al., 2018). For our mechanism,  $\text{OH}/\text{HO}_2$  is larger than T1 but smaller than MCM, implying too slow  $\text{HO}_x$  cycling. In response to low OH in our mechanism relative to MCM, isoprene is overestimated relative to MCM. The similarity in OH between our mechanism and T1 leads to similar isoprene. There is agreement across the mechanisms in [MVK + MACR] and HCHO. Our mechanism overestimates ISOP- $\text{RO}_2$  (the overestimate is larger relative to MCM than T1) but captures that ISOP- $\text{RO}_2$  fate is dominated by reaction with  $\text{HO}_2$ . Compared to MCM, our mechanism (and T1) slightly underestimates ISOP- $\text{RO}_2$  fate by reaction with  $\text{HO}_2$ . Likely the difference follows slower reaction rates in T1 and our mechanism versus MCM (factor of 1.9 at 298 K). Our mechanism agrees with T1 and MCM that reaction with NO is second in importance in terms of ISOP- $\text{RO}_2$  fate in the first 2 hours, but then loss via other reactions become more important.

## 2.3. Biogenic Emissions

### 2.3.1. Leaves

There are only biogenic emissions of isoprene from leaves in our model. Leaf-level emissions of isoprene are parameterized at each vertical layer of the canopy. We leverage the Model of Emissions of Gases and Aerosols



**Figure 3.** Chemical evolution from box-model simulations of the different chemical mechanisms as a function of time since the simulation start for clean conditions. For plots with multiple quantities, the first quantity listed on the y-axis label is orange, the second is teal, and the third is gold. The first hour is dictated by initial conditions and a few reactions and does not represent the full behavior of the mechanisms so we focus on the final 2 hr, highlighted in gray.

from Nature (MEGAN) framework (e.g., Guenther et al., 2006, 2012). Simulated isoprene emission ( $E_{\text{isop}}$ ) at each layer follows:

$$E_{\text{isop}} = \text{LAD} \epsilon_i \gamma_{T_i} \gamma_{\text{PPFD}} \quad (2)$$

The emission capacity ( $\epsilon_i$ ) is a constant equal to  $10,000 \mu\text{g m}^{-2} \text{hr}^{-1}$  for isoprene emission from broadleaf deciduous temperate trees (Guenther et al., 2012). However, we use  $\epsilon_i$  equal to  $3,000 \mu\text{g m}^{-2} \text{hr}^{-1}$  to obtain reasonable mean in-canopy mixing ratios of isoprene (i.e., less than 10 ppbv). As stated by Guenther et al. (2012), there often needs to be a model specific tuning parameter to get reasonable emissions from MEGAN for a particular model (i.e., the canopy environment coefficient), and we interpret our reduction in  $\epsilon_i$  as resulting from that need. The activity factors ( $\gamma$ ) for leaf temperature ( $T_i$ ) and leaf-level photosynthetic photon flux density (PPFD) are calculated separately for sunlit versus shaded leaves and then averaged using the fraction of sunlit leaves at each layer.

The  $T_i$  activity factor ( $\gamma_{T_i}$ ) follows:

$$\gamma_{T_i} = E_{\text{opt}} \frac{C_{T_1} e^{\frac{C_{T_1}}{R} \left( \frac{1}{T_{\text{opt}}} - \frac{1}{T_i} \right)}}{C_{T_2} - C_{T_1} \left( 1 - e^{\frac{C_{T_2}}{R} \left( \frac{1}{T_{\text{opt}}} - \frac{1}{T_i} \right)} \right)} \quad (3)$$

Parameters  $C_{T_1}$  (95 kJ mol<sup>-1</sup>) and  $C_{T_2}$  (230 kJ mol<sup>-1</sup>) are the activation and deactivation energies;  $R$  is the ideal gas constant (kJ K<sup>-1</sup> mol<sup>-1</sup>);  $T_1$  is in units K;  $E_{\text{opt}}$  (1.9; unitless) is the maximum normalized emission capacity;  $T_{\text{opt}}$  (312.5 K) is the  $T_1$  at which  $E_{\text{opt}}$  occurs. Values for  $E_{\text{opt}}$  and  $T_{\text{opt}}$  follow Guenther et al. (1999).

The PPFD activity factor ( $\gamma_P$ ) follows:

$$\gamma_P = C_P \frac{\alpha \text{PPFD}}{\sqrt{1 + \alpha^2 \text{PPFD}^2}} \quad (4)$$

The variable PPFD is in  $\mu\text{mol m}^{-2} \text{s}^{-1}$ ;  $C_P$  (unitless) and  $\alpha$  (mol mol<sup>-1</sup>) follow Guenther et al. (2006):

$$\alpha = 0.004 - 0.0005 \ln(\text{PPFD}_{240}) \quad (5)$$

$$C_P = 0.0468 \text{PPFD}_{240}^{0.6} e^{0.005(\text{PPFD}_{24} - \text{PPFD}_0)} \quad (6)$$

The variable  $C_P$  is the normalized emission capacity at PPFD = 1000  $\mu\text{mol m}^{-2} \text{s}^{-1}$  and  $\alpha$  is the quantum yield of isoprene. The variable  $\text{PPFD}_{240}$  is the average PPFD over the past 240 hr and  $\text{PPFD}_{24}$  is the average PPFD over the past 24 hr. We halve the instantaneous PPFD to estimate  $\text{PPFD}_{240}$  and  $\text{PPFD}_{24}$ —given that we simulate mid-day conditions, halving instantaneous values should approximate daily averages. The constant  $\text{PPFD}_0$  is equal to 200  $\mu\text{mol m}^{-2} \text{s}^{-1}$  for sunlit leaves and 50  $\mu\text{mol m}^{-2} \text{s}^{-1}$  for shaded leaves (Guenther et al., 2006).

### 2.3.2. Soil

Only NO is emitted by the soil in our model. We perform two NCAR-LES-MLM-CHEM simulations: the first case is a lower bound and the second is an upper bound of observed soil NO emissions for temperate deciduous forests. The micrometeorology of the two cases can be considered the same (Table 1). The case high\_soil\_NO has 15 pptv m s<sup>-1</sup> soil NO emissions and low\_soil\_NO has 3 pptv m s<sup>-1</sup>. The slim constraints on chamber soil NO fluxes from deciduous forests in the southeast US range from 0.1 to 15 pptv m s<sup>-1</sup> (Thornton et al., 1997; Valente & Thornton, 1993; Williams & Fehsenfeld, 1991).

We note that NO emitted by the soil faces parameterized transport from the soil surface to 1 m (the lowest grid cell of the LES). Because we did not modify MOST to account for chemical reactivity, no chemistry occurs from 0 to 1 m. Therefore, the mean soil NO source to the model should be regarded as the mean NO<sub>x</sub> flux at 1 m (14.0 pptv m s<sup>-1</sup> in high\_soil\_NO and 2.6 pptv m s<sup>-1</sup> in low\_soil\_NO). We do not tie our soil NO emission to soil moisture and temperature given that they do not change much on the timescales investigated here.

## 2.4. Dry Deposition

### 2.4.1. Description of Framework and Processes in Multilayer Dry Deposition Scheme

Dry deposition is parameterized at each vertical layer of the canopy using a resistance framework, with dry deposition for different gases varying with their molecular weight, reactivity, and solubility, as is common in dry deposition schemes (e.g., Wesely, 1989; Wesely & Hicks, 2000; Zhang et al., 2002). Other than in this broad sense, however, our scheme differs from widely used approaches (we discuss the differences in Section 2.4.2). Below we outline exactly how we calculate dry deposition. In our model, all reactants deposit to leaves and soil. Note that the parameterization for dry deposition described here differs from that described by Clifton and Patton (2021) for ozone because here we need to simulate dry deposition of many species within the same framework.

Leaf deposition terms are calculated separately at each canopy layer for sunlit versus shaded leaves, and then averaged according to the local sunlit versus shaded leaf fraction. Resulting from the need for parallel stomatal and cuticular uptake pathways on the lower side of the leaves but only cuticular uptake on the upper side (see next paragraph), leaf uptake at each height in the canopy (“effective leaf conductance” or  $eg_{\text{leaf}}$  [s<sup>-1</sup>]) follows:

$$eg_{\text{leaf}} = eg_{\text{leaf,up}} + eg_{\text{leaf,lo}} \quad (7)$$

The subscripts “up” and “lo” refer to the upper and lower sides of leaves.

$$eg_{\text{leaf,up}} = \left[ \frac{1}{g_b} + \frac{1}{g_{\text{leaf,up}}} \right]^{-1} \text{LAD} \quad (8)$$

**Table 5**  
The Initial Resistances Used for All Species in the Dry Deposition Scheme

Initial resistances	Value (s m <sup>-1</sup> )
$R_{i,s,cut}$	$1 \times 10^8$
$R_{i,r,cut}$	$1 \times 10^3$
$R_{i,s,m}$	$5 \times 10^5$
$R_{i,r,m}$	$5 \times 10^0$
$R_{i,s,soil}$	$1 \times 10^8$
$R_{i,r,soil,wet}$	$1 \times 10^4$
$R_{i,r,soil,dry}$	$1 \times 10^2$

$$e g_{leaf,lo} = \left[ \frac{1}{g_b} + \frac{1}{g_{leaf,lo}} \right]^{-1} LAD \quad (9)$$

The variable  $g_b$  (m s<sup>-1</sup>) is the conductance for transport of a gas through the leaf boundary layer, and  $g_{leaf}$  (m s<sup>-1</sup>) is the conductance for removal of a gas on the upper or lower side of the leaf.

We consider that the trees only have stomata on one side of their leaves, consistent with broadleaf deciduous vegetation, but have cuticles on both sides of their leaves. The structure of our dry deposition scheme with regards to the upper and lower sides of the leaves follows Monteith and Unsworth (1990) and Nobel (1991).

The variable  $g_b$  is calculated with the same model as heat: the maximum of values for free and forced convection (Leuning et al., 1995; Nikolov et al., 1995):

$$g_{b,H,free} = 0.921 \frac{D_H (\kappa g |T_a - T_l| l_1^3 / \nu^2)^{0.25}}{l_1} \quad (10)$$

$$g_{b,H,forced} = 0.003 \sqrt{\frac{\|\mathbf{u}\|}{l_w}} \quad (11)$$

$$g_b = \left( \frac{D_X}{D_{heat}} \right)^{\frac{2}{3}} \max [g_{b,H,free}, g_{b,H,forced}] \quad (12)$$

The variable  $D_X$  is molecular diffusivity of gas  $X$  in air (m<sup>2</sup> s<sup>-1</sup>);  $D_H$  is the molecular diffusivity for heat in air (m<sup>2</sup> s<sup>-1</sup>);  $\kappa$  is a coefficient for the volumetric thermal expansion of air (K<sup>-1</sup>);  $g$  is the acceleration due to gravity (m s<sup>-2</sup>);  $T_a$  is air temperature (K);  $l_1$  is the leaf length (0.15 m);  $\nu$  is the kinematic viscosity of dry air (m<sup>2</sup> s<sup>-1</sup>);  $\mathbf{u}$  is the wind vector (m s<sup>-1</sup>);  $l_w$  is the leaf width (0.05 m).

