

Ice Engineering Guide for 📱 Design and **Construction of Small Craft Harbors**

C. Allen Wortley

ICE ENGINEERING GUIDE

FOR

DESIGN AND CONSTRUCTION

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SMALL CRAFT HARBORS

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ABSTRACT

This ice engineering guide is an aid to designers and builders of small craft harbors in northern climates. In these areas ice does great damage to structures. The guide has been prepared from a review of the technical literature, consultation with knowledgeable engineers, contractors, and scientists, and field observations. It is applicable to lake ice in protected harbors. Information on the behavior and properties of ice and ice covers is presented. Descriptions of ice pressures of thermal origin and estimated thrusts on structures are given. The operation and design of compressed air suppression systems are explained. Safe bearing loads on ice sheets are estimated, as well as ice uplift loads for design. The information presented answers many questions about ice behavior and serves as a basis for innovative and economical small craft harbor design.

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PREFACE

More than a hundred years ago Dumble (1858) stated,

"I may add, that the ignorance, or want of proper appreciation, of the properties of ice, evinced in the construction of numerous wharves, piers, and bridges on the inland lakes and rivers of Canada and the northern States, has proved a source of infinite annoyance and of immense expense".

Dumble was an engineer with the Cobourg and Peterborough Railway. Today, most designers of small craft harbors in northern climates have an appreciation of the behavior of ice. However, the level of relative ignorance about ice behavior is still high.

The successful design and construction of small craft harbor structures (such as docks, piers, and piles) subjected to ice forces and motions continues to be an uncertain science. The four methods currently being used are: (1) Predict the ice forces and motions, both horizontal and vertical, and design the marina structures to withstand them, (2) Predict the ice prevention or suppression capabilities of various systems to remove the adverse effects of ice or render them ineffective, (3) Estimate the relative scope of structure required by local comparative studies of successful and unsuccessful constructions and build accordingly, and (4) Remove the dockage by either beaching floating docks or retracting fixed piers.

Designers and builders should continue innovations and improvements in the state of the art of small craft harbor construction in northern climates. This publication endeavors to aid them and serve as a guide. It is a first step. It is incomplete. And likely will never be totally complete due to the uncertainty existing in the field of ice engineering. It is hoped the U.S. Cold Regions Research and Engineering Laboratory (US CRREL) will subsequently expand and update ice information necessary to construct small craft harbors in ice conditions.

Madison, Wisconsin January, 1978 C. Allen Wortley Principal Investigator

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The project has been undertaken with the assistance of an Advisory Committee. Members were

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The principal investigator is especially grateful for the suggestions and assistance received from Messrs. G. D. Ashton, G. E. Frankenstein, and D. E. Nevel, all of the US CRREL.

INTRODUCTION

This publication is a guide for designers and builders of small craft harbors in northern climates, where ice forces and motions do great damage to man's works.

It has been prepared from a review of the literature, consultation with knowledgeable engineers, contractors and scientists, and the experiences and observations of the principal investigator and members of an Advisory Committee.

Zarling (1974) prepared an annotated bibliography entitled "Ice Engineering in Small Craft Marinas". This work served as the starting point for reviewing the literature. Additional and more recent publications were examined as well as many of those noted in the bibliography. There is a large amount of information in the overall field of ice engineering and ice physics. However, very little is easily applied by the small craft harbor designer or builder because of its research-technical nature. Also much of the information doesn't relate directly to the problems of designers and builders, or is of such orders of magnitude as to render it useless for safe and economical design.

However it is from the more useful, more current, and more applied information in the literature that much of this guide has been prepared. The literature has been freely cited. No additional laboratory experimental work has been undertaken. Field observations and measurements have been made and have influenced the recommendations derived from the information in the literature.

Ice for purposes of this guide is primarily stationary lake ice. River ice, ice floes and sea ice are not specifically dealt with. They can and do present different problems. Small craft harbors are customarily built in sheltered areas, away from dynamic ice forces.

The guide is divided into five parts. Part I presents information on the characteristics, behavior and properties of ice and ice covers. It presents as much specific quantitative information as is available. The other information presented is of a general nature but useful in understanding ice and its behavior.

Part II deals with ice pressures of thermal origin. It is derived largely from current laboratory research on thrusts on hydraulic structures. Certain information however is applicable to small craft harbors.

Part III presents recommendations for design and construction of compressed air ice suppression systems and other information on weakening ice. These suppression systems are effective in small craft harbors.

Part IV discusses the bearing capacity of ice and presents recommended maximum short term construction loadings.

Part V gives minimum and expected maximum values for ice uplift forces. Information about other ice forces is also presented.

As stated initially this is a guide for designers and builders of marinas in northern climates. It by no means answers all the pertinent questions, nor does it present all the "how to's". This is beyond our present scope and our present abilities.

It is hoped that the reader will review all Parts of this guide and form an overall understanding of ice behavior. Successful, innovative and economical design of small craft harbor structures is our goal.

