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Key Points:

- Strong EMIC waves were observed by Van Allen Probe A in association with phase space density dips over the similar L shell extent
- FIREBIRD-II observed electron precipitation from ~250 keV up to ~1 MeV over the similar L shell extent of EMIC waves but at a later MLT
- Quasi-linear theory predicts efficient precipitation at > 0.8–1 MeV but requires other mechanisms to explain the subrelativistic one

Supporting Information:

- Supporting Information S1
- Figure S1
- Figure S2
- Figure S3

Correspondence to:

L. Capannolo and W. Li,
luisacap@bu.edu;
wenli77@bu.edu

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Direct Observation of Subrelativistic Electron Precipitation Potentially Driven by EMIC Waves

L. Capannolo¹ , W. Li¹ , Q. Ma^{1,2} , L. Chen³ , X.-C. Shen¹, H. E. Spence⁴ , J. Sample⁵ , A. Johnson⁵ , M. Shumko⁵ , D. M. Klumpar⁵, and R. J. Redmon⁶ 

¹Center for Space Physics, Boston University, Boston, MA, USA, ²Department of Atmospheric and Oceanic Sciences, University of California, Los Angeles, CA, USA, ³William B. Hanson Center for Space Sciences, University of Texas at Dallas, Richardson, TX, USA, ⁴Institute for the Study of Earth, Oceans, and Space, University of New Hampshire, Durham, NH, USA, ⁵Department of Physics, Montana State University, Bozeman, MT, USA, ⁶Cooperative Institute for Research in Environmental Sciences, University of Colorado Boulder, Boulder, CO, USA

Abstract Electromagnetic ion cyclotron (EMIC) waves are known to typically cause electron losses into Earth's upper atmosphere at >~1 MeV, while the minimum energy of electrons subject to efficient EMIC-driven precipitation loss is unresolved. This letter reports electron precipitation from subrelativistic energies of ~250 keV up to ~1 MeV observed by the Focused Investigations of Relativistic Electron Burst Intensity, Range and Dynamics (FIREBIRD-II) CubeSats, while two Polar Operational Environmental Satellites (POES) observed proton precipitation nearby. Van Allen Probe A detected EMIC waves (~0.7–2.0 nT) over the similar L shell extent of electron precipitation observed by FIREBIRD-II, albeit with a ~1.6 magnetic local time (MLT) difference. Although plasmaspheric hiss and magnetosonic waves were also observed, quasi-linear calculations indicate that EMIC waves were the most efficient in driving the electron precipitation. Quasi-linear theory predicts efficient precipitation at >0.8–1 MeV (due to H-band EMIC waves), suggesting that other mechanisms are required to explain the observed subrelativistic electron precipitation.

Plain Language Summary Plasma waves in the Earth's magnetosphere can alter the trajectory of particles traveling along geomagnetic field lines. Specifically, electromagnetic ion cyclotron (EMIC) waves can interact with both electrons and protons and cause them to fall into the upper atmosphere of Earth. From past studies and theories, it is known that EMIC waves drive precipitation of ultrarelativistic (>~MeV) electrons and tens to hundreds of keV protons. Such electron precipitation can lead to atmospheric changes and potentially aid ozone depletion. In this work, we show a direct observation of electron precipitation from ~250 keV up to ~1 MeV, potentially driven by EMIC waves using multipoint measurements primarily from a CubeSat mission (FIREBIRD-II) and Van Allen Probes. Quasi-linear calculations indicate that EMIC waves are efficient in driving the electron precipitation at >0.8–1 MeV, but other mechanisms are needed to explain the observed electron precipitation down to ~250 keV. Our study also highlights the capabilities of FIREBIRD-II studying not only microbursts, but also other precipitation patterns (e.g., driven by EMIC waves).

1. Introduction

Electromagnetic ion cyclotron (EMIC) waves occur in the Earth's magnetosphere, typically below the hydrogen, helium, and oxygen ion gyrofrequencies (f_{cH} , f_{cHe} , and f_{cO}), with frequencies of 0.1–5 Hz (Erlandson & Ukhorskiy, 2001; Fraser et al., 2010). As EMIC waves mainly propagate along geomagnetic field lines, they are observable at ground as Pc1–Pc2 waves (Engebretson et al., 2008, 2018), although they can spread from the original source region due to ionospheric ducting (Engebretson et al., 2015; Kim et al., 2010). EMIC waves are generally excited during enhanced geomagnetic activity, when ring current ions are injected from the magnetotail (Fraser et al., 2010; Jordanova et al., 2008), thus occur more frequently from postnoon to dusk in the inner magnetosphere (Allen et al., 2016; Blum et al., 2017; Min et al., 2012).