For upper and lower sides of the leaf,  $g_{leaf}$  (m s<sup>-1</sup>) follows:

$$g_{leaf,up} = g_{cut} \quad (13)$$

$$g_{leaf,lo} = \frac{1}{\left( \frac{1}{g_s} + \frac{1}{g_m} \right)} + g_{cut} \quad (14)$$

The variable  $g_{cut}$  is the cuticular conductance (m s<sup>-1</sup>), which describes the ability of leaf cuticle to remove the gas via the gas's solubility and reactivity, follows:

$$g_{cut} = f_{wet,cut} \left( \frac{H_X^{eff} R T}{R_{i,s,cut}} + \frac{f_{0,X}}{R_{i,r,cut}} \right) + (1 - f_{wet,cut}) \left( \frac{f_{0,X}}{R_{i,r,cut}} \right) \quad (15)$$

The variable  $H_X^{eff}$  is the effective Henry's Law constant at 298.15 K (M atm<sup>-1</sup>) for gas  $X$ ;  $R$  is the universal gas constant (atm M<sup>-1</sup> K<sup>-1</sup>);  $T = 298.15$  K;  $f_{0,X}$  is the reactivity of gas  $X$  with the surface, which scales from 0 to 1 (unitless). The parameter  $R_{i,s,cut}$  (s m<sup>-1</sup>) is the initial resistance for solubility on cuticles and  $R_{i,r,cut}$  (s m<sup>-1</sup>) is the initial resistance for reactivity on cuticles (Table 5).

The variable  $f_{wet,cut}$  is the fraction of wet leaf cuticle (unitless) and estimated via:

$$f_{wet,cut} = e^{RH_l} / e \quad (16)$$

$e$  is the natural number (Euler's number);  $RH_l$  is leaf-level fractional relative humidity because thin water films form on leaf cuticles as humidity increases (Burkhardt & Eiden, 1994; Burkhardt & Hunsche, 2013; Eiden et al., 1994). The assumption in Equation 16 is that a greater fraction of the leaf becomes covered by the water film as  $RH_l$  increases. We use an exponential dependence because nonstomatal dry deposition of ozone has been

shown to depend on relative humidity exponentially—current understanding is that this dependence is due to thin water film formation on leaf cuticles (Clifton et al., 2020; Sun et al., 2016).

The stomatal conductance ( $g_s$ ) ( $\text{m s}^{-1}$ ) describes the ability of the gas to diffuse through plant stomata, and follows:

$$g_s = \frac{D_x}{D_{\text{H}_2\text{O}}} g_{s,\text{H}_2\text{O}} \quad (17)$$

The parameter  $D_{\text{H}_2\text{O}}$  is the molecular diffusivity of water vapor in air. The variable  $g_{s,\text{H}_2\text{O}}$  is the stomatal conductance for water vapor and follows the Ball-Berry model (Ball et al., 1987):

$$g_{s,\text{H}_2\text{O}} = \left[ g_0 + g_1 \frac{A_{\text{net}}}{c_s} p_s \text{RH}_l \right] \frac{R T_a}{p_s} \quad (18)$$

The parameters  $g_0$  ( $0.01 \text{ mol m}^{-2} \text{ s}^{-1}$ ) and  $g_1$  (9; unitless) are vegetation type specific and empirical;  $A_{\text{net}}$  is net photosynthesis ( $\text{mol m}^{-2} \text{ s}^{-1}$ );  $c_s$  is the carbon dioxide partial pressure at the leaf surface (Pa);  $p_s$  is the atmospheric surface pressure (Pa);  $R$  is the universal gas constant ( $\text{J mol}^{-1} \text{ K}^{-1}$ ). An iterative framework for the carbon dioxide partial pressure inside the leaf couples  $A_{\text{net}}$  to  $g_s$  (see Text S1 of Clifton and Patton (2021) for full details).

The mesophyll conductance ( $g_m$ ) ( $\text{m s}^{-1}$ ) is calculated via:

$$g_m = \frac{H_X^{\text{eff}} R T}{R_{i,s,m}} + \frac{f_{0,X}}{R_{i,r,m}} \quad (19)$$

The variable  $g_m$  describes the ability of the fluids and tissues internal to the leaf (e.g., the mesophyll) to remove the gas via the gas's reactivity and solubility. The parameter  $R_{i,s,m}$  ( $\text{s m}^{-1}$ ) is the initial resistance for solubility inside the leaf and  $R_{i,r,m}$  ( $\text{s m}^{-1}$ ) is the initial resistance for reactivity instead the leaf (Table 5).

The conductance for soil ( $g_{\text{soil}}$ ) ( $\text{m s}^{-1}$ ) follows:

$$g_{\text{soil}} = f_{\text{wet,soil}} \left( \frac{H_X^{\text{eff}} R T}{R_{i,s,\text{soil}}} + \frac{f_{0,X}}{R_{i,r,\text{soil,wet}}} \right) + (1 - f_{\text{wet,soil}}) \left( \frac{f_{0,X}}{R_{i,r,\text{soil,dry}}} \right) \quad (20)$$

The parameter  $R_{i,s,\text{soil}}$  ( $\text{s m}^{-1}$ ) is the initial resistance for solubility on soil,  $R_{i,r,\text{soil,wet}}$  ( $\text{s m}^{-1}$ ) is the initial resistance for reactivity on wet soil, and  $R_{i,r,\text{soil,dry}}$  is the initial resistance for reactivity on dry soil (Table 5). The construction of our model, specifically the use of different initial resistances for  $f_0$  for wet versus dry conditions is based on evidence for ozone that dry deposition to soil decreases with increasing soil moisture (Clifton et al., 2020; Fumagalli et al., 2016; Massman, 2004; Stella et al., 2019). The variable  $f_{\text{wet,soil}}$  is the fraction of wet soil (unitless) and follows:

$$f_{\text{wet,soil}} = \text{vsm}_{0.5\text{m}} \quad (21)$$

The variable  $\text{vsm}_{0.5\text{m}}$  is the volumetric soil moisture at 0.5-m depth in the soil ( $\text{m}^3 \text{ m}^{-3}$ ).

Parameterized transport through the soil boundary layer is combined with parameterized turbulent transport from 1 m to the soil boundary layer. Thus, we do not adjust transport through the soil boundary layer by the gas's diffusivity.

#### 2.4.2. Comparison of Dry Deposition Scheme to Wesely (1989)

The Wesely (1989) dry deposition scheme (hereinafter, Wesely) is the most widely used and recognized dry deposition scheme and has been a key building block for both big leaf and multilayer canopy parameterization development over the past 30 years. Our dry deposition scheme differs from Wesely so we identify the differences to illustrate that regardless of scale, the processes and dependencies are represented in distinct ways:

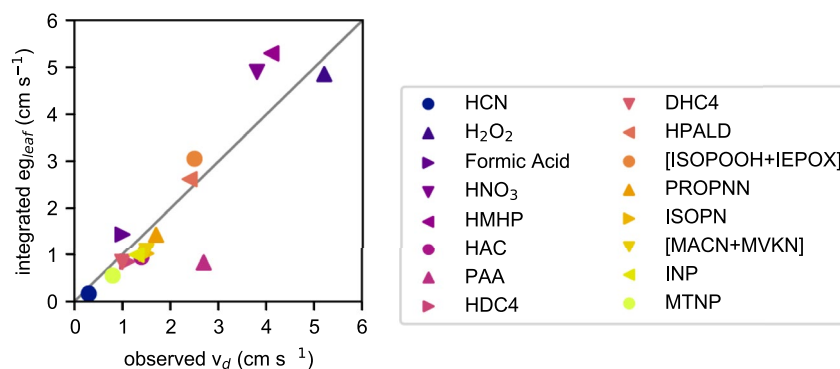
1. Our model is an MLM so we estimate leaf uptake at multiple vertical canopy layers (currently 10) in contrast to Wesely who only estimate leaf uptake at two canopy layers.
2. Turbulence is resolved in our model, which eliminates the need for parameterizations of turbulent transport above or inside the canopy except at the ground where our model has a rough-wall boundary condition.

3. Wesely uses a single bulk term to describe transport through the thin boundary layer around all surfaces. In our model, transport through the boundary layer around surfaces differs for soil and leaves (and for leaves is calculated at each canopy layer). As such, our models for boundary layers around surfaces are also different from Wesely.
4. The original big-leaf Wesely scheme employs seasonal parameters instead of leaf area index (LAI; i.e., LAD integrated over the canopy) while our leaf conductances are scaled by LAD.
5. Our model provides an estimate of carbon dioxide exchange across stomata driven by the carbon supply and demand for photosynthesis, rather than a limited Jarvis (1976) type stomatal conductance model where a prescribed maximum stomatal conductance is scaled by meteorological factors as used in Wesely.
6. Wesely assumes dry deposition on a given surface depends on the gas's solubility and reactivity but does not distinguish between wet versus dry surfaces (i.e., a gas's solubility always increases uptake, even if the part of the leaf is not explicitly wet; e.g., Knote et al., 2015; Hardacre et al., 2021). However, the surface needs to be wet for a gas to dry deposit due to its high solubility. While we can assume that after diffusion through stomata, the surfaces inside the leaf are wet, soil and cuticles are not always wet. In our model, a gas's solubility only impacts uptake to wet surfaces while a gas's reactivity impacts dry deposition to wet and dry surfaces.
7. While Wesely includes adjustments to cuticular uptake in the upper canopy due to leaf wetness from rain and dew, the formation of thin water films on leaves (Burkhardt & Eiden, 1994; Burkhardt & Hunsche, 2013; Eiden et al., 1994) is not represented by Wesely. In our model, we include the formation of thin water films on leaves.
8. Leaf wetness does not reduce stomatal uptake in our model, in contrast to Wesely. Leaf wetness also does not reduce cuticular uptake of ozone, as suggested by Wesely for some cases, as is consistent with evidence showing enhancements of cuticular uptake of ozone on wet leaves (Clifton et al., 2020; Massman, 2004; Sun et al., 2016).
9. In our model, cuticular uptake happens on both sides of the leaf, but stomatal uptake happens only on one side, as consistent with broadleaf deciduous vegetation. In Wesely, any distinction as to whether cuticular uptake happens on a single side or both sides is not explicit. Theoretically, cuticular uptake could happen on both sides, but the reactivity of the cuticle and the water films or droplets on the cuticle may differ between the sides (Clifton et al., 2020). Regardless, we assume that cuticular reactivity and thin water film formation are the same on both sides. Without this assumption, we do not capture the highest  $v_d$  observed across gases for  $\text{H}_2\text{O}_2$ ,  $\text{HNO}_3$ , and HMHP (see Section 2.4.3).
10. Wesely does not consider changes in soil uptake with soil moisture. We use a similar framework for soil uptake as cuticular uptake and we parameterize the fraction of dry versus wet soil with volumetric soil moisture at 0.05-m depth.

In summary, our dry deposition scheme is fundamentally different from Wesely aside from a general resistance framework incorporating molecular weight, solubility, and reactivity influences on uptake. While the exact Wesely scheme determines uptake at two layers in a canopy, our scheme can estimate uptake for several layers in a canopy providing a detailed vertical distribution of dry deposition with key differences in the representation of individual processes.

### 2.4.3. Tuning Strategy for Initial Resistances

The initial resistances in our dry deposition scheme are the same across all gases and need to be tuned. Without a comprehensive flux data set for all the gases in our chemical mechanism, or even constraints on the fluxes of many of them, we turned to one of the most comprehensive reactive flux datasets for a deciduous forest to define the initial resistances: the Nguyen et al. (2015) data set from 5 days in June 2013 at Centerville, Alabama. Nguyen et al. (2015) report fluxes for hydrogen cyanide (HCN), hydrogen peroxide ( $\text{H}_2\text{O}_2$ ), formic acid, nitric acid ( $\text{HNO}_3$ ), hydroxymethyl hydroperoxide (HMHP), hydroxyacetone (HAC), peroxyacetic acid (PAA), the  $\text{C}_4$  hydroxy dicarbonyl from IEPOX oxidation (HDC4), the  $\text{C}_4$  dihydroxy carbonyl from IEPOX oxidation (DHC4), isoprene hydroperoxy aldehydes (HPALD), the sum of isoprene hydroxy hydroperoxides and epoxydiols ([ISOPOOH + IEPOX]), propanone nitrate or propanal nitrate (PROPNN), isoprene hydroxy nitrates (ISOPN), the sum of hydroxy nitrates with carbon backbones of methacrolein and methyl-vinyl ketone ([MACN + MVKN]), isoprene nitrooxy hydroperoxide (INP), and monoterpene nitrooxy hydroperoxide (MTNP).