UNITS

Conversion Factors: U.S. Customary to SI Metric Units

Measurements in this guide are shown in customary U.S. units with SI metric equivalents. Specific quantities are converted accurately, but where the measurements are general the conversion is approximate only. As much of the source data was in the metric system and not the U.S. customary or the SI systems, some data appear to have "odd" values. For example, pressures in kilograms per square centimeter (kg/cm²) have been converted to pounds per square inch (psi) and kilonewtons per square meter (kN/m²). The SI unit for pressure is pascal (Pa). One Pa is equal to one N/m². Because we are dealing with forces, and forces per unit length, as well as pressures, we have elected to leave pressures in force-length units rather than the derived pressure unit Pa.

Some U.S. Customary to SI Metric conversion factors are given below. The ASTM E380-76 <u>Standard for Metric Practice</u> gives additional conversion factors.

To Convert	To	Multiply By
<pre>inches (in) inches (in) feet (ft) yards (yd) miles (mi) inches per yard (in/yd)</pre>	millimeters (mm) centimeters (cm) meters (m) meters (m) kilometers (km) millimeters per meter (mm/m)	25.40 2.540 0.3048 0.9144 1.609 27.78
square inches (sq in) square feet (sq ft)	square centimeters (cm^2) square meters (m^2)	6.45 0.093
cubic inches (cu in) cubic feet (cu ft)	cubic centimeters (cm ³) cubic meters (m ³)	16.39 0.0283
pounds (1b) tons (ton) tons (ton)	kilograms (kg) kilograms (kg) metric tons (mt)	0.453 907.185 0.907

To Convert	<u>To</u>	Multiply By
one pound force (lb) one kip (kip) one kilogram force (kg)	newtons (N) kilonewtons (kN) newtons (N)	4.448 4.448 9.806
kips per foot (kips/ft)	kilonewtons per meter (kN/m)	14.59
pounds per square foot (psf) pounds per square inch (psi) pounds per square inch (psi) pounds per square inch (psi)	newtons per square meter (N/m ²) kilonewtons per square meter (kN/m ²) meganewtons per square meter (MN/m ²) kilograms per square centimeter (kq/cm ²)*	47.88 6.89 0.00689 0.07031
kips per square inch (ksi)	meganewtons per square meter (MN/m ²)	6.89
pounds per cubic foot (pef)	kilograms per cubic meter (kg/m ³)*	16.02
Degree Fahrenheit (F) Temperature Fahrenheit (F)	Degree Celsius (C)* Temperature Celsius (C)*	0.555 (F-32)/1.8
square feet per minute (ft²/min)	square meters per second (m ² /sec)	0.001548
square feet per hour (ft ² /hr)	square centimeters per second (cm ² /hr)	929.03
Btu per hour per square foot (Btu hr ^{-l} ft ⁻²)	watts per square meter (W m ⁻²)	3.152
Btu per pound per degree Fahrenheit (Btu]b ⁻¹ F ⁻¹)	joule per kilogram per degree Celsius (J kg ^{-l} C ^{-l})*	4186.8
Btu per hour per square foot per degree Fahrenheit (Btu hr ⁻¹ ft ⁻² F ⁻¹)	watts per square meter per degree Celsius (W m $^{-2}$ C $^{-1}$)	5.678
Stu inch per hour per square foot per degree Fahrenheit (Btu in hr ⁻¹ ft ⁻² F ⁻¹)	watts per meter per degree Celsius (W m $^{-1}$ C $^{-1}$)*	0.1442
miles per hour (mph)	meters per second	0.4470

^{*}not an SI unit

PART T

TOF AND ICE COVERS

This section of the guide contains information on the crystal structure of ice, the polycrystalline forms of ice, the nucleation and growth of ice, the freeze-up and break-up processes on fresh-water lakes, the thicknesses and behavior of ice covers and some introductory material on the deformation and strength of polycrystalline ice, a visco-elastic material.

Some of the material is pertinent to present design techniques for small craft harbors and some is only of general background interest. The general information may be useful in developing new design techniques, explaining behavior of ice covers, understanding other ice research papers and monographs that readers may come in contact with, and in other ways in the future. The material should answer some of the readers' questions about ice and confirm some of their observations in the field.

Simply stated, ice is not well understood. At the present time straight-forward and proven design approaches for small craft harbors are not available. Therefore the best approach for design is to begin with an understanding of the present state of knowledge about ice, especially in the context of small craft harbors, and go from there.

Those wishing more detailed technical and theoretical information on the physics and mechanics of ice should consult references such as the two monographs by Glen (1974, 1975) and textbooks by Pounder (1965) and Hobbs (1974).

Crystal Structure of Ice

Ice is a solid consisting of a crystallographic arrangement of water molecules. Each molecule has an oxygen atom with chemical bonds with two hydrogen atoms.

Although several forms of ice can exist, we shall be concerned only with the ordinary form of ice. It is the form with which most people are familiar. This form is hexagonal in crystal symmetry. The symmetry is clearly demonstrated in the symmetry of many snow crystals.