EMIC waves scatter ~MeV electrons and ~10s–100s keV protons into the loss cone via cyclotron resonance, potentially leading to losses of radiation belt electrons and ring current ions (Cao et al., 2016; Jordanova et al., 2008; Jordanova & Miyoshi, 2005; Li et al., 2007; Meredith et al., 2014; Summers et al., 2017;

Summers & Thorne, 2003). While precipitating protons can produce proton aurora (Miyoshi et al., 2008; Yuan et al., 2010), ~MeV electrons can interact with atmospheric nitrogen and hydrogen oxides, leading to ozone reduction (e.g., Meraner & Schmidt, 2018). The efficacy of EMIC waves in scattering MeV electrons and ring current protons has been confirmed in multiple studies, both observationally and theoretically (Blum et al., 2015; Capannolo et al., 2018, 2019; Hirai et al., 2018; Qin et al., 2018, 2019; Shekhar et al., 2017; Woodger et al., 2018; Yuan et al., 2018). However, the techniques and the satellites used have limitations that still leave a few open questions, for example, on the minimum energy of electrons (E_{\min}) that can be scattered into the loss cone. Many previous studies indicate that only ~MeV electrons can be scattered by EMIC waves (Chen et al., 2011; Meredith et al., 2003; Ni et al., 2015; Summers & Thorne, 2003). However, more recent studies, such as nonresonant interactions, suggest that the electron precipitation can extend to sub-MeV energies (Chen et al., 2016). Although some studies reported sub-MeV electron precipitation events potentially driven by EMIC waves (Capannolo et al., 2019; Clilverd et al., 2015; Hendry et al., 2017), a direct observation of subrelativistic electron precipitation due to EMIC waves is still limited. The often used National Oceanic and Atmospheric Administration (NOAA) Polar Operational Environmental Satellites (POES), in fact, only detect electrons in integral energy channels. Therefore, even if precipitation in the >300-keV channel occurred (Capannolo et al., 2019; Hendry et al., 2017), it is difficult to determine E_{\min} of the precipitating electrons, unless by using some assumptions and modeling. In addition, proton contamination often affects POES lower-energy channels and, despite removal (e.g., Peck et al., 2015), precipitation can still sometimes be misinterpreted. Among the satellites available so far, the Focused Investigations of Relativistic Electron Burst Intensity, Range and Dynamics (FIREBIRD-II) CubeSats provide an excellent opportunity to analyze E_{\min} of precipitating electrons driven by EMIC waves, given their high energy resolution.

Our study focuses on the electron precipitation (from ~250 keV to ~1 MeV) observed by FIREBIRD-II CubeSats. By taking full advantage of available multipoint measurements, we discuss several indicators suggesting that the subrelativistic electron precipitation is potentially caused by EMIC waves. We follow up the observations with quasi-linear modeling to determine if EMIC waves can drive the observed electron precipitation.

2. Observations

Figures 1a and 1b show the electron observations by FIREBIRD-II (hereafter simply FIREBIRD) on 9 July 2017 over 2–3 UT. The twin FIREBIRD 1.5 U CubeSats (Flight Units, FU) were launched on 31 January 2015 as a continuation of the original FIREBIRD mission, operational from December 2013 to June 2014 (Klumpar et al., 2015; Spence et al., 2012). Previous studies have been successful in investigating characteristics and spatial/temporal structure of electron microbursts (Breneman et al., 2017; Crew et al., 2016; Shumko et al., 2018). However, FIREBIRD has also the potential of studying other precipitation patterns, such as EMIC-driven precipitation as shown in this letter. The two units measure electron precipitation in five differential channels and one integral channel, along a 632×433 -km orbit with 99° inclination and a ~90-min orbital period. Count rates are converted to differential fluxes at the central (lowest) energy for the differential (integral) energy channels using a GEometry And Tracking (GEANT4) simulation. GEANT4 also shows minimal proton contamination, significant only for > MeV protons. Figures 1i and 1ii present the L shell and magnetic local time (MLT) values (solid lines) of FU4 and FU3 for different magnetic field models (T89 in black (Tsyganenko, 1989) and T05 in blue (Tsyganenko & Sitnov, 2005)). The MLT values for T89 and T05 overlap. Panels 1iii show the high-resolution electron flux color coded with energy. FU4 and FU3 crossed the same region within ~11 min and both observed electron precipitation across all energy channels, indicating that it occurred not only for relativistic electrons but also for subrelativistic ones, at least down to ~250 keV. While FIREBIRD does not have a fixed pointing direction, it is reasonable to assume that the observations in this case correspond to precipitating electrons rather than trapped ones, given the FIREBIRD location West of the South Atlantic Anomaly (lat~ -66° , lon~ -148° , within a region of quasi-trapped particles as in Greeley et al. (2019), Figure 1), and no other electron or proton enhancements west of FIREBIRD (NOAA-19 at ~2:30 UT, MLT ~ 13–14, not shown), which excludes the presence of drifting electrons from west.