**Figure 4.** Observed deposition velocity ( $v_d$ ) at Centerville, Alabama versus effective leaf conductance ( $eg_{\text{leaf}}$ ) integrated through the canopy as estimated with our multilayer dry deposition scheme. Daytime average  $v_d$  are from Nguyen et al. (2015) Table S2.

Specifically, we estimated  $eg_{\text{leaf}}$  with our dry deposition scheme for the species in the Nguyen et al. (2015) data set. We tuned the initial resistances needed for our model by comparing  $eg_{\text{leaf}}$  integrated through the canopy (hereinafter, integrated  $eg_{\text{leaf}}$ ) calculated with our scheme to observed  $v_d$  from Nguyen et al. (2015) (see their Table S2 for exact values used). In general, observed  $v_d$  include the influences of not only leaf uptake (i.e.,  $eg_{\text{leaf}}$ ) but also turbulent transport, fast chemical losses, dry deposition to soil, and changes due to nonstationarity. However, Nguyen et al. (2015) suggest that the observed ranking of  $v_d$  across the gases in their data set from high to low values is dominated by leaf uptake, justifying our comparison of integrated  $eg_{\text{leaf}}$  with the observed  $v_d$  as the basis of our tuning.

We estimate  $eg_{\text{leaf}}$  for the gases from Nguyen et al. (2015) using simple Henry's Law constants, diffusivities, and reactivities as defined in Nguyen et al. (2015)'s Table S2 as well as archived vertical profiles of meteorological and biophysical variables from an existing simulation with the NCAR-LES-MLM (Clifton & Patton, 2021) that has similar midday conditions as observed at Centerville during the period of measurements. For example, potential temperature at canopy top is 30.5°C in our simulation within the range of the observed midday virtual temperature around 28–35°C, relative humidity is 61%–62% in our simulation within the range of observed daytime relative humidity of 50%–80%, and friction velocity at canopy top is 0.39 m s<sup>−1</sup> in our simulation and about 0.36–0.46 m s<sup>−1</sup> observed midday. Sensible heat flux at canopy top in our simulation is 133 W m<sup>−2</sup> compared to an observed 140–160 W m<sup>−2</sup> midday and latent heat flux is 422 W m<sup>−2</sup> compared to an observed 450 W m<sup>−2</sup> midday. Our simulation is for a 20-m high homogeneous broadleaf deciduous forest with LAI of 4 m<sup>2</sup> m<sup>−2</sup>, which differs from Talladega National Forest that surrounds the Centerville site on three sides with an LAI of 4.7 m<sup>2</sup> m<sup>−2</sup> and 10-m canopy height. However, similar micrometeorology suggests that LAI and canopy height differences should not preclude development of an empirical dry deposition scheme appropriate for use here.

Figure 4 shows good agreement between the observed  $v_d$  and the integrated  $eg_{\text{leaf}}$  using the initial resistances in Table 5. Our tuning is primarily guided by the range in solubilities across the gases in the observational data set. The parameters related to solubility and leaf uptake are  $R_{i,s,\text{cut}}$  and  $R_{i,s,\text{meso}}$ . The Nguyen et al. (2015) data set has only four gases with non-zero  $f_0$ : H<sub>2</sub>O<sub>2</sub> and HNO<sub>3</sub> ( $f_0 = 1$ ) as well as PAA and HMHP ( $f_0 = 0.1$ ). H<sub>2</sub>O<sub>2</sub>, HNO<sub>3</sub> and HMHP are also highly soluble. Thus, it is challenging to tune based on reactivity relative to solubilities and molecular weights, for which there are more even distributions of different values across the gases. Our offline estimate slightly overestimates dry deposition of HNO<sub>3</sub> and HMHP but underestimates dry deposition of H<sub>2</sub>O<sub>2</sub> (slightly) and PAA (majorly) (Figure 4), suggesting that the reactivity influence is underestimated, and the solubility influence is overestimated across these gases. However, increasing the influence of reactivity would lead to less-than-ideal ozone dry deposition and we thus do not tune the initial resistances further. For example, ozone has  $f_0 = 1$  and very low solubility. With our offline dry deposition scheme employing the parameters values in Table 5, integrated  $eg_{\text{leaf}}$  for ozone is 1.4 cm s<sup>−1</sup> with 58% stomatal uptake. Ozone fluxes were not measured by Nguyen et al. (2015) but observed midday  $v_d$  for ozone span ~0.2–1.4 cm s<sup>−1</sup> at mid-latitude forested monitoring sites (Bauer et al., 2000; Clifton et al., 2017, 2019; Finkelstein et al., 2000; Mikkelsen et al., 2004; Neirynck & Verstraeten, 2018; Turnipseed et al., 2009; Wu et al., 2018), suggesting integrated  $eg_{\text{leaf}}$  falls on the higher end of observations. The daytime stomatal fraction of ozone dry deposition inferred from observations

at forested sites averages 57% (Clifton et al., 2020), and so we capture the stomatal fraction. Because stomatal uptake of ozone is already the highest possible value for a given  $g_s$  (i.e., there no mesophyll resistance for ozone), increasing the influence of the reactivity would increase cuticular uptake of ozone, decreasing the stomatal fraction and most importantly increasing  $eg_{\text{leaf}}$  to unrealistic values, which we wish to avoid.

For the reactants in the chemical mechanism in the NCAR-LES-MLM-CHEM, we use  $f_0$  from the Community Earth System Model (CESM; Emmons et al., 2020; Karl et al., 2010; Schwantes et al., 2020). Table 2 gives  $f_0$  for each reactant. For CO, NO, methane, and isoprene,  $f_0 = 0$ . For HO<sub>2</sub>, NO<sub>2</sub>, and N<sub>2</sub>O<sub>5</sub>,  $f_0 = 0.1$ . For other reactants,  $f_0 = 1$ . Note that while we do not use  $f_0 = 0$  for HPALD and [ISOPOOH + IEPOX] as suggested by Nguyen et al. (2015) and used in our offline tuning, employing  $f_0 = 1$  for these reactants does not change leaf uptake, which is mostly dominated by cuticular uptake due to the moderate-to-high solubilities as well as the high molecular weights that slow diffusion through stomata.

NO<sub>2</sub> reacts inside stomata, but there is not much evidence suggesting cuticular uptake of NO<sub>2</sub> (Chaparro-Suarez et al., 2011; Delaria et al., 2020; Sparks et al., 2001). With  $f_0 = 0.1$ , the stomatal fraction of leaf uptake for NO<sub>2</sub> is 93%. To our knowledge, only Delaria et al. (2020) infers mesophyll resistances for NO<sub>2</sub>—their values are 20–130 s m<sup>−1</sup> for Californian trees. Our choice of  $R_{i,r,\text{meso}}$  allows the mesophyll resistance for NO<sub>2</sub> to fall within this range and stomatal uptake of ozone to be unobstructed by the mesophyll, as supported by laboratory evidence (Laisk et al., 1989; Omasa et al., 2000; Wang et al., 1995).

To parameterize dry deposition to the soil, we started with the same initial resistances for soil uptake as cuticular uptake in our tuning but increased the initial resistance for reactivity on wet soil by a factor of 10, reducing the uptake by surface reaction to wet soil. To get a reasonably high  $g_{\text{soil}}$  for ozone (i.e., 1 cm s<sup>−1</sup> for completely dry soil moisture), we reduced the initial resistance for reactivity on dry surfaces by 10. Our parameterization works so that dry deposition of ozone to soil decreases with soil moisture as reactions on dry soil surfaces outweigh reactions on wet soil surfaces (Clifton et al., 2020).

In the NCAR-LES-MLM-CHEM, we use the effective Henry's Law constant like Wesely to indicate the gas's solubility, in contrast to Nguyen et al. (2015) who suggest using the simple Henry's Law constant. In our effective Henry's Law constants, we consider hydrolysis, but not dissociation of acids and bases. Except for H<sub>2</sub>O<sub>2</sub>, a weak acid, the gases in our mechanism should not have strong protonation tendencies and thus our effective Henry's Law constants can likely adequately describe the solubility.

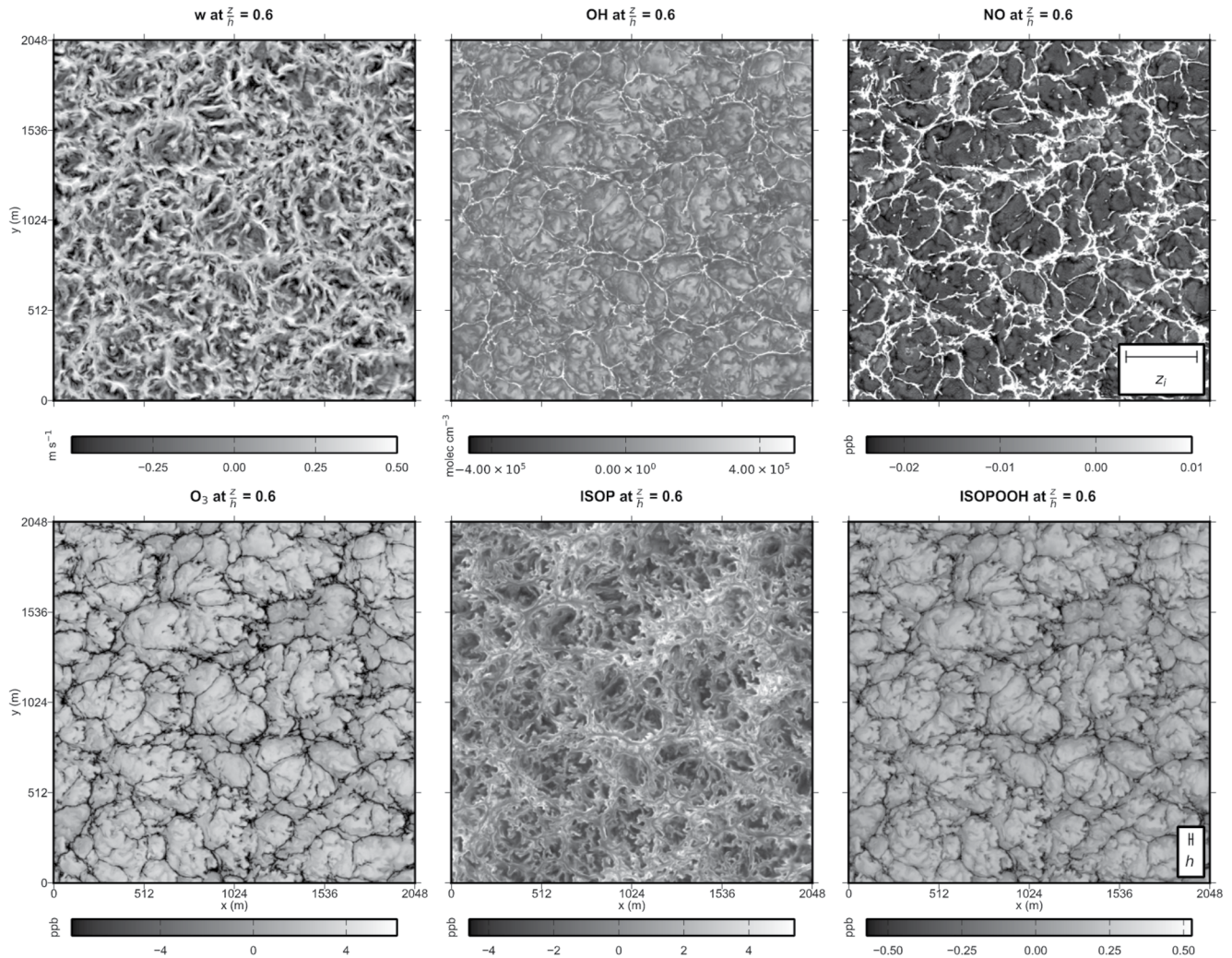
In general, there is a need for further research into the reactivities (e.g., Niinemets et al., 2014) used in classic approaches for modeling dry deposition (e.g., Wesely, 1989; Zhang et al., 2002). The reactivity term in Wesely is based on “the negative log of electron activity for half redox reactions in neutral aqueous solutions and the overall second order reaction rates with S(IV)” “in water with near neutral pH” “for the oxidation of biological substances” (Wesely, 1989). Wesely likely used the reactivity for aqueous reactions because there were more constraints on aqueous reactions at the time and/or the reactivity term was largely designed for inside plant stomata (e.g., mesophyll) where reactions are aqueous. Thus, there are at least three topics that should be addressed in future work. First, whether the basis for the reactivity term of aqueous reactions is appropriate for both wet and dry surfaces should be evaluated. Second, applying the same reactivity to wet versus dry mesophyll, cuticles, and soil may be inappropriate, given that the composition of these surfaces may vary. Third, using redox reactions to define the reactivity does not make sense for many compounds, including the oxidized biogenic compounds examined here, so a critical re-examination of the definition of the reactivity is needed.