Each crystal of ice is made up of water molecules arranged in a regular geometric pattern forming hexagon prisms. The axis of the prism is known as the optic axis and the surface perpendicular to this axis is called the basal plane. Crystallographic terminology labels planes and directions. In the hexagonal system, there are three directions in the basal plane that are crystallographically equivalent. These directions are called a, b and d (sometimes al alequation) and the axis perpendicular to the basal plane is called the c-axis (i.e. the optic axis).

Runnels (1966) suggests that at the molecular level the seemingly rigid perfection of the ice crystal is disrupted by an astonishingly busy traffic of molecules and migrating lattice faults. There is a good deal of empty space between molecules and hence ice has a rather open structure. Because of this fact, water, in contrast to almost all other substances, is less dense in the solid state than in the liquid state. This results in ice floating on water. The density of pure ice is given by Pounder (1965) as $57.2~\mathrm{pcf}$ (916.8 kg/m³).

Ice encountered in nature, such as on a lake, is in a high temperature state. It is floating on its own melt with its bottom surface at a temperature equal to its melting point and the upper surface varying with ambient conditions. Predicting the properties and behavior of a material in such a state is quite difficult.

Polycrystalline Ice

Most bulk ice in nature consists of polycrystalline ice, that is, ice composed of a large number of single crystals in different orientations. Polycrystalline ice behaves differently from single crystals of ice because of the existence of grain boundaries and the fact that single crystals are anisotropic. A single crystal of ice has physical properties that are different in directions parallel and perpendicular to its c-axis. If the constituent crystals of polycrystalline ice are randomly oriented, then the bulk properties will be some kind of average of the properties of the single crystals of which it is composed. For some properties such as elasticity, this averaging is fairly straightforward, but for others, such as plastic deformation, there is such a radical difference in the behavior of ice crystals in different directions that the physical properties of the polycrystal are quite different from that of the single crystal.

Moreover, the grain size as well as the crystal orientation of polycrystalline ice has an effect on the strength and behavior of an ice sheet. Fortunately, from a design point of view, those effects may be minor for cold polycrystalline ice, the type that normally would present the most severe forces on structures placed in ice. This may not always be true for all kinds of polycrystalline ice encountered in lakes and rivers. The kind of ice which forms depends upon its genesis.

Nucleation and Growth of Ice Crystals

Ice in the absence of very great supercoolings or supersaturations will only form if some nucleus is present on which it can grow. This can be ice itself, and if an ice seed is present, ice will normally develop and grow. It may also be a nucleus consisting of some other material.

Once an ice nucleus is present, ice can grow from the surrounding water provided the temperature remains favorable. How ice will grow will depend on whether the water surface is supercooled, and also on the amount of impurities it contains.

Michel and Ramseier (1969) have **classified** river and lake ice based upon its genesis, structure and texture. An ice sheet or ice cover is divided into three basic ice formations: primary ice, secondary ice, and snow ice.

Primary ice is the first type of ice of uniform structure and texture which forms on a water body. On a calm water body the primary ice is in the form of an ice skim which grows horizontally in the supercooled layer and is about a hundredth of an inch (a few tenths of a mm) thick. Usually when water freezes a little plate-like disc forms. It may rapidly develop dendritic extensions depending on the amount of supercooling and the rate of heat loss. Since the planar structure will naturally float with its plane parallel to the surface, the ice which forms when water freezes often has its surface crystals oriented with their c-axes vertical, i.e. perpendicular to the free surface. These initial crystals spread over the surface until they interact and form a layer over the whole surface of the water. The subsequent growth of this layer downwards into the water is not usually dendritic.

In rough and turbulent water the primary ice consists of congealed frazil slush which can be an inch (several centimeters) thick. Frazil slush is a floating agglomerate of loosely packed individual ice crystals having the form of small discoids or spicules which are formed in supercooled turbulent water. If the nucleation occurred by snow, the resulting congealed snow slush (i.e. a loosely packed agglomerate of floating snow particles) would also be part of the primary ice.

Secondary ice forms parallel to the heat flow which is perpendicular to the primary ice. As the ice grows down into the water under conditions of calm formation and growth the grain boundaries are almost perpendicular to the surface. The result is a columnar-grained structure with the c-axes

vertical in the long direction of the grains. This type of secondary ice is referred to as type S1. The crystal size increases with depth, because unfavorably oriented crystals disappear. S1 ice is found in lakes, reservoirs and in rivers with low flow velocities.

S2 columnar ice is a secondary ice form with the c-axes horizontal. It forms under similar conditions to S1 ice but where the primary ice has a random or preferred vertical orientation superimposed on a random orientation. As the ice grows down into the water, the c-axes vertical preferred orientation gives way to one in which the c-axes are horizontal. Some crystals are wedged out at the expense of others. This type of ice is found in lakes, reservoirs, shores and rivers.

Another type of secondary ice, but graunular rather than columnar in form, is S4 Congealed Frazil Slush. Its crystal boundaries are irregular and the grains highly angular and randomly oriented. Frazil is found in rivers where it is taken downstream and deposited upwards under the secondary ice overlying more slowly moving water. S4 ice is found in rivers and reservoirs, and lakes fed by turbulent waters. An entire ice cover can consist of S4 ice or it can be found layered between columnar forms of ice.