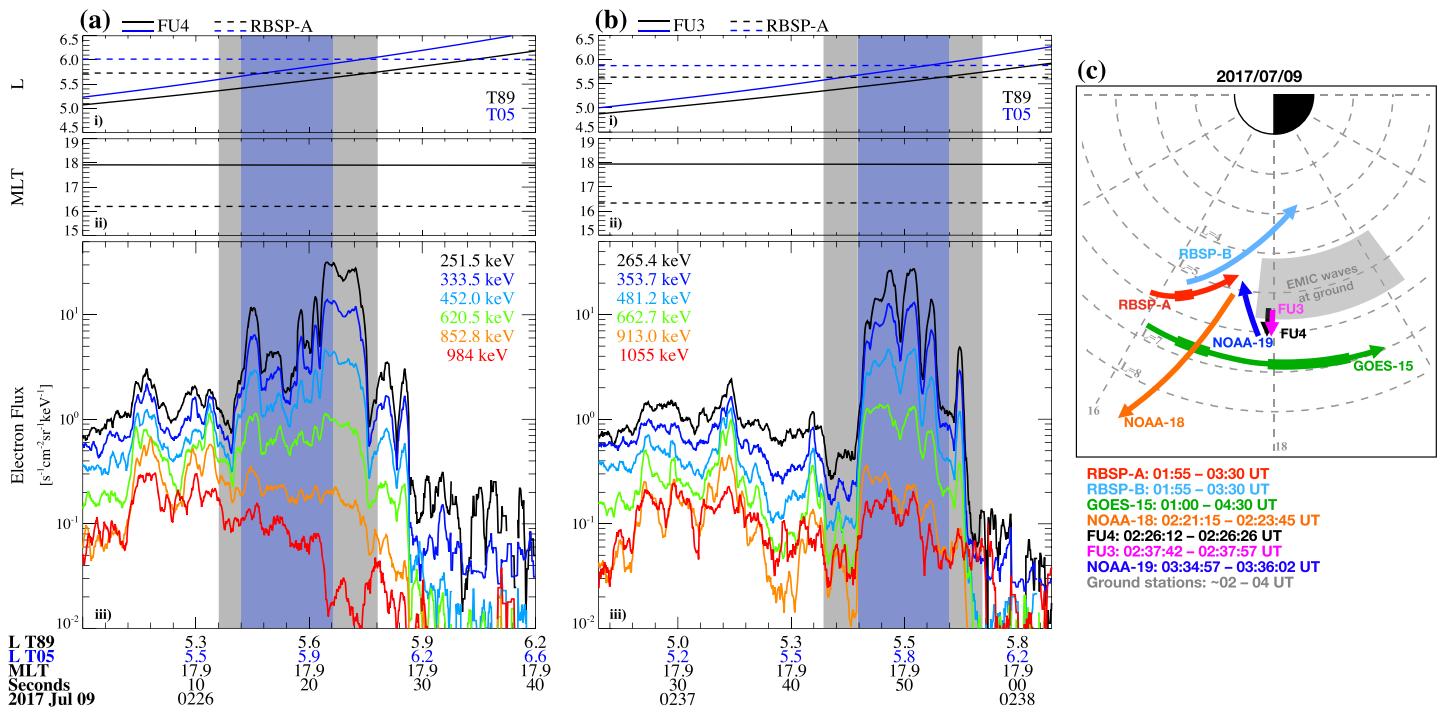


Figure 1. FIREBIRD-II observations: FU4 (a) and FU3 (b). (i) L shell, (ii) MLT for FU (solid line) and RBSP-A (dashed line); The black (blue) line indicates the values from the T89 (T05) model, as marked in the x axes; (iii) electron fluxes color coded with energy channels from the 50-ms high-resolution data, smoothed for better visualization with a 450-ms window. The gray (blue) shaded block shows the T89 (T05) L shell extent of EMIC waves observed by RBSP-A mapped onto the FIREBIRD trajectories. The x axes values refer only to the FIREBIRD units. (c) Color-coded trajectories of the satellites observing the event and location of the collection of ground-based stations observing EMIC waves (gray region) mapped onto the L (T89)-MLT plane. Thicker lines on RBSP-A and GOES-15 indicate the intervals when EMIC waves were detected (~02:30–02:55 UT, and ~01:30–2:00 UT and 3:00–04:10 UT). The legend shows the time intervals of the trajectories plotted; for the POES satellites, the time stamps indicate when proton precipitation was observed; for FU, the time stamps refer to the UT intervals of the gray blocks in (a) and (b).

Figure 1c shows a schematic overview of the trajectories of the satellites used in this study and the collection of the ground-based observations onto the L-MLT plane. Both FU were in close L shell conjunction with Van Allen Probe A (RBSP-A, dashed lines, Figures 1i and 1ii) at ~02:26:26 and ~02:37:54 UT respectively, with $\Delta\text{MLT} \sim 1.6$. Observations by RBSP-A (Figure 2) are useful to understand the driver of the precipitation observed by FIREBIRD. Van Allen Probes have orbital period, perigee and apogee of ~9 hr, ~600 km, and ~6 R_E , respectively (Mauk et al., 2013). The event occurred during a strong substorm ($\text{AL} \sim -900$ nT, Figure 2a) but over a nonstorm period ($\text{Dst} > -20$ nT). Over 02:00–03:30 UT, RBSP-A was traveling inbound within the plasmasphere (density $\sim 130 \text{ cm}^{-3}$, Figure 2b, inferred from the upper hybrid resonance frequency observed by the High-Frequency Receiver, HFR; Kurth et al., 2015) on the dusk side. We use Electric and Magnetic Field Instrument Suite and Integrated Science (EMFISIS, Kletzing et al., 2013) to analyze the magnetospheric wave properties over 10 Hz to 400 kHz. Figures 2c–2f show the enhancement of magnetosonic (MS) waves ($\sim 37 \text{ pT}$) with frequency below the lower hybrid resonance frequency (f_{LH}) identified based on their linear polarization (ellipticity ~ 0 , Figure 2e) and oblique propagation (Figure 2f). Moreover, very weak hiss emissions ($\sim 8 \text{ pT}$) were observed inside the plasmasphere. EMIC waves (from the triaxial fluxgate magnetometer) were observed simultaneously (Figure 2g), both in the helium and hydrogen band over ~02:30–02:55 UT. RBSP-B, traveling at lower L shells (~ 3.0 – 5.2), did not observe EMIC waves. Proton flux (Figure 2h, from the Helium, Oxygen, Proton, Electron Mass Spectrometer, HOPE in the Energetic Particle, Composition, and Thermal Plasma Suite, ECT; Funsten et al., 2013; Spence et al., 2013) showed injection of ~ 10 s keV protons from the magnetotail, associated with enhanced EMIC wave activity. The L shell extent of EMIC waves observed by RBSP-A was $L \sim 5.4$ – 5.7 for T89 ($L \sim 5.7$ – 5.9 for T05), projected along the FIREBIRD trajectories as a gray (blue) vertical colored block in Panels 1iii. The strongest electron precipitation observed by FIREBIRD was consistent with the L shell extent and location of the EMIC waves observed by RBSP-A, although the two spacecraft were ~ 1.6 MLT apart. In addition, given that the