## 2.5. Chemical Segregation

We consider chemical segregation in the following framework, with the mean chemical sink of  $X_1$  by reaction with  $X_2$  and reaction rate coefficient  $k$  at a given height following:

$$\left. \frac{\partial \overline{X_1}}{\partial t} \right|_{\text{chem}} = -\overline{k X_1 X_2} = -\overline{k} \left( \overline{X_1 X_2} + \overline{X'_1 X'_2} \right) = -\overline{k X_1 X_2} (1 + I_s) \quad (22)$$

The overbars represent horizontal averages, and the primes represent deviations from the horizontal average. The covariance between  $X_1$  and  $X_2$  ( $\overline{X'_1 X'_2}$ ) reflects similar or dissimilar variability in  $X_1$  and  $X_2$  due to organization in the turbulence transporting  $X_1$  and  $X_2$  and the spatially varying sources and/or sinks of  $X_1$  and  $X_2$ .



**Figure 5.** Instantaneous horizontal slices of vertical velocity ( $w$ ), OH, NO, ozone, isoprene, and ISOPOOH at  $0.6 \times$  canopy height ( $h$ ) for high\_soil\_NO. Values shown are anomalies from the instantaneous horizontal average. Insets show scale for  $h$  and atmospheric boundary layer height ( $z_i$ ) at the time of the slice.

The segregation intensity ( $I_s$ ) reflects the magnitude of the covariance relative to  $\overline{X_1 X_2}$ . Models that assume well-mixed conditions implicitly assume negligible  $\overline{X_1' X_2'}$  and thus negligible  $I_s$ . Positive  $I_s$  means that models assuming complete mixing of reactants underestimate reaction rates while negative  $I_s$  implies an overestimate.

Note that the mean and fluctuating components of  $k$  could be considered in the framework of segregation and the NCAR-LES-MLM-CHEM considers varying  $k$ . However, we expect  $k$  fluctuations to be small relative to reactant fluctuations, and therefore do not examine  $k$  within the framework of segregation, as done by prior studies (e.g., Kim et al., 2016; Krol et al., 2000). Also note that Equation 22 ignores any sub-grid covariance between  $X_1$  and  $X_2$  because we do not have an appropriate model for this.

### 3. Spatial Distribution of Reactants Simulated by the NCAR-LES-MLM-CHEM

#### 3.1. Instantaneous Horizontal Variability

Figure 5 shows instantaneous horizontal variability in OH, NO, ozone, isoprene, and ISOPOOH as well as vertical velocity ( $w$ ) at the height of maximum canopy density in high\_soil\_NO (i.e., a single instance in time in our midday simulation). Features range in size from canopy height ( $h$ ) to the ABL height ( $z_i$ ) for  $w$  and all reactants.

For  $w$ , upward-moving motions occur in thin-filament-like structures and downward-moving motions occur in broad regions of negative  $w$ .

Ozone and ISOPOOH are depleted in the thin upwelling motions, but OH and NO are enriched in them because the canopy is an ISOPOOH and ozone sink but NO source and  $\text{NO} + \text{HO}_2$  is an important OH source (Figure S2 in Supporting Information S1). Horizontal variability in isoprene is particularly distinct compared to the other reactants, closely following horizontal variability in leaf temperature (not shown). Above the canopy at  $2h$ , the structure of the horizontal variability is much larger for  $w$  and all reactants (Figure S3 in Supporting Information S1). The small-scale variations observed at  $0.6h$  (Figure 5) and even more exaggeratedly so at the ground (Figure S4 in Supporting Information S1) are not as clearly present at  $2h$ , highlighting the impact of the canopy on turbulence.

In low\_soil\_NO, some features of OH and NO horizontal variability change relative to high\_soil\_NO (compare Figure S5 in Supporting Information S1 with the upper right two panels of Figure 5). For example, the filament-like upwelling motions with enriched NO are slightly thinner in low\_soil\_NO at  $0.6h$ . Additionally, horizontal variability in OH at  $0.6h$  tends to have larger structure in low\_soil\_NO (like the horizontal variability in ISOPOOH and ozone in high\_soil\_NO in Figure 5). In contrast to high\_soil\_NO, OH is not enriched in the thin upwelling motions which contain high NO, rather OH is depleted in low\_soil\_NO. The difference between cases likely results from enhanced secondary OH production in high\_soil\_NO (Figure S2 in Supporting Information S1).

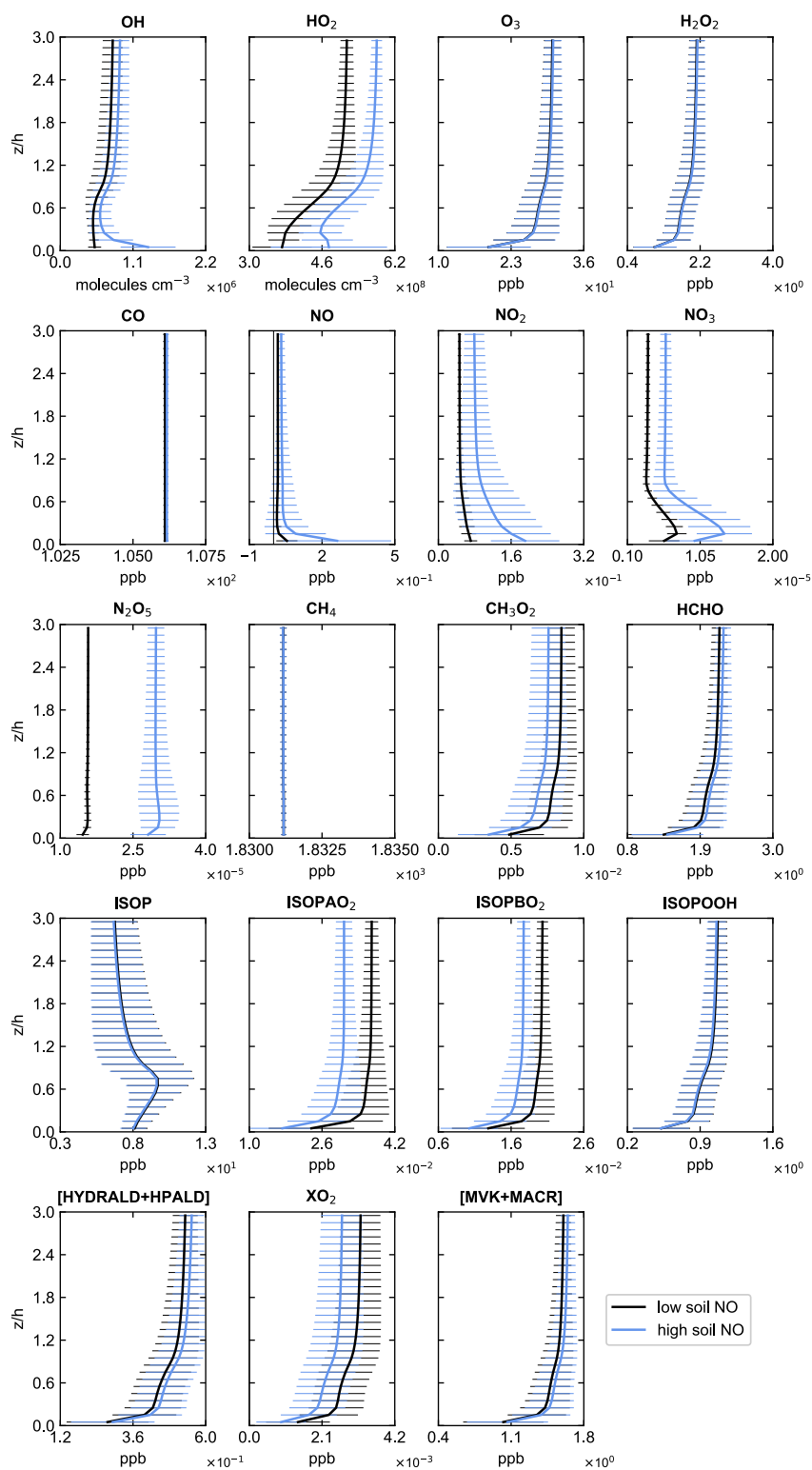
### 3.2. Mean Vertical Profiles

Figure 6 shows the vertical profiles of the mean abundances for the reactants in the chemical mechanism. The effect of the canopy is apparent in most of the profiles. Most profiles are relatively constant above  $h$ , except isoprene, which decreases above  $h$ . The reactants with canopy-bottom minima in their profiles are  $\text{HO}_2$ ,  $\text{O}_3$ ,  $\text{H}_2\text{O}_2$ ,  $\text{N}_2\text{O}_5$ , and the reactive carbon oxidation products ( $\text{CH}_3\text{O}_2$ , HCHO, ISOPAO<sub>2</sub>, ISOPBO<sub>2</sub>, ISOPOOH, XO<sub>2</sub>, [HYDRALD + HPALD], and [MVK + MACR]). The oxidation products maximize above  $h$ . We expect the maxima in the oxidation products to occur above  $h$  because the canopy is an efficient sink of these compounds through dry deposition. Isoprene's maximum happens at  $0.8h$ , where isoprene emissions maximize (not shown), which is slightly above maximum LAD at  $0.6h$  but where the sunlight fraction of leaves exceeds the shaded fraction of leaves (Figure 1). Horizontal mean isoprene is maximum 9.7 ppbv, which is on the higher side of daytime hourly isoprene mixing ratios observed at temperate deciduous forests (Goldstein et al., 1998; Helmig et al., 1998; McKinney et al., 2011; Wei et al., 2020). Both NO and NO<sub>2</sub> peak at the ground, but NO<sub>3</sub> and N<sub>2</sub>O<sub>5</sub> peak just above the ground in the lower canopy. Elevated NO<sub>3</sub> below canopy relative to above canopy was also observed at a mixed forest and was attributed to lower below-canopy NO<sub>3</sub> photolysis (Schulze et al., 2017).