An interesting ice form, anchor ice, results when frazil is being formed in supercooled water. The frazil particles stick and grow on bottom weeds and stones. After a night of frazil formation the morning sun can cause anchor ice on stones and other objects to float to the surface and be carried downstream.

Snow ice is the third basic formation of ice. This type of ice always forms on top of the primary ice due to snow deposits which lie on the ice cover. Snow ice may form due to variations in discharge of water, by melt or rain, or by the depression of the ice cover due to a heavy snow load. In general snow ice is formed from any imaginable kind of water source. It is the designation for Snow Ice and its grains are round to angular and randomly oriented. Being granular it is similar in appearance to the S4 Congealed Frazil Slush ice found usually under the secondary ice layer.

The crystal size of ice depends greatly on meteorological parameters. If the water surface is calm and the temperature gradient is not too steep extra large crystals will develop. Their diameters can be greater than 0.8 inches (20 mm). If the temperature gradient becomes larger then the crystal size will decrease. If snow causes the nucleation, crystals will be of the 0.2 inches (5 mm) size and smaller. Extra large and giant size crystals can form under extremely calm conditions such as in a pool during the formation of the primary ice layer. Crystal dimensions in yards (meters) are possible.

During the evolution of an ice cover the crystal size will change. In lake ice it is possible to go from fine sized to extra large sized crystals in a columnar layer if nucleation has occurred by snow.

Gow and Langston (1977) describe the growth history of lake ice in relation to its stratigraphic, crystalline and mechanical structure. They have

successfully correlated specific increments of ice growth with major weather events. This was accomplished by diagnosing bubble layering observed in strationaphic samples of lake ice. The ice covers of the two New Hampshire lakes used in their studies contained two major components: lake ice formed solely from the freezing of lake water and snow ice formed from the freezing of water-saturated snow. The lake ice component consisted of predominately fibrous-textured crystals with a vertical c-axis orientation. Columnar crystals elongated in the vertical but with horizontal c-axes fabrics. were observed only in the top 4 to 6 inches (10 to 15 cm) of ice that fringed the shores of both lakes. These fabrics had reverted entirely to a vertical c-axis condition at 6 inches (15 cm) depth. The reason for this reversal in orientation texture has not clearly been established. However, indications are that the temperature gradient in the water may be a significant factor in initially controlling the growth textures, since continued growth of ice near shore resulted in rapid elimination of the horizontal c-axis in favor of crystals with vertical c-axes.

Several variants of snow ice were distinguished, depending upon the source of infiltrating water. Both ice types were readily distinguished on the basis of differences in their stratigraphic and crystalline structure.

Michel and Ramseier (1969) have also found that studying ice profiles from lakes and rivers enables one to determine what some of the meteorological conditions were upon formation of the primary ice. Similar analyses can be obtained from snow ice and overlying snow. They believe that in the future it may even be possible to predict directly the texture at a particular site from meteorological and hydrodynamic data.

The predominant inclusion in ice is air. During solidification the air is rejected at the ice-water interface and is supplemented by gases arising from biochemical processes which take place in a natural body of water. Air can also be incorporated from the atmosphere through cracks, drained snow ice or an agitated water surface. If the growth is slow most of the air will be rejected and the ice will be transparent like glass. This transparent ice is sometimes called black ice. (The term blue ice is sometimes used in connection with glacial ice, but not transparent lake ice.)

If the solidification process proceeds rapidly air in the form of bubbles will be distributed throughout some parts of the ice cover. If air is being emitted continuously, the bubbles take on a cylindrical shape. It is the air entrapped in ice that causes it to appear white. This type of ice is sometimes called white ice. Snow ice is an example of white ice because of entrapped air. Grain size does not determine the color of ice. Frozen frazil ice can appear transparent in spite of its granular texture.

Freeze-Up and Break-Up of Fresh Water Lakes

Water, in the temperature range of 39°F to 32°F (4°C to 0°C), increases in volume with decreasing temperature. This is contrary to the behavior of most substances. In this temperature range the coefficient of expansion of water is negative, that is, water expands as it becomes cooler. Since the volume of a given mass of water is smallest at 39°F (4°C), it has its maximum density at this temperature. The negative coefficient of expansion of water in this range is one of nature's anomalies. This explains why lakes first freeze at their surfaces and why somewhat warmer (and denser) water is found near the bottom.

Williams (1966) has prepared a detailed and comprehensive description of how fresh water lakes freeze and thaw. Michel (1971) also describes the winter regime of rivers and lakes in a monograph. It too is comprehensive and contains information on rivers as well as lakes. For our purposes we will recite the lake ice processes described by Williams (1966).

His descriptions are applicable to sheet ice, the type that forms on relatively still lakes, slow-moving rivers, and sheltered harbors. Williams designates three distinct periods in the life of sheet ice: the freeze-up period, the ice-growth period, and the ice-melting period. These periods roughly correspond with the fall, winter and spring seasons.

During the fall, or early winter, smaller lakes freeze over. A solid ice sheet forms. In larger lakes the sheet may only cover harbors and the area near the shoreline, but is generally extensive. The Great Lakes rarely completely freeze over, at least for any extended period of time.