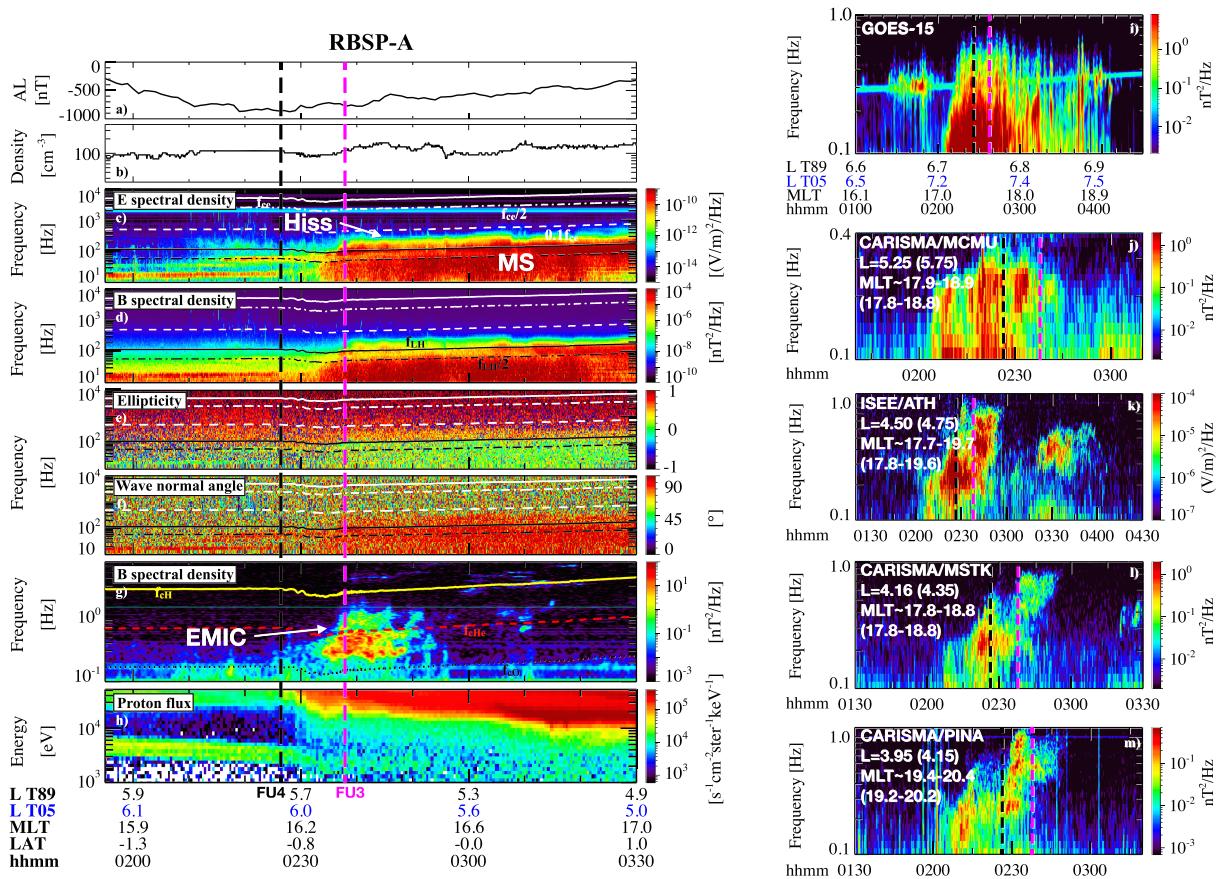


Figure 2. Overview of wave observations. From RBSP-A: (a) AL index, (b) total electron density, (c and d) frequency-time spectrogram of electric and magnetic spectral density, (e) ellipticity, (f) wave normal angle, (g) spectrogram of magnetic spectral density from MAG, and (h) proton flux from HOPE. (i) GOES-15 EMIC wave observations. EMIC wave observations at ground-based stations: (j) CARISMA/MCMU, (k) ISEE/ATH, (l) CARISMA/MSTK, and (m) CARISMA/PINA, with indicated L and MLT values from IGRF and T89 (in parentheses). The vertical dashed lines indicate the UT of magnetic conjunction with FU4 (black) and FU3 (magenta). L T89 (black) and L T05 (blue) on the x axes indicate the L shell values from the T89 and T05 magnetic field model, respectively.