Differences with soil NO emissions are also apparent for many reactants. In high\_soil\_NO relative to low\_soil\_NO, the abundances of NO, NO<sub>2</sub>, NO<sub>3</sub>, N<sub>2</sub>O<sub>5</sub>, OH, and HO<sub>2</sub> are higher but CH<sub>3</sub>O<sub>2</sub>, ISOPAO<sub>2</sub>, ISOPBO<sub>2</sub>, and XO<sub>2</sub> are lower. Other species have similar abundances between cases. Differences between the cases in N<sub>2</sub>O<sub>5</sub>, HO<sub>2</sub>, CH<sub>3</sub>O<sub>2</sub>, ISOPAO<sub>2</sub>, ISOPBO<sub>2</sub>, and XO<sub>2</sub> are largely consistent with height, whereas for OH, NO, NO<sub>2</sub>, and NO<sub>3</sub> differences are strongest in the canopy. In high\_soil\_NO, OH strongly peaks at the ground, and there is a strong decrease in OH upwards from the ground to a local minimum mid-canopy. In low\_soil\_NO, OH weakly peaks above  $h$  and is relatively constant in the lower canopy. In both cases, OH increases with height within the canopy above  $0.6h$  but becomes constant above  $h$ . High OH near the ground in high\_soil\_NO follows high local secondary OH production (Figure S2 in Supporting Information S1).

The horizontal mean NO/NO<sub>2</sub> ratio is constant above  $h$  around 0.36–0.41 in both cases but varies substantially within the canopy (Figure S6 in Supporting Information S1), another line of evidence that the chemistry is very different at different heights of the canopy. In low\_soil\_NO, the ratio varies from 0.23 to 0.77 inside the canopy. In high\_soil\_NO, the ratio varies from 0.33 to 1.37 inside the canopy. Thus, with higher soil NO, more NO<sub>x</sub> is NO inside the canopy. In both cases, the ratio peaks at ground due to NO emission at the soil surface and decreases from the ground to the lower canopy with increasing NO + O<sub>3</sub> forming NO<sub>2</sub> and little NO<sub>2</sub> photolysis in the shaded canopy. The ratio then increases from the lower to upper canopy with increasing NO<sub>2</sub> photolysis and less shading.

Horizontal variability in the reactants also varies with height in the canopy. For reactants other than CO and methane, which do not vary much relative to the magnitude of their mixing ratios, horizontal variability in reactant



**Figure 6.** Vertical profiles of the horizontal mean mixing ratios for the reactants in the chemical mechanism (concentrations for OH and HO<sub>2</sub>) from the ground to 3  $\times$  canopy height ( $h$ ). Error bars indicate one standard deviation (horizontal variability). Thin gray line on NO panel denotes  $x = 0$ .

abundances is largest inside the canopy (Figure 6 and Figure S7 in Supporting Information S1). For isoprene, horizontal variability is largest at  $h$ . There is also a peak in the coefficient of variation (standard deviation divided by mean at same height) for OH at  $h$  in both cases (with a primary peak at the ground in high\_soil\_NO). Otherwise, reactants mostly have horizontal variability maxima near or at the ground. NO has the strongest horizontal variability, with the peak in its variability in the lower canopy (coefficient of variation is up to 173% in high\_soil\_NO and up to 113% in low\_soil\_NO). NO<sub>2</sub> also has large horizontal variability compared to other gases in high\_soil\_NO—the coefficient of variation for NO<sub>2</sub> is up to 70% (in contrast with up to 28% in low\_soil\_NO).

NO variability is also highly skewed, causing the standard deviation error bars on the horizontal mean NO mixing ratio in high\_soil\_NO to span zero on Figure 6 (note that NO never attains negative values in the simulations). Edburg et al. (2012) also find high skewness for a passive scalar emitted by the ground, compared to passive scalars emitted by or deposited to the canopy. Like Edburg et al. (2012), we find that NO skewness maximizes at canopy top (not shown). Maximal skewness near canopy top likely results from shear-driven motions generated near  $h$  intermittently penetrating the canopy and transporting NO upwards (Edburg et al., 2012).

## 4. Vertical Reactant Flux Profiles Simulated by the NCAR-LES-MLM-CHEM

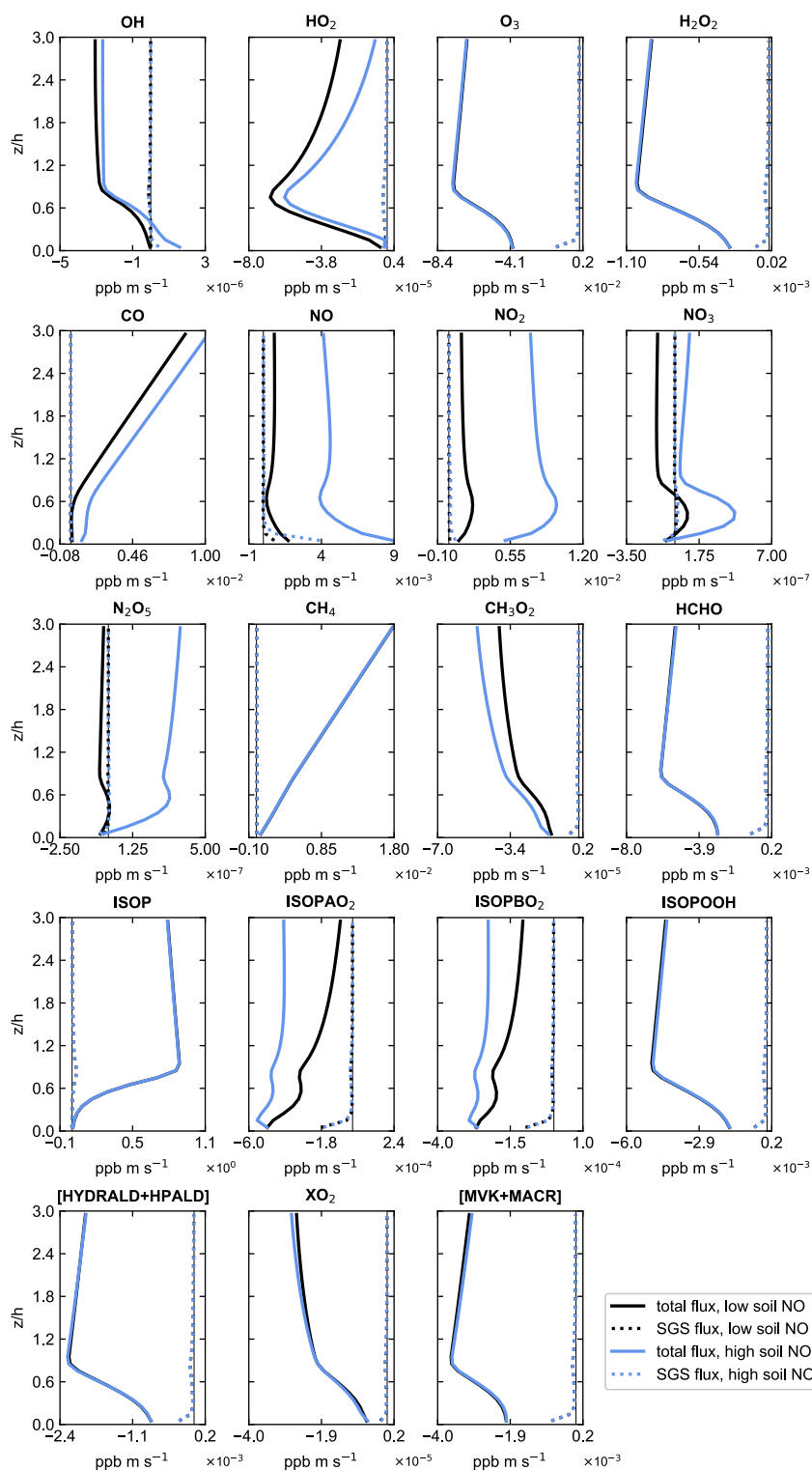
### 4.1. Total Fluxes

Figure 7 shows profiles of the vertical total fluxes for all reactants in the mechanism. A vertical reactant flux profile reflects the net vertical transport of the reactant by turbulence; this transport can be altered by chemistry depending on the rate at which chemistry produces or destroys the reactant compared to the rate at which turbulent transport occurs. The total fluxes are the sum of the resolved and sub-grid scale (SGS) fluxes. The sign of the fluxes at  $h$  indicates whether the canopy layer acts as a net source or sink of the reactant. The canopy layer is a net sink of reactive carbon oxidation products and ozone but net source of isoprene, NO, and NO<sub>2</sub>.

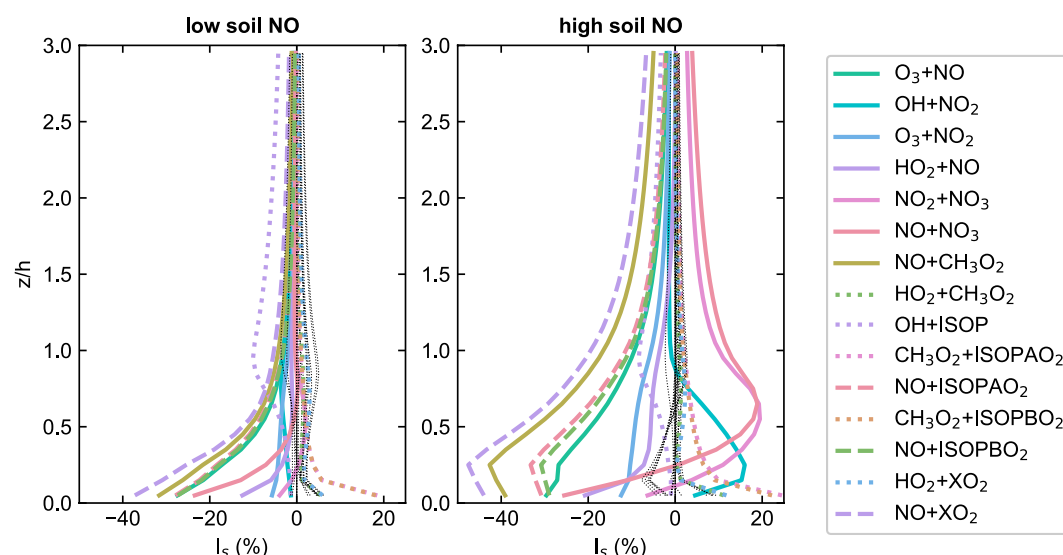
Below  $h$ , there are variations in the fluxes of many species with height. Above  $h$ , most flux profiles change linearly with height, as expected for reactants with lifetimes longer than the time it takes for turbulence to mix the ABL. For example, methane and CO are very well mixed, and the positive vertical fluxes simply reflect entrainment of air with lower methane and CO abundances from the FT into the ABL. Interestingly, OH, which has a very short lifetime compared to the turbulence mixing time, exhibits nearly linear flux profiles with height above  $h$ . The linear OH flux profile above the canopy suggests that local OH production compensates for the inability for turbulence to transport OH before it reacts. For ISOPAO<sub>2</sub> and ISOPBO<sub>2</sub>, the fluxes are negative within the canopy and increasingly become closer to zero with increasing height, with the maximum negative flux at the ground. We expect that fast chemistry between ISOPAO<sub>2</sub> or ISOPBO<sub>2</sub> and NO and/or HO<sub>2</sub> contributes to the high downward fluxes at the ground. The CH<sub>3</sub>O<sub>2</sub> and XO<sub>2</sub> fluxes are negative from the ground to  $3h$ . The CH<sub>3</sub>O<sub>2</sub> and XO<sub>2</sub> flux profiles linearly increase with height, indicating that fast losses are even greater above  $h$  than below.

Comparing the soil NO cases, the fluxes are substantially different for NO, NO<sub>2</sub>, NO<sub>3</sub>, N<sub>2</sub>O<sub>5</sub>, CH<sub>3</sub>O<sub>2</sub>, ISOPAO<sub>2</sub>, and ISOPBO<sub>2</sub>, but only slightly different for OH, HO<sub>2</sub>, CO, and XO<sub>2</sub>. However, for OH and CO, as well as NO<sub>3</sub> and N<sub>2</sub>O<sub>5</sub>, there are differences between cases in the sign of the flux, at least at some heights.