The formation of an ice cover on a lake is a function of the water's heat exchange with the atmosphere, the initial amount of heat stored in the water body, and the amount of inflow of warm water (heat) to the site. The amount of heat lost to the atmosphere is a function of the air temperature, the wind velocity, and solar radiation. The amount of heat that can be stored in a water body is a function of depth. Usually, the deeper the lake, the deeper the convective mixing and the slower the rate of water cooling for a given surface heat loss. (In a man-made harbor, depths are usually shallow and hence significant thermal reserves because of depth are not present.) Inflow of heat by currents from warmer water in rivers or from deep reservoirs can prevent or delay sheet ice formation at certain sites.

The cooling of a fresh water lake occurs in two stages: gradual cooling until all the water is at a temperature of about $39^{\circ}F$ ($4^{\circ}C$), and cooling of the surface water from the time the water is isothermal (i.e. the same temperature, top to bottom) at $39^{\circ}F$ ($4^{\circ}C$) until sheet ice forms.

First the warm lake surface water cools down (even steams or mists when giving up its heat to the cool atmosphere) and thereby contracts and becomes more dense. This water sinks in the less dense lake and forces up the lower less dense, warmer water to in turn be cooled. This process is repeated until the lake is isothermal at 39° F (4° C). The lake has then "turned over".

From this point on the lighter, cooler surface water continues to cool down until it freezes as sheet ice. (A thin layer of ice first forms along the edges of the lake. A convective air motion from land toward rising warm air over the lake cools the lake edges more rapidly than the center.) Since the density of ice is even less—than that of freezing water, the ice floats on the water below it, and further freezing can only result from heat flow upward by conduction.

The mid-winter period, when the permanent ice cover on lakes gradually increases to a maximum thickness late in the winter, is a period of interest to harbor engineers concerned with ice pressures, ice uplifts, and bearing capacities.

When the ice cover is thin, and the air very cold, there will be a relatively large heat flow from the water, through the ice, to the atmosphere. This will cause a rapid increase in ice thickness at the interface between the ice and the water. As the ice cover thickens, an insulating effect occurs, and ice grows at a slower rate. A snow cover on the initial ice formation can effectively prevent any further ice growth.

Ice normally grows downward into the water, but not always. The upper surface of ice can be submerged below the water level due to the weight of a snow cover. Snow ice then forms and freezes on top of the ice sheet. Sometimes this causes rapid increases in the total thickness.

A frozen snow covered lake can be highly variable in its underlying stratigraphy. The snow cover can conceal slush and weak layers of ice, as well as strong sound ice.

The clearing of ice from bodies of water at break-up is affected by heat gain from the atmosphere, snow and ice conditions, wind and currents, and inflow or runoff of warm water to a site. The heat gained from the atmosphere weakens and melts the ice. The amount of solar radiation absorbed is determined by the reflectivity of the snow or ice surface. The properties of the ice cover determine the depth solar radiation will penetrate, causing internal melting. The thickness of the snow and ice is a measure of the amount of ice to be melted. The mechanical action of wind and currents is of great importance in breaking up the ice after it has been weakened by surface and internal melting.

At the beginning of break-up, snow melt runoff from surrounding land areas may flood onto part of the lake surface. At the same time the snow cover on the ice densifies and begins to melt. The runoff usually accumulates near the mouths of streams draining into the lakes and along the shoreline, producing dark-looking slush on the ice surface. At this stage the darkened ice cover surface absorbs a large proportion of the incoming solar radiation.

The ice around the shoreline tends to melt first, partly because of the darkened surface layers but also because the ice thickness is often less close to the shore, where it can only grow as thick as the depth of the water. Eventually the ice cover completely melts around the shoreline, leaving the main body of ice floating free. This body of ice can still have considerable strength.

During melting, drainage holes can develop in the ice sheet. In the early stages of melt these holes develop near the shore often where runoff flows into old thermal cracks. The flow is due to difference in level between the lake water supporting the ice and the melt water on the ice surface. Such drainage holes can enlarge rapidly, sometimes developing into holes one to two feet (0.3 to 0.6 m) in diameter.

Later on during melting, the main body of ice, floating free, melts at the surface. This melt water flows along the surface of the ice and eventually creates surface drainage patterns. It can drain to the open water adjacent to the shoreline, or down the holes that appear to develop preferentially along old thermal cracks. If the surface melt water can drain away, the surface of the ice becomes a porous, crumbly white structure which reflects solar radiation readily, thus retarding melting by solar radiation. As the melt season progresses, penetrating solar radiation causes internal melting in the ice sheet. At this time the ice sheet may consist of a shallow, porous surface layer several inches (several cm) thick, a layer of water-logged ice also several inches (several cm) thick, and then solid unmelted ice (whose lower surface is melting in the lake water).