electron flux outside this L shell extent was still larger than the background level, especially at lower L shells, precipitation may have occurred even beyond this region, probably due to a larger extent of EMIC waves than that observed by RBSP-A alone. In fact, GOES-15 (part of the Geostationary Operational Environmental Satellite constellation) also recorded EMIC wave activity near dusk over L shells of ~ 6.5 – 7.5 (Figure 2i) before 2 UT and after ~ 3 UT (the waves observed over 2–3 UT are probably relevant to other ultralow frequency waves (ULF), such as Pi2). Additionally, EMIC wave activity was recorded at ground, as shown in Figures 2j–2m by the stations MCMU, MSTK, and PINA in the Canadian Array for Realtime Investigations of Magnetic Activity (CARISMA, Mann et al., 2008) and ATH magnetometer, part of the Institute for Space-Earth Environmental Research (ISEE). The extent of these observations is likely relevant to ionospheric ducting (Kim et al., 2010). The collection of such ground magnetometer observations ($L \sim 4.1$ – 5.3 , $MLT \sim 17.7$ – 20.4), together with the observations from RBSP-A and GOES-15, showed that EMIC wave activity occurred over a rather broad region, although these observations are not necessarily relevant to the same EMIC waves stretching in L and MLT.

As mentioned in section 1, EMIC waves typically scatter not only electrons but also ~ 10 s keV protons. Since FIREBIRD only provides electron measurements, we take advantage of the NOAA/POES constellation to investigate the evidence of proton precipitation. NOAA-18 and NOAA-19 (Figure S1 in the supporting information) were traveling in a similar L-MLT region to RBSP-A and FIREBIRD (Figure 1c), and both observed proton precipitation at energies over 10s–100s keV (more details in the Supporting Information), roughly consistent with the L shell extent and location of EMIC waves observed by RBSP-A. These observations, together with ground-based measurements, provide additional evidence

that EMIC waves were active longitudinally and driving scattering at MLTs between RBSP-A and FIREBIRD locations.

3. Evolution of Electron PSD

Electron phase space density (PSD) evolution provides key information on the nonadiabatic radiation belt electron dynamics to evaluate whether the EMIC waves played a significant role in electron loss in this event. Figure 3 shows the PSD as a function of L^* (T05, Tsyganenko & Sitnov, 2005) for fixed values of the first (μ) and second (K) adiabatic invariants, calculated using the method by Chen et al. (2006) over the interval including RBSP-A EMIC wave observations (00–05 UT). PSD showed decreases from the first pass (black to dark blue lines) to the second one (cyan to green) within a narrow region of $L^* \sim 4.8\text{--}5.1$ (highlighted by gray blocks), with more evident PSD dips for increasing values of μ and K . For example, at low μ and K in Figure 3a (~60° pitch angle and energy ~300 keV, at $L^* \sim 4.8\text{--}5.1$), even though the electron PSD dropout was present, it was not as clear as that at higher μ in Figure 3g, where the electron energy is ~1 MeV. Similarly, the PSD dip in Figure 3a is less evident than that in Figure 3c where higher values of K correspond to lower pitch angles (~30°). These PSD dips that occurred within such a spatially narrow region and deepened with increasing values of μ and K are neither the signature of magnetopause shadowing nor that of outward radial diffusion loss, instead are consistent with local pitch angle scattering losses caused by EMIC waves (Shprits et al., 2017; Su et al., 2016; Xiang et al., 2017). EMIC waves are known to be more efficient in scattering electrons at higher energies and low-intermediate pitch angles (Jordanova et al., 2008; Meredith et al., 2003; Usanova et al., 2014), which accounts for the PSD evolution trends as those in Figure 3. Reasonably, the region where these PSD minima occurred for this event overall agrees with the location, spatial extent ($L^* \sim 4.7\text{--}4.9$, and UT (~2–3 UT) of EMIC waves observed by RBSP-A (and also electron precipitation observed by FIREBIRD), suggesting that these waves are important in this event for depleting the outer radiation belt electrons through scattering them into the loss cone. Since the dips were already visible at relatively low values of μ (~200 MeV/G), such a loss was effective at subrelativistic energies as well.

4. Quantifying Electron Precipitation Using Quasi-Linear Theory

The location of wave observations from RBSP-A and at ground, the associated electron precipitation, as well as the results of the PSD analysis, suggest that EMIC waves potentially drove such electron precipitation. Among all available satellites, RBSP-A was in the closest conjunction with FIREBIRD-II, despite the ~1.6 MLT difference. Therefore, we adopt the plasma and wave observations from RBSP-A (averaged over 02:37–02:42 UT around the FU3 conjunction) to determine which waves (among plasmaspheric hiss, MS waves, and EMIC waves) were primarily driving the electron precipitation, assuming that the wave and plasma conditions (near the equator) at the perfectly conjugate location to FIREBIRD are similar to those at the location of RBSP-A. The ratio of plasma frequency and electron gyrofrequency (f_{pe}/f_{ce}) is ~22.9. The EMIC wave spectrum used in the simulation has lower cutoffs of 0.26 f_{ci} (i: H and He), and upper cutoffs of 0.99 f_{cHe} for He-band and 0.65 f_{cH} for H-band, by only considering the wave power well above the background noise level ($>10^{-2}$ nT²/Hz). EMIC waves are assumed to be field-aligned (consistent with the observations by RBSP-A), which leads to similar results by using the latitudinally varying wave normal distribution (Ni et al., 2015). We estimated the cold ion composition as 84%:3%:13% (H⁺:He⁺:O⁺) by calculating the ion density from the two lowest HOPE channels following the method used in Blum et al. (2019) and Gamayunov et al. (2018).