The only study reporting HCHO eddy covariance fluxes at a forest show positive daytime fluxes (24 pptv m s<sup>−1</sup> at noon at a coniferous forest in Colorado during August 2010; DiGangi et al., 2011). Simulated HCHO fluxes are negative, contrasting with the positive observed fluxes. Negative HCHO fluxes follow large sinks at soil and leaf surfaces in our model due to the high solubility and reactivity of HCHO. While there is laboratory and field evidence for HCHO depositional sinks (Filella et al., 2006; Gray et al., 2014; Kramshög et al., 2018; Li, Su, et al., 2016; Seco et al., 2008; Trowbridge et al., 2020), there is also field evidence for leaf and soil HCHO emissions (DiGangi et al., 2011; Trowbridge et al., 2020). While the presence of large HCHO sinks at canopy and soil surfaces in our model may be justifiable, the sinks as parameterized may overestimate loss because we do not account for desorption (Li, Su, et al., 2016). Too high HCHO dry deposition in our model may contribute to too low HCHO mixing ratios in the canopy. For example, observed afternoon HCHO at a mixed forest in upper Michigan during 8 days in August 2008 are maximum 4.2 ppb above  $h$  in the afternoon (Galloway et al., 2012). While observed HCHO mixing ratios are as low as 2 ppb during the afternoon (comparable to our simulated mixing ratios), the higher observed levels of HCHO mostly correspond to the higher isoprene levels simulated by our model. Better understanding of HCHO sources and sinks in the canopy will advance our ability to accurately represent HCHO in chemistry MLMs (e.g., Wei et al., 2021; Wolfe et al., 2011).



**Figure 7.** Vertical profiles of the vertical total and sub-grid scale fluxes for the reactants in the chemical mechanism from the ground to 3 x canopy height ( $h$ ). Thin gray lines denote  $x = 0$ .



**Figure 8.** Vertical profiles of the intensities of segregation ( $I_s$ ) for all the reactions in the chemical mechanism from the ground to 3  $\times$  canopy height ( $h$ ). Dotted lines denote reactions that do not contain NO or  $\text{NO}_2$ . Color lines denote reactions with substantial segregation intensities in either of the cases (absolute value  $>10\%$  at least one height within the ground to  $3h$ ) or OH + ISOP which we discuss in the text. Other reactions are shown in thin black lines.

## 4.2. Sub-Grid Scale Flux Contributions

A parameterization in the NCAR-LES-MLM-CHEM is needed to estimate SGS vertical transport. Generally, the impact of the SGS flux on the total flux is minor except at the ground (Figure 7). At the ground, the NO SGS flux is 45% of the total flux. For other gases with non-zero total fluxes at the ground except  $\text{NO}_3$  and  $\text{N}_2\text{O}_5$ , the SGS fluxes are less than 40% of the total fluxes at the ground. For  $\text{NO}_3$  and  $\text{N}_2\text{O}_5$ , near-surface SGS fluxes are up to 185% of the total fluxes in high\_soil\_NO, suggesting that  $\text{NO}_3$  and  $\text{N}_2\text{O}_5$  transport there is poorly estimated. Nonetheless, the total  $\text{NO}_3$  and  $\text{N}_2\text{O}_5$  fluxes at the ground are relatively small compared to elsewhere, and the flux divergences are small compared to net chemistry so the poor estimation likely does not matter.

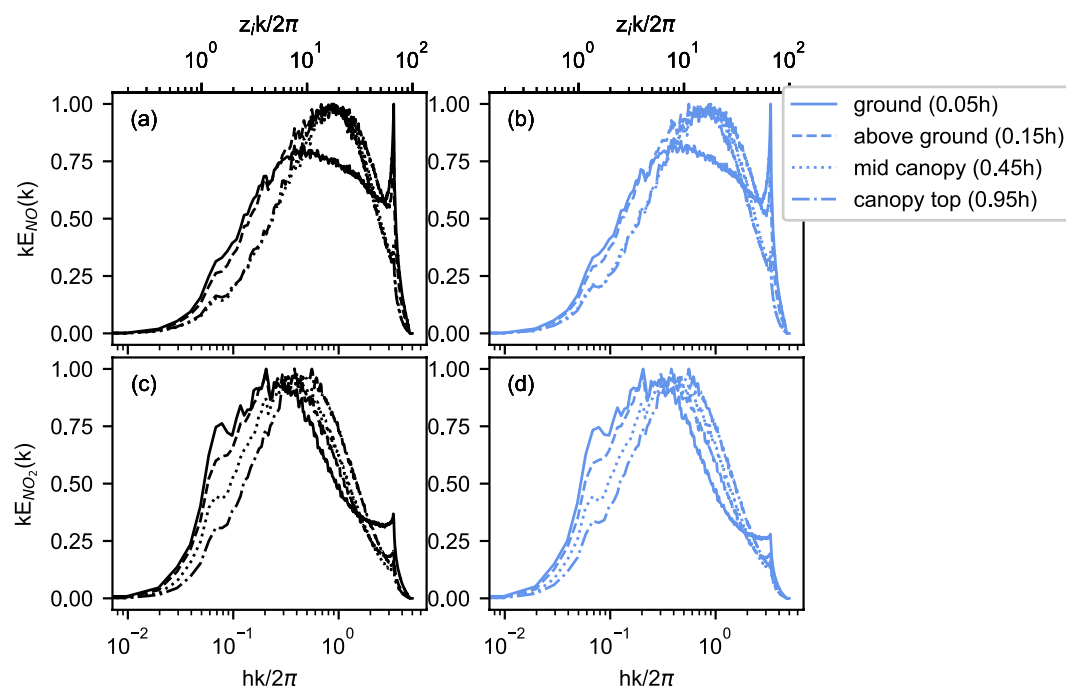
## 5. Chemical Segregation Inside and Just Above the Forest Canopy Simulated by the NCAR-LES-MLM-CHEM

### 5.1. Changes With Height and High Versus Low Soil NO Emissions

Figure 8 shows  $I_s$  for all the reactions in our mechanism. At least some degree of segregation (i.e.,  $|I_s| > 1\%$  at some point from 0 to  $3h$  in one of the cases) occurs in each reaction except OH + CO and OH +  $\text{CH}_4$ . Across all reactions, the highest  $I_s$  from 0 to  $3h$  occur within the canopy. In high\_soil\_NO,  $I_s$  range from  $-19\%$  to  $+13\%$  across reactions above  $h$  but from  $-48\%$  to  $+23\%$  below. In low\_soil\_NO,  $I_s$  range from  $-10\%$  to  $+4\%$  above  $h$  but from  $-38\%$  to  $+19\%$  below. Comparing low\_soil\_NO to high\_soil\_NO, fewer reactions have substantial in-canopy  $I_s$ . In low\_soil\_NO, large  $I_s$  primarily occur in the lower canopy, whereas in high\_soil\_NO, large  $I_s$  occurs in both the lower and upper canopy.

Many reactions have some degree of segregation at the ground, whether positive or negative. Strong segregation at the ground is likely due to shear-driven organized motions generated near the top of the canopy intermittently penetrating the depth of the canopy coupled with differing spatially varying sources and/or sinks between the reactants. In low\_soil\_NO, no reactions have positive  $I_s$  greater than  $+5\%$  aside from at the ground. In high\_soil\_NO, only three reactions have positive  $I_s$  greater than  $+10\%$  above the ground:  $\text{NO} + \text{NO}_3$ ,  $\text{NO}_2 + \text{NO}_3$ , and OH +  $\text{NO}_2$ .

In high\_soil\_NO, the most negative  $I_s$  in the canopy (and the largest absolute  $I_s$  out of all reactions) occurs for the reactions  $\text{NO} + \text{XO}_2$  and  $\text{NO} + \text{CH}_3\text{O}_2$  (ranging from  $-16\%$  to  $-47\%$  in the canopy).  $I_s$  for the reactions  $\text{NO} + \text{XO}_2$  and  $\text{NO} + \text{CH}_3\text{O}_2$  maximizes at  $0.25h$ .  $I_s$  is also substantially negative for the reactions  $\text{NO} + \text{ISOPBO}_2$ ,  $\text{NO} + \text{ISOPAO}_2$ , and  $\text{O}_3 + \text{NO}$  in the canopy. In low\_soil\_NO, negative  $I_s$  is largest above  $0.75h$  for reaction between OH and isoprene. Below  $0.75h$ , the reactions with the largest  $I_s$  are  $\text{NO} + \text{XO}_2$ ,  $\text{NO} + \text{CH}_3\text{O}_2$ ,  $\text{O}_3 + \text{NO}$ ,



**Figure 9.** Azimuthally- and time-averaged two-dimensional spectra for NO (a, b) and NO<sub>2</sub> (c, d) at four heights in the canopy for low\_soil\_NO (a, c) and high\_soil\_NO (b, d). Because we aim to compare the shapes of the spectra across heights and not the absolute variability at a given height, each spectrum is normalized by its peak. Spectra are multiplied by horizontal wavenumber  $k$  such that the area under the curve is proportional to the total variance. The bottom  $x$ -axis is such that  $x = 1$  is the scale of the canopy height ( $h$ ), and the top  $x$ -axis is such that  $x = 1$  is the scale of the ABL height ( $z_i$ ).

NO + ISOPAO<sub>2</sub>, NO + ISOPBO<sub>2</sub>, and NO + NO<sub>3</sub>. For these reactions except NO + NO<sub>3</sub>,  $I_s$  increases substantially from low\_soil\_NO to high\_soil\_NO but remains the same sign. Excluding the reaction NO + NO<sub>3</sub>, the reactions with high negative segregation in both cases are reactions between NO, which is emitted from the soil, and species that have large canopy sinks (i.e., O<sub>3</sub>, XO<sub>2</sub>, CH<sub>3</sub>O<sub>2</sub>, ISOPAO<sub>2</sub>, ISOPBO<sub>2</sub>; Figure 7).

For reaction between OH and isoprene,  $I_s$  is negative and up to  $-8\%$  in high\_soil\_NO and  $-9\%$  in low\_soil\_NO. In contrast to most other reactions,  $I_s$  for OH and isoprene maximizes at  $h$ . This is likely due to the peak in horizontal variability in isoprene at  $h$ ; other reactants mostly have absolute peaks in their variability in the lower canopy (Figure S7 in Supporting Information S1).

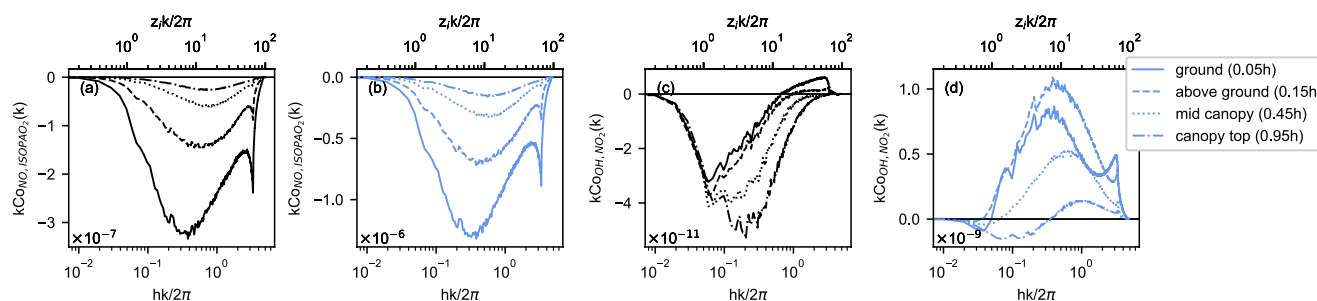
In summary, segregation is highest in the canopy for heights up to  $3h$ , and specifically at the ground for most reactions. Segregation is particularly important for reactions involving NO and NO<sub>2</sub>. Reactions involving NO have greater than  $\pm 10\%$   $I_s$  somewhere in the canopy. Reactions involving NO<sub>2</sub> have greater than  $\pm 10\%$   $I_s$  somewhere in the canopy in high\_soil\_NO. The segregation of reactions involving NO and NO<sub>2</sub> mostly increase when soil NO emissions reflect the upper bound of observed values for temperate forests.