In the final stages of break-up the underlying entrapped water layers result in darkened surface ice and most of the incoming solar radiation is absorbed. The ice is now ripe for breaking by wind and currents and is unsafe for over-ice transportation. The current created by strong winds will break up the ice cover and induce circulation that brings the warmer subsurface water to the surface. This can cause rapid melting of ice from a lake. Indeed the final disappearance of ice covers has occurred so quickly at times that early observers believed that the ice actually sank.

The term rotten ice is sometimes used to refer to disintegrating ice at break-up time. If the ice has a columnar structure it becomes candled and is referred to as candle ice. Melting is concentrated at the boundaries of the columnar prisms leaving a weak candle-like structure.

Williams (1966) and Michel (1971), as well as others, have proposed methods to estimate the date of freeze-up, maximum ice thickness, and the date of break-up of fresh water lake ice. The method of doing this is approached in two ways: develop formulas based on physical principles, and make an analysis of past records to give statistical limits within which maximum ice thickness or freeze-up or break-up can be expected to occur.

The operational use of formulas is usually limited because of difficulties in making dependable forcasts of such variables as air temperature and snow depth. It is also difficult to allow for natural variability, as the rate at which ice forms, grows, and melts varies not only from lake to lake but on an individual lake.

The statistical climatological approach also has limitations. The required long-term records are only available for some lakes and it is difficult to use the records collected to estimate freeze-up or break-up for other lakes with different sites and thermal regimes. In addition, even for lakes where some long-term records are available, the statistical

approach will only define upper and lower limits within which freeze-up or break-up will occur.

We believe direct on-site observations by the small craft harbor designer coupled with any information from local residents is the best approach to predicting freeze-up or break-up. Ice thickness determination can be handled in the same manner but of course modified by recommended design criteria for minimum thicknesses and strengths.

The U.S. Great Lakes Environmental Research Laboratories (GLERL) in cooperation with Canadian and other agencies are involved in efforts to understand and forcast behavior of ice formations on the Great Lakes. Samples and measurements of ice thicknesses from many places on the Great Lakes are underway. This information together with other published measurements and observations will be useful to marina designers.

Anecdotal Behavior of Natural Ice Covers

Many anecdotal accounts of the behavior of ice covers are contained in the literature. Some of the explanations about its behavior have since been proven wrong. Some have been substantiated with research and additional observations. During the nineteenth century Dumble (1858), who was an engineer with a railway company, published accounts of his observations on inland lakes and rivers in the northern states and Canada. A few of his observations are repeated below together with some of our parenthetical comments.

The most violent shoves of ice occur previous to rainstorms. (This would be associated with warm winter weather and thermal expansion.)

Contraction generally occurs at night, and is accompanied by sharp reports. (This would be associated with cooling nighttime weather and thermal contraction cracking.)

A coating of snow of any depth over six inches (15 cm) effectively prevents any motion of ice. (Snow is a very effective insulator and prevents thermal motion of ice.)

It is but reasonable to suppose that any solid, equally dense throughout its dimensions, and susceptible of expansion, would, when equally acted upon by the active agent or moving cause, expand from its center towards its circumference. This is the effect produced on any large field of ice of equal thickness and density, when acted upon uniformly by either the midday sun or warm winds. It is a fact, however, that it moves from other directions than from the center of the lake. Shoves are sometimes witnessed from the east and sometimes from the west, to the north and to the south.

Ice owing to the peculiar circumstances under which it sometimes forms, is not found to be equally pure or dense, neither is it of uniform thickness. This ice irregularly acted upon by warm winds, or by the slanting rays of the sun at different altitudes, shoves, or expands from various directions other than from the center of the lake.

Shoving is from the stronger and most susceptible ice toward the weaker and less expansive.

Ice on any large and irregular sheet of water studded with islands must naturally be of unequal thickness and density. There is therefore no doubt whatever, that the phenomenon of ice expanding and shoving from various directions is caused by the unequal thickness, density and glareness of the ice and likewise by the manner in which the heated atmosphere strikes it.

(Ice sheets and covers are variable in their composition and frequently expand, contract and move in erratic ways.)

It is observed, that when a large extent of field ice expands towards the shore it does not shove into deep bays but fractures from point to point in a zig-zag manner, across the chord at the mouth. A thrust of the main field must find less resistance across this chord than around the area of the bay. (Crack patterns can now easily be observed from the air and indicate where ice has failed as it moved into or by some shore configuration such as a bay.)

Much more recently Striegl (1952) described the formation of ice in the Great Lakes, and particularly in Lake Michigan.

Ice normally starts to form along the shores of the southern part of Lake Michigan about the middle of December. By the early part of January the initial ice sheet may reach a thickness of 6 to 8 inches (15 to 20 cm) and extend off shore distance of one to five miles (1.6 km to 8.0 km) depending on the continuity of cold weather and the strength of the winds. Frequently during this initial ice period the thin ice cover may be broken up by the wind and wave action. If the wind is toward shore, one of a series of windrows may develop where the drifting cake of ice will pile up on or under the sheet of ice until it may reach a depth of many feet (meters) above and below water, and may ultimately pack to the bottom where shallower water is reached. If the ice along the shore is weak the windrow may be pushed ashore until a large part of it is above the water level. If the wind is off shore the fields of floating ice are driven out into the lake where they may drift for long periods or pile up in ice windrows or along shore in another part of the lake, depending on the continuity and direction of the wind.