Figures 4a–4c show the bounce-averaged pitch angle diffusion coefficients for EMIC waves at $L \sim 5.9$, calculated from cold plasma quasi-linear theory (Ni et al., 2015; Summers, 2005). The diffusion coefficients of hiss and MS waves are shown in Figure S3 in the supporting information. The diffusion coefficients for EMIC waves are higher than hiss and MS waves at energies above several hundred keV, indicating that EMIC waves dominated the pitch angle scattering and primarily drove the observed electron precipitation. While He-band EMIC waves efficiently scatter electrons at $>\sim 1$ MeV at rates $>10^{-3}$ s⁻¹ (Figure 4a), H-band ones scatter sub-MeV electrons at $>\sim 800$ keV (Figure 4b).

Cold plasma theory may become invalid at frequencies close to f_{cHe} , where the He-band EMIC waves are strong (Figure 2g). Therefore, we also adopt a wavenumber analysis (Chen et al., 2019), which contains the physics of hot plasma, based only on the electromagnetic wave field data to extract the

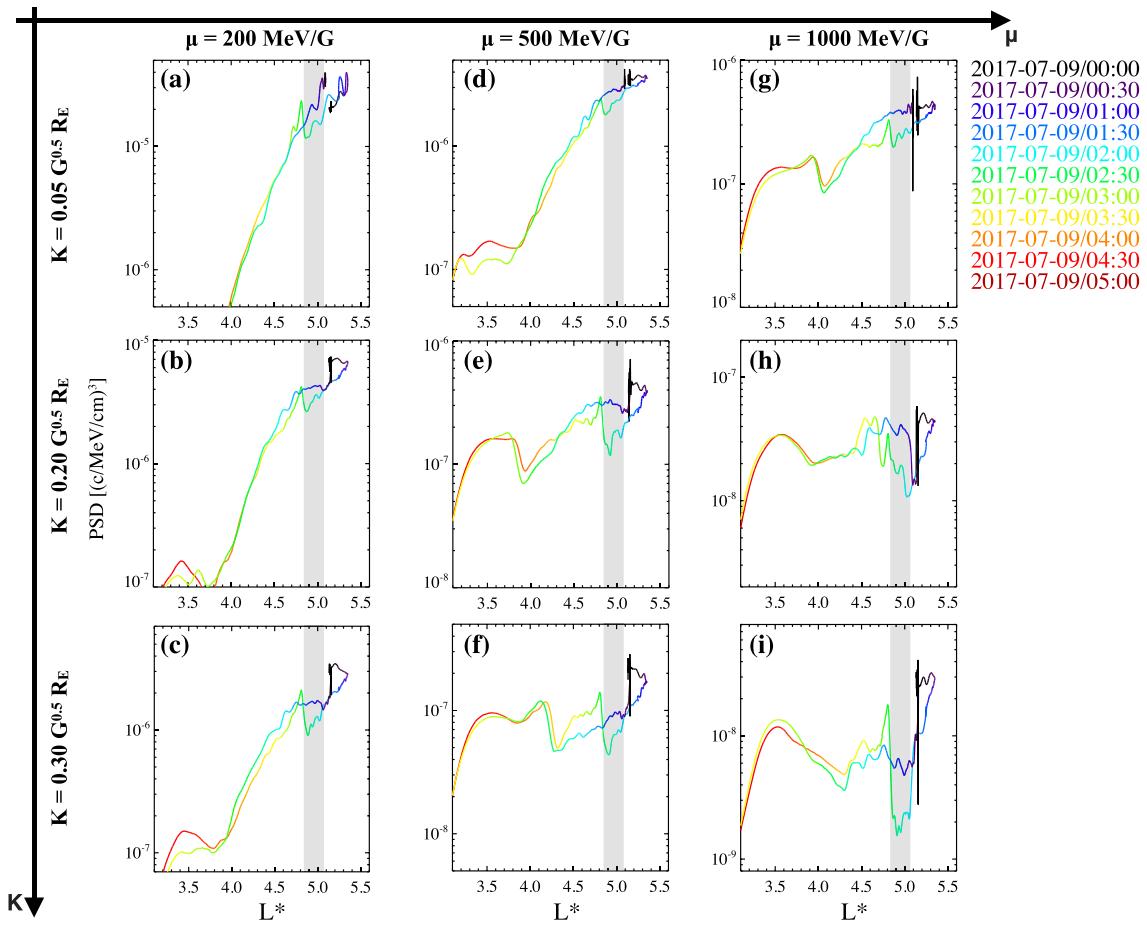


Figure 3. Evolution of electron phase space density (PSD) as a function of L^* (T05). Columns (rows) show the PSD at constant values of the first (second) adiabatic invariant μ (K) with varying values of the second (first) adiabatic invariant K (μ). PSD is color coded every 30 min. Vertical gray blocks highlight the PSD dips.