## 5.2. Understanding Segregation Changes With Height and High Versus Low Soil NO Emissions

Differences in  $I_s$  with species, height, and soil NO emissions only provide vague clues as to the controlling scales and thus processes. Given that we find that segregation is particularly important for reactions involving NO and NO<sub>2</sub>, here we examine variability in NO and NO<sub>2</sub> and covariability in select reactions involving NO or NO<sub>2</sub> using spectral analysis in order to better understand the scales and processes controlling segregation.

### 5.2.1. Variability in NO and NO<sub>2</sub> Across Scales

Figure 9 shows the spectra of NO and NO<sub>2</sub> at different heights in the canopy. We already know that there are differences in the magnitude of variability in NO and NO<sub>2</sub> between the soil NO cases and with height (Figure 6 and Figure S7 in Supporting Information S1) so we normalize each spectrum by its highest value to isolate how the shape of the spectra, or the spatial character of variability, changes between cases and with height.



**Figure 10.** Azimuthally- and time-averaged two-dimensional co-spectra for the reactions  $\text{NO} + \text{ISOPAO}_2$  and  $\text{OH} + \text{NO}_2$  at four heights in the canopy for low\_soil\_NO (a, c) and high\_soil\_NO (b, d). Co-spectra are multiplied by horizontal wavenumber  $k$  such that the area under the curve is proportional to the total covariance. The bottom  $x$ -axis is such that  $x = 1$  is the scale of the canopy height ( $h$ ), and the top  $x$ -axis is such that  $x = 1$  is the scale of the ABL height ( $z_i$ ). The true  $y$ -values of the data are the numbers on the  $y$ -axes multiplied by the numbers in the lower left corner.

Between the soil NO cases, there is very little difference in the normalized NO and  $\text{NO}_2$  spectra for a given height. There are, however, differences in the normalized NO and  $\text{NO}_2$  spectra with height. First, there are thin spectral peaks at very high wavenumbers (i.e., corresponding to very small-scale processes) at and near the ground. Figure 9 shows that very small-scale processes are an important component of the variance for NO (less so for  $\text{NO}_2$ ). Second, the spectra for NO and  $\text{NO}_2$  near the ground are shifted to larger scales compared to higher in the canopy.  $\text{NO}_x$  variability near the ground may be more strongly influenced by the constant NO emission at the surface, which corresponds to the largest length scales in the domain. Thus, higher in the canopy, the spectral peaks shifting to smaller scales likely reflect a more dominant influence of fast chemistry on  $\text{NO}_x$  variability there, with a weakened role for NO emissions at the ground (and very small-scale processes).

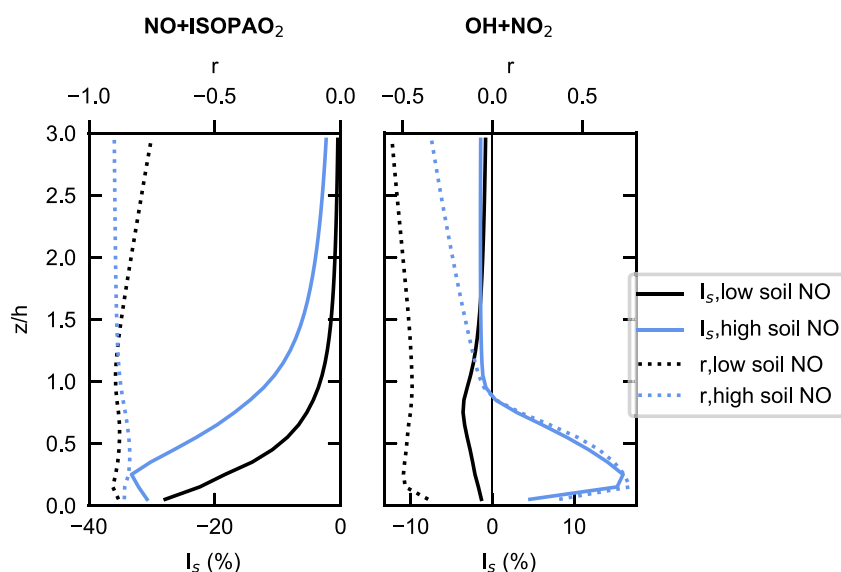
That the normalized spectra are similar in shape between the soil NO cases at a given height suggests that the spatial character of the variation in NO and  $\text{NO}_2$  remains the same between the cases. Similar spatial character in NO and  $\text{NO}_2$  between cases despite changes in segregation in NO and  $\text{NO}_2$  reactions implies that changes in the magnitude of variation in NO and  $\text{NO}_2$  (Figure 6 and Figure S7 in Supporting Information S1) and/or how the NO and  $\text{NO}_2$  change the spatial character or magnitude of variability in other quantities drives the differences in segregation between cases.

### 5.2.2. Select Reactions With High Segregation

Here we investigate  $I_s$  for select reactions by leveraging context from the co-spectra between reactants as well as the correlation coefficients between reactants and the individual reactant standard deviations. Note that the covariance between two quantities can be considered as the product of their correlation coefficient and the standard deviations in each quantity. We focus on two reactions that are important for photochemistry and have high segregation:  $\text{OH} + \text{NO}_2$  and  $\text{NO} + \text{ISOPAO}_2$ . Figure 10 shows the absolute co-spectra at four heights in both cases (we show absolute co-spectra because we aim to discern how changes alter  $I_s$ ). Figure 11 shows  $I_s$  and correlation coefficients for the reactions in both cases.

For  $\text{NO} + \text{ISOPAO}_2$ , we present evidence that changes in  $I_s$  with NO are largely due to the magnitude of NO variability. Changes with height are driven by the magnitude of NO and  $\text{ISOPAO}_2$  variability. Our findings for the reaction  $\text{NO} + \text{ISOPAO}_2$  are the same as for  $\text{NO} + \text{ISOPBO}_2$ ,  $\text{NO} + \text{CH}_3\text{O}_2$ ,  $\text{NO} + \text{XO}_2$ , and  $\text{NO} + \text{O}_3$  (not shown), implying that there may be some generalizability about in-canopy segregation of reactant emitted at the ground versus reactants largely deposited to the canopy. Previous work suggests but does not show with interactive chemistry in an LES MLM that this may be the case (Edburg et al., 2012; Patton et al., 2001).

For  $\text{OH} + \text{NO}_2$ , changes in  $I_s$  with NO and height are influenced by changes in the magnitude of variability and/or spatial character of OH, as well as the magnitude of  $\text{NO}_2$  variability. Generally, we find that low segregation above the canopy stems from low reactant variability, rather than a poor correlation between reactants (the correlations can be moderate to strong). Although it is challenging to determine what exactly drives changes in the correlation between reactants or in the magnitude of variability or spatial character of a given reactant, this analysis demonstrates that constraints on the distributions of reactants across scales in the canopy can aid in building mechanistic understanding of segregation changes with environmental conditions.



**Figure 11.** Vertical profiles from the ground to  $3 \times$  canopy height ( $h$ ). Solid lines show intensities of segregation ( $I_s$ ) for the reactions  $\text{NO} + \text{ISOPAO}_2$  and  $\text{OH} + \text{NO}_2$ . Dotted lines show correlation coefficients ( $r$ ) between reactants.

We first discuss the thin co-spectral peaks at high wavenumbers (i.e., that correspond to very small-scale processes). The peaks are especially apparent at the ground (Figure 10). Because there is high segregation at the ground, and we do not know what drives these peaks, we probe whether they drive the high segregation at the ground by calculating  $\overline{X_1 X_2}$  with low pass filtering of the reactant mixing ratios corresponding to an eddy turnover time of 10 s (a couple seconds longer than the very small-scale scales) for the reactions  $\text{NO} + \text{ISOPAO}_2$  and  $\text{OH} + \text{NO}_2$ . At the ground,  $\overline{X_1 X_2}$  calculated with low pass filtering is within 1%–2% of  $\overline{X_1 X_2}$  calculated without low pass filtering. Thus, while these peaks may contribute to segregation, they do not drive the high segregation.

#### 5.2.2.1. $\text{NO} + \text{ISOPAO}_2$

For the reaction  $\text{NO} + \text{ISOPAO}_2$ ,  $I_s$  is negative and peaks in the lower canopy. The  $\text{NO}$ - $\text{ISOPAO}_2$  correlations do not have strong height variations, suggesting that vertical standard deviation profiles of  $\text{NO}$  and/or  $\text{ISOPAO}_2$  (Figure S7 in Supporting Information S1) drive the vertical  $I_s$  profiles. Both reactants' variability peaks in the lower canopy, and thus both could contribute to height variation in  $I_s$ .

In terms of differences between cases,  $I_s$  is more strongly negative in high\_soil\_NO relative to low\_soil\_NO, with the difference decreasing with height. Together with the  $I_s$  difference between cases despite largely similar correlations between the cases, the similar-in-shape co-spectra at a given height between the cases (Figure 10) imply that the higher  $I_s$  from low\_soil\_NO to high\_soil\_NO stems from stronger reactant variability in high\_soil\_NO, rather than a different relationship between  $\text{NO}$  and  $\text{ISOPAO}_2$  between the cases. While the standard deviations for  $\text{ISOPAO}_2$  are slightly larger in high\_soil\_NO than low\_soil\_NO above the ground (by factor of 1.03–1.1), the increase in the standard deviation of  $\text{NO}$  (by a factor of 5–6) likely dominates the increase in  $I_s$  (Figure S7 in Supporting Information S1).

There may be more variability in  $\text{NO}$  with higher soil  $\text{NO}$  emissions simply because there is more  $\text{NO}$ . However, the coefficient of variation for  $\text{NO}$  increases substantially from low\_soil\_NO to high\_soil\_NO (from 15%–113% to 44%–173%), suggesting that the change in  $\text{NO}$  variability is not a simple function of the change in the amount of  $\text{NO}$ . The  $\text{NO}/\text{NO}_2$  ratio changes such that more  $\text{NO}_x$  is  $\text{NO}$  in high\_soil\_NO (Figure S6 in Supporting Information S1), which may play a role increasing relative  $\text{NO}$  variability.

#### 5.2.2.2. $\text{OH} + \text{NO}_2$

For the reaction  $\text{OH} + \text{NO}_2$ ,  $I_s$  below  $h$  is very different between cases, both in terms of sign and magnitude. Above  $h$ ,  $I_s$  is small and positive in both cases. In this section, we first discuss in-canopy changes with height in

low\_soil\_NO, then in-canopy changes with height in high\_soil\_NO, and finally differences between the cases at  $h$  and then in the canopy.

In low\_soil\_NO,  $I_s$  is negative and small below  $h$ , with a peak in the upper canopy. The highest  $I_s$  in the upper canopy likely reflects OH variability maximizing there (Figure S7 in Supporting Information S1) because the OH-NO<sub>2</sub> correlation is largely constant throughout the canopy (Figure 11) and NO<sub>2</sub> variability maximizes in the lower canopy (Figure S7 in Supporting Information S1). Additional evidence for a driving role for OH on changes in  $I_s$  with height for this case is that the co-spectra only increase with height at scales smaller than  $z_i$ , implying that the cause of the changes with height occurs at smaller scales (consistent with a role for OH because OH variability is largely driven by fast chemistry).