Engineering structures built in the lakes for harbor protection or other purposes must be designed to support enormous ice loads due to the ice build up on the structure by continued freezing of wave wash and spray or by loose ice piled on top of the structure by ice jams and then consolidated by continued freezing action. Those ice loads may pile up to 20 to 25 feet (6 to 8 m) above the top of the structures. Usually this load, does little damage to overloading individual structural members because of the extra support and bridging effect of the ice itself. However, it may cause considerable damage in other ways. Stones or concrete may be broken, through freezing and thawing action to such sizes that they will later be displaced by wave action. Also large breakwater stones, up to several tons each, may be carried bodily from the structure when encased in large ice

blocks. Such stones have been carried by ice floes miles from structures and dropped in navigable channels where they have become obstructions to navigation. Smaller rocks have been found in the channels of harbors which were picked up by the ice on the beach or from the structures and floated with the ice to areas where they were dropped on melting or breaking of the ice.

In Chicago in January 1948, a strong easterly wind storm, accompanied by sleet, snow and low temperatures, piled ice in a solid mass along the entire Illinois shoreline. The ice was as much as 15 feet (5 m) above the water surface and extended 50 to 100 feet (15 to 30 m) behind sea walls and bulkheads in Chicago, and from 100 to 300 feet (30 to 90 m) in front of the beaches and natural shoreline and other points.

The section of this guide dealing with pressures of thermal origin exerted by ice sheets on structures will give additional and current information on the behavior of ice covers.

Thicknesses of Great Lakes Ice

Striegl (1952), Aune, Beaudin, and Borrowman (1957), and others have reported observations on natural ice thicknesses on the Great Lakes. Because of the wide geographic extent of the lakes area, as well as the differences in the characteristics of individual lakes, the ice season varies considerably. Climatic conditions obviously are the major factors in the amount and rapidity of ice formation and disintegration.

The following typical cases indicate the range in temperatures and the differences in length of the cold season which are primary factors in connection with the formation of ice. As these figures are based on normal daily mean temperatures published by the U.S. Weather Bureau it is obvious that there will be wide variations from these means depending on the general character of the season. Please refer to Table 1.

Locality	Normal Lo Freezing From		Total Days	Ave. Temp. During Freezing Season	Freezing Index* F°(C°) Degree Days of Freezing
Duluth	Nov. 12	Apr. 2	142	16.73°F (-8.48°C)	2168 (1204)
Sault Ste. Marie	Nov. 16	Apr. 4	140	18.76 F (-7.36°C)	1854 (1030)
Milwaukee	Nov. 28	Mar. 15	108	24.22°F (-4.32°C)	840 (467)
Port Huron	Dec. 1	Mar. 20	110	24.98°F (-3.90°C)	772 (429)
Buffalo	Dec. 8	Mar. 19	102	26.35°F (-3.14°C)	576 (320)

Table l (Striegl--1952)

It will be noted that these figures indicate 3.76 times as much of the air temperature condition which, when in contact with undisturbed water, causes the formation of ice at Duluth as at Buffalo. Naturally one would expect the formation of more ice at Duluth, and the northerly areas of the Great Lakes, than in the southern areas. Other conditions such as sunshine, winds, volume of the body of water, etc., of course, modify the relationship between temperature and ice formation so there is no direct relationship that may be determined for natural conditions.

^{*}Freezing Index is the total number of Degree Days using the freezing temperature as the reference. A Degree Day is defined as the departure of the daily mean temperature from the freezing temperature.

Aune, Beaudin, and Borrowman (1957) indicate maximum thicknesses of solid lake ice observed in locations from 1899 through 1951 and average daily temperatures for the month of January for a period of 20 years. They described solid lake ice as occurring in protected harbors, bays, and channels of low velocity current. Ice of this type begins to form in November and continues to increase in thickness throughout the winter occasionally obtaining a depth of three and one half feet (1.1 meters) with depths of two feet (0.6 meters) being quite ordinary in northern localities. Solid lake ice conditions usually prevail to the later part of March or early April at which time the ice begins to honeycomb and break up. Their data is presented in Table 2.

Location	North Latitude	Average Daily January Temperature	Maximum Thickness of Ice
Duluth, Minnesota	47 °	10°F (-12°C)	38 in (97 cm)
Marquette, Michigan	46°	15°F (- 9°C)	27 in (69 cm)
Escanaba, Michigan	46°	15°F (- 9°C)	35 in (89 cm)
Green Bay, Wisconsin	44°	17°F (- 8°C)	36 in (91 cm)
Milwaukee, Wisconsin	43 ^c	22°F (- 6°C)	
Chicago, Illinois	42 °	25°F (- 4°C)	
Mackinaw City, Michigan	n 46°	17°F (- 8°C)	
Saginaw Bay, Michigan	43°	22°F (~ 6°C)	35 in (89 cm)
Detroit, Michigan	42°	26°F (- 3°C)	
Cleveland, Ohio	41°	27°F (- 3°C)	17 in (43 cm)
Buffalo, New York	43°	25°F (- 4°C)	24 in (61 cm)
Oswego, New York	43 °	23°F (- 5°C)	25 in (64 cm)
Ogdensburg, New York	45°	17°F (- 8°C)	30 in (76 cm)
Kingston, Ontario	44 °	23°F (- 5°C)	25 in (64 cm)
Sault Ste. Marie, Michigan	46°		27 in (69 cm)
Port Arthur, Ontario	48 °	10°F (-12°C)	41 in (104 cm)