wavenumber spectrum, where background plasma approximations or dispersion relation assumptions are not required. Two of the Maxwell's equations for the components of the electromagnetic fields are solved in the frequency domain (for details, see Chen et al., 2019) to obtain the median wavenumber spectrum (Figure 4g, black). After performing a power law fit to the obtained wave dispersion (Figure 4g, gray) and assuming constant wavenumber spectrum along the field line, it is possible to calculate the bounce-averaged pitch angle diffusion coefficients (Figures 4d–4f) and the minimum resonant energy (Figure 4h) for the observed parallel EMIC waves. As expected from hot plasma effects, He-band E_{\min} increases (to ~ 4 MeV) compared to that predicted by the cold plasma calculation. Although the wavenumber spectrum shows high refractive index (kc/ω) in the H-band, lowering the hot plasma E_{\min} down to ~ 700 keV, H-band EMIC waves can scatter electrons efficiently ($>10^{-3} \text{ s}^{-1}$) only at $>\sim 1$ MeV.

In order to evaluate the scattering efficiency of each wave type, we compare the diffusion coefficients at the bounce loss cone to the strong diffusion limit (Figure 4i). We define electrons subject to “efficient pitch angle scattering” if their pitch angle diffusion coefficient is larger than $0.1 D_{SD}$, where D_{SD} is the strong diffusion limit. With this definition, the minimum energy of electrons subject to efficient pitch angle scattering due to EMIC waves is ~ 800 keV (~ 1 MeV) for the results without (with) hot plasma effects included. Pitch angle scattering due to hiss and MS waves is well below the strong diffusion limit. All above findings indicate that H-band (He-band) EMIC waves play a dominant role in scattering electrons at energies $\sim 0.8\text{--}4$ MeV (>4 MeV) into the loss cone, suggesting that other mechanisms are required to precipitate electrons at $<\sim 800$ keV.