In high\_soil\_NO,  $I_s$  is positive below  $h$  and peaks strongly in the lower canopy. The co-spectra are mostly positive with the highest values just above the ground (but not at the ground), consistent with  $I_s$  maximizing in the lower canopy. The consistency of changes in height in  $I_s$  and the correlation suggests that changes in the relationship between OH and NO<sub>2</sub> with height drives changes in segregation with height in this case, but the vertical profile of variability in NO<sub>2</sub> (Figure S7 in Supporting Information S1) may also play a role. Regardless, this begs the question: why are the correlation and NO<sub>2</sub> variability strongest in lower canopy? To which, the answer may have to do with NO becoming a more important source of OH in high\_soil\_NO in the lower canopy (Figure S2 in Supporting Information S1).

At  $h$ , the co-spectra are different between cases. For example, on high\_soil\_NO at  $h$ , the co-spectrum is positive (like other heights) at small scales but is negative at large scales, which causes the overall correlation to become negative there. In low\_soil\_NO, the co-spectrum at  $h$  is negative at all scales. Despite differences between the cases in the co-spectra at  $h$  (as well as the correlation),  $I_s$  is similar and low, which may reflect reductions in OH and/or NO<sub>2</sub> variability at and above  $h$ .

Between the cases, large differences in shape and sign of the co-spectra are likely driven by differences in OH variability. There are not only strong changes in the magnitude of the OH standard deviation, but also the vertical profile of the OH standard deviation below  $0.75h$  (Figure S7 in Supporting Information S1). In high\_soil\_NO, the OH standard deviation strongly peaks at the ground (in low\_soil\_NO the OH standard deviation peaks weakly at  $h$ ). Together with the difference in the co-spectra between cases (and the correlation) despite similar NO<sub>2</sub> spatial character, the difference in the OH standard deviation vertical profile suggests a change in the OH spatial character as the underlying cause of the strong sign changes in OH-NO<sub>2</sub> segregation between cases.

We hypothesize that changes in the spatial character of OH between the soil NO cases result from the reaction HO<sub>2</sub> + NO becoming a much more important source of OH (Figure S2 in Supporting Information S1). Evidence includes a complete change in correlation coefficient between OH and NO in the canopy (Figure S8 in Supporting Information S1). The OH-NO correlation goes from negative and weak in low\_soil\_NO to positive and strong in high\_soil\_NO, which may be the result of an important source of OH from NO in high\_soil\_NO.

## 6. Comparison to Previous Work and Implications for Atmospheric Chemistry

In the context of key reactions for photochemistry, the reactions initiating oxidation are OH + CO and OH + CH<sub>4</sub> and OH + ISOP. While OH + CO and OH + CH<sub>4</sub> have very low  $I_s$ , negative  $I_s$  between OH and isoprene implies that chemistry following oxidation slows down relative to well-mixed conditions. The vertical profile and the magnitude of OH-ISOP  $I_s$  simulated here are similar to the Patton et al. (2001) estimates with simple proxies for OH and isoprene in a previous version of the NCAR-LES-MLM. However, our estimates are on the lower end of the Dlugi et al. (2010) estimates of  $-1$  to  $-15\%$  in the roughness sublayer of the forest (specifically, 7 m above canopy height) that employ 5-s measurements during the day at mixed forest in Germany during July 2003. Our estimates are also much lower than previous work estimating OH-ISOP  $I_s$  of up to  $-13\%$  in a clear ABL using airborne measurements (Butler et al., 2008; Kaser et al., 2015). Other estimates with LES or airborne measurements suggests that heterogeneous isoprene emissions and clouds can alter OH-ISOP  $I_s$  (Brosse et al., 2018; Karl et al., 2007; Kaser et al., 2015; Kim et al., 2016; Krol et al., 2000; Li, Barth et al., 2016; Ouwersloot et al., 2011), but only clear-sky conditions and a canopy with homogenous structure and composition are investigated here.

In terms of changes with NO<sub>x</sub>, we find that segregation between OH and isoprene is slightly stronger with lower soil NO. To our knowledge, only two studies probe changes in OH-ISOP  $I_s$  with changes in NO<sub>x</sub>. With

the Dutch Atmospheric Large-Eddy Simulation (DALES) simulating the ABL above a typical tropical forest, Ouwersloot et al. (2011) investigate the impact of an order of magnitude change in surface NO emissions (originally  $0.5 \text{ ppt m s}^{-1}$ ), finding a very low impact on OH-ISOP  $I_s$ . However, in another set of sensitivity simulations, Ouwersloot et al. (2011) find increases in OH-ISOP  $I_s$  (factor of 1.6–1.7) with FT  $\text{NO}_2$  increasing from 0 to 0.5 ppb. As simulated by the NCAR-LES-CHEM with updated version of the Horowitz et al. (2003) MOZART2.2 chemical mechanism by Kim et al. (2012) configured for a summer day in the southeast US, OH-ISOP  $I_s$  is negative and less than 10% near the surface but weakens (i.e., becomes closer to zero) as  $\text{NO}_x$  increases but remains less than 5 ppb (Kim et al., 2016). The  $I_s$  increase with lower  $\text{NO}_x$  is attributed to lower OH and thus higher isoprene and stronger variability (Kim et al., 2016). While we also find that  $I_s$  slightly increases with lower  $\text{NO}_x$ , isoprene variability is very similar between our soil NO cases. Thus, changes in OH variability or the OH-ISOP correlation likely drive differences. Future work will examine the causes and impacts of differences of OH-ISOP across our simulations.

Other reactions that are key for photochemistry include reactions between peroxy radicals ( $\text{XO}_2$ ,  $\text{CH}_3\text{O}_2$ ,  $\text{ISOPAO}_2$ ,  $\text{ISOPBO}_2$ ) and NO as well as reaction between OH and  $\text{NO}_2$ . Most previous work examining segregation examines isoprene and OH, but we find large negative  $I_s$  between peroxy radicals and NO, which implies a slowing down of the cycling of peroxy radicals with  $\text{HO}_x$  relative to well-mixed conditions. We also find segregation between OH and  $\text{NO}_2$ , which alters terminal loss of both  $\text{HO}_x$  and  $\text{NO}_x$ . When NO is lower, we find negative segregation between  $\text{NO}_2$  and OH, which implies that there may be more  $\text{NO}_2$  to photolyze to NO and ozone compared to well-mixed conditions. When NO is higher, there is positive segregation between  $\text{NO}_2$  and OH, implying accelerated terminal loss and perhaps reduced ozone production relative to well-mixed conditions.

To estimate the impact of segregation on chemistry, we perform box-model simulations with our simplified mechanism similar to the simulations performed in Section 2.2.2. In a second simulation, we scale the rate coefficients by  $(1 + I_s)$  for all reactions, with  $I_s$  simulated by the NCAR-LES-MLM-CHEM. We use  $I_s$  from 0.35h because segregation for most reactions maximizes around this height. Inserting the  $I_s$  from high\_soil\_NO into the polluted scenario (Figure S9 in Supporting Information S1), we find  $-3$  to  $3\%$  relative changes in OH, but lower  $\text{HO}_2$  (by  $-10.4\%$  to  $-19.7\%$ ) and higher OH/ $\text{HO}_2$  (by  $+8.5$  to  $+28\%$ ). Ozone is lower by up to  $-2\%$  (up to  $-1.2$  ppb) when accounting for segregation. On the other hand, inserting the  $I_s$  from low\_soil\_NO into the clean scenario (Figure S10 in Supporting Information S1), we find  $+2.4$  to  $+4.8\%$  relative increases in OH and lower but then higher  $\text{HO}_2$  ( $-4.2$  to  $+3\%$ ) with mostly higher OH/ $\text{HO}_2$  (up to  $+9.4\%$ ). Ozone is the same between simulations. In summary, while in both clean and polluted scenarios, segregation speeds up  $\text{HO}_x$  cycling, substantially, the effects on OH,  $\text{HO}_2$ , and ozone are different and can vary in sign between the scenarios. In general, differences in OH,  $\text{HO}_2$ , and ozone between the box model simulations with and without segregation implies that segregation may impact interpretations of observations and measurement-model comparisons.

## 7. Comparison to Previous Work and Implications for Ozone Dry Deposition

Observed ozone fluxes measured near the canopy top are commonly interpreted as dry deposition only but may be influenced by fast chemical losses of ozone such as the  $\text{NO} + \text{O}_3$  reaction (Finco et al., 2018; Kramm et al., 1995; Plake et al., 2015). Negative  $\text{NO-O}_3$   $I_s$  means that the influence of NO on the ozone flux is even weaker than under the assumption of well-mixed conditions. To our knowledge, Edburg et al. (2012) is the only study to estimate  $\text{NO-O}_3$  segregation in a canopy; however, they ignored the influence of chemistry. Using a previous version of the NCAR LES-MLM, with a sparser canopy and smaller domain, Edburg et al. (2012) estimate segregation between a scalar emitted from the ground (a proxy for NO) and a scalar deposited to the canopy (a proxy for ozone), finding  $I_s$  up to  $-8\%$  in the canopy, with the largest segregation in the lower canopy and negligible above the canopy. Thus, in-canopy segregation is much larger in our simulations (up to  $-30\%$  in the canopy), which likely result from interactive chemistry, interactive sources and sinks, and the larger domain size that allows ABL scales to influence turbulent transport in the canopy (Patton et al., 2016). The peak in the  $I_s$  vertical profile in Edburg et al. (2012) is also broader and occurs higher in the canopy than in our simulations; in addition to the other differences between our model configurations, this may also be due to the sparser canopy in Edburg et al. (2012) leading to more well-mixed air near the ground.

## 8. Conclusion

Previous work documents the NCAR-LES (e.g., Moeng, 1984; Moeng & Wyngaard, 1988; Sullivan et al., 1996) including the version coupled to a multilayer canopy model (Patton et al., 2016). We present a new version of the NCAR multilayer canopy LES coupled to a simplified chemical mechanism. We show that our mechanism can reproduce the features of first-generation isoprene chemistry and describe how we leverage a multi-species flux data set to parameterize dry deposition for all the gases in the mechanism. For the first time with both a vertically resolved and coupled forest and interactive chemistry in an LES, we demonstrate height-varying vertical profiles of reactant mixing ratios and fluxes as well as chemical segregation inside and above the canopy. We find large segregation for many reactions, in particular for reactions involving NO. Mostly reactant segregation is highest in the lower canopy and lowest above the canopy. With two simulations with differing soil NO emissions, we find that segregation generally becomes more important when soil NO emissions increase to the upper bound of observed values at temperate deciduous forests. For some reactions, segregation switches sign between low and high soil NO cases.

We use spectral analysis to gain a more complete understanding of the factors controlling reactant segregation. We show that not only are changes in the magnitude of reactant variability important for changes in segregation with soil NO emissions, but changes in the spatial character of reactants are important as well. However, capturing segregation accurately in our model requires a full understanding of reactant variability on fast timescales. Thus, there is a need for coupling model development with short-term field campaigns with simultaneous multi-height high-frequency observations of a variety of reactive gases' abundances and micrometeorology, as well as observational constraints on leaf and soil sources and sinks of reactants.

We close by emphasizing that our segregation estimates are closely tied to the atmospheric configuration and chemical mechanism. Changes to the atmospheric configuration will change the character of the organized structures within the turbulence, and changes to the chemical mechanism will alter the pathways by which species are transformed. With access to sufficient computational time, the model introduced here can be expanded to include more reactants and/or a more complex chemical mechanism and simulate other environments (e.g., different atmospheric stability, forest type, chemical composition, and clouds). Such simulations will be key for advancing predictive ability of turbulence-chemistry interactions at regional-to-global scales.

## Data Availability Statement

Processed data from the LES is provided in the NCAR Geoscience Data Exchange repository (<https://doi.org/10.5065/ky81-f733>).

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