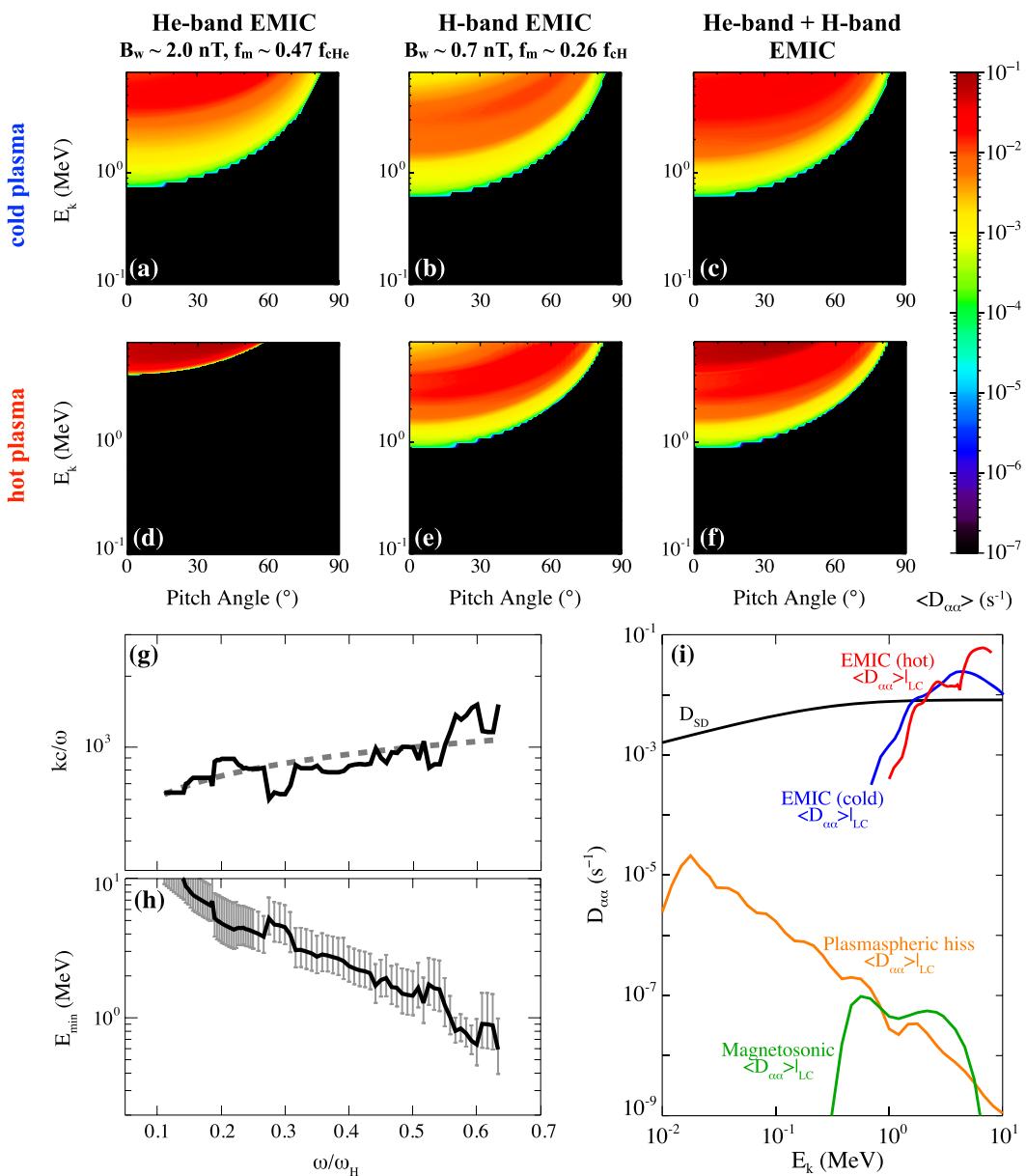


Figure 4. Bounce-averaged pitch angle diffusion coefficients without (a–c) and with (d–f) including hot plasma effects: He-band (a, d), H-band (b, e), and total EMIC waves (c, f). B_w is the root-mean-square wave amplitude during the EMIC wave interval 02:37–02:42 UT and f_m is the frequency corresponding to the highest EMIC wave intensity. (g) Median (black) and fit (dashed gray) of the wave refractive index from the wavenumber analysis and (h) median E_{\min} with the first and third quartiles (gray vertical bars) versus normalized frequency. (i) Comparison between the strong diffusion limit (black) and the diffusion coefficients at the bounce loss cone color coded by each wave mode.

5. Summary and Discussion

In this letter, we report a direct observation of subrelativistic electron precipitation from FIREBIRD-II, potentially driven by EMIC waves. The principal findings of this study are summarized as follows.

1. Both FIREBIRD-II units observed electron precipitation from subrelativistic energies of ~250 keV up to ~1 MeV, strongest within L (T89) ~5.4–5.7 at ~17.9 MLT lasting at least ~11 min.
2. Near the time interval and location of precipitation, multipoint measurements from satellites and ground-based magnetometers recorded strong EMIC wave activity and EMIC-driven precipitation signatures. Among them, RBSP-A (the closest to FIREBIRD-II, albeit with $\Delta MLT \sim 1.6$) observed intense He-band waves ($B_w \sim 2.0$ nT) and modest H-band waves ($B_w \sim 0.7$ nT). NOAA-18 and NOAA-19, crossing

similar L shells at MLTs between FIREBIRD and RBSP-A, observed \sim 10s–100s keV proton precipitation as well, a typical signature of EMIC-driven precipitation.

3. Simultaneously with EMIC waves, plasmaspheric hiss and MS waves were also observed by RBSP-A, although much weaker (8 and 37 pT, respectively). The calculation of the quasi-linear pitch angle diffusion coefficients indicates that while hiss and MS waves did not scatter electrons efficiently, H-band (He-band) EMIC waves played a major role in driving the electron precipitation at energies \sim 0.8–4 MeV (>4 MeV).
4. The EMIC waves observed by RBSP-A over \sim 02:30–03:00 UT were also associated with electron flux dropouts in the outer radiation belt. The PSD profiles revealed by RBSP indeed show localized dips coincident with EMIC waves, which become more evident as μ and K increase. Interestingly, the PSD dips were evident at relatively low values of μ , indicating that EMIC waves played a role in precipitating radiation belt electrons even at a few 100s keV. The scale size and location of PSD dips agree well with the L shell extent of EMIC waves and electron precipitation measured by FIREBIRD-II.

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It is important to note that the simulation results without (with) including hot plasma effects due to EMIC waves can only explain the electron precipitation at energies $>\sim$ 800 keV ($>\sim$ 1 MeV), requiring other mechanisms to explain the low energy electron precipitation down to \sim 250 keV. For example, nonresonant scattering (e.g., Chen et al., 2016) and nonlinear effects (Denton et al., 2019; Hendry et al., 2019; Kubota et al., 2015; Omura & Zhao, 2012, 2013) are suggested to lower the E_{\min} down to a few 100s keV. However, the evaluation of these effects is beyond the scope of this case study and left as future investigation.

The results of our study are consistent with previous observations of sub-MeV electron precipitation (e.g., Capannolo et al., 2019; Clilverd et al., 2015; Hendry et al., 2017). However, the differential energy channels extending down to \sim 250 keV, available on FIREBIRD-II, allowed us to directly and unambiguously detect subrelativistic electron precipitation without further assumptions or modeling, and without proton contamination, which has been limiting reliable observations from POES at subrelativistic energies. Owing to a combination of multiple satellites and ground-based stations, our study provides the direct observation of subrelativistic (at least down to \sim 250 keV) electron precipitation due to efficient pitch angle scattering potentially driven by the observed EMIC waves as summarized in (1)–(4), although quasi-linear theory is insufficient to explain the full precipitation energy range and no wave measurements were available at the exact conjugate location to FIREBIRD. Nevertheless, this case study highlights FIREBIRD-II capabilities to provide electron measurements with high resolution in time and energy not only for microbursts but also for other precipitation patterns, such as EMIC-driven electron precipitation, and also potentially provide a valuable interpretation of electron precipitation data from POES and other LEO satellites.

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Erratum

Due to a production error, an outdated version of Figure 2 appeared in the originally published version of this article. This error has been corrected, and the present version may be considered the official version of record.