Table 2 (Aune, Beaudin and Borrowman--1957)

Freshwater lakes and rivers lying between $35^{\circ}F$ ($2^{\circ}C$) and $30^{\circ}F$ ($-1^{\circ}C$) January isothermals will freeze on rare occasions of extremely cold temperatures (about $0^{\circ}F$ or $-18^{\circ}C$) which are usually of only several day's duration.

Between $30\,^{\circ}\text{F}$ (-1°C) and $25\,^{\circ}\text{F}$ ($-4\,^{\circ}\text{C}$) isothermals some ice normally less than 4 inches ($10\,\text{cm}$) in thickness can be expected almost every winter. Northward of the $25\,^{\circ}\text{F}$ ($-4\,^{\circ}\text{C}$) isothermal ice can be expected every winter. The thickness will vary from 6 inches ($15\,\text{cm}$) to $12\,\text{inches}$ ($30\,\text{cm}$) in the vicinity of $25\,^{\circ}\text{F}$ ($-4\,^{\circ}\text{C}$) isothermal to $36\,\text{inches}$ ($91\,\text{cm}$) near the $10\,^{\circ}\text{F}$ ($-12\,^{\circ}\text{C}$) isothermal.

During January and February 1977 we made ice thickness measurements in harbors on Lakes Superior, Michigan and Huron. This was an unusually cold winter in the western Great Lakes area. Many thicknesses in the range of 30 to 39 inches (75 to 100 cm) were observed between 42° and 47° North latitudes. At 42° North latitude on Lake Michigan, 35 inches (90 cm) was measured. There seems little doubt but that three feet (one meter) of ice is probable in the western Great Lakes. Also, as was previously mentioned, the U.S. Great Lakes Environmental Research Laboratories is measuring and recording ice thickness in the Great Lakes.

Deformation and Strength Behavior of Polycrystalline Ice

Polycrystalline ice is a viscoelastic material. A comprehensive description of a complete linear viscoelastic model representing the stress-strain behavior is given by Nevel (1976).

For our purposes we will review deformation behavior by describing a laboratory unconfined uniaxial compressive strength test of ice. The test is run at a constant stress, i.e. a single load is placed on the sample. As it remains on the sample, the sample will compress or deform. The amount of deformation divided by the sample length is the strain.

The longer the load is left on the sample, the more the sample will deform. Strain is therefore a function of time. This is termed creep behavior. (Strain is also a function of other variables, e.g. the magnitude of the stress.)

The total creep occurring before the sample breaks is described by three creep periods. These periods are termed primary, secondary, and tertiary.

When the load is first placed on the ice, the ice has an instantaneous elastic response which can be thought of as an elastic compression in a spring. If the load is too great, the sample may break before creep can take place. Assuming however that the load is moderate, the ice will next exhibit a delayed elastic response. This response can be thought of as the travel in a dashpot, i.e. a plunger moving in a viscous fluid. Together these two responses make up the creep period called primary creep.

During secondary creep the strain increases with time and if the ice behaves as a linear viscoelastic material this strain increase is linear.

At some point in time the strain will accelerate. This increasing straining is the tertiary creep. The sample compresses at a faster and faster rate until it fails. At lower stresses ice may not exhibit tertiary creep, but may fail during secondary creep.

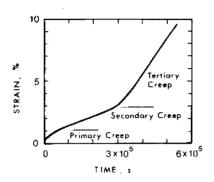


Figure 1 Creep Curve for Granular Ice (Gold, 1973)

Figure 1 is a creep curve for granular ice reported by Gold (1973). The ice was stressed at 145 psi (1 MN/m²). Its temperature was 14°F (-10°C). The primary creep represents about one percent strain. Secondary creep occurred until three and one-half days (3 x 10^5 s) had elapsed. The strain increased to an increasing rate and at the end of about six days had deformed ten percent of its original length.

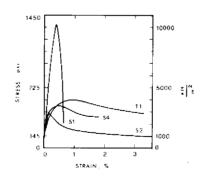


Figure 2 Strain Dependence on Stress (Gold, 1973)

If now instead of applying a constant stress to the sample we apply a constant rate of strain we can obtain stress-strain curves like those in Figure 2 by Gold (1973). Some ice engineering problems involve conditions approaching constant strain. S1 and S2 are columnar ice and were tested with the load perpendicular to the long direction of the grains. It is granular ice and S4 is frazil ice. The temperature of the ice was 15° F (-9.4°C). The strain rate 1.67 x 10^{-5} s⁻¹, i.e. 0.0000167 per second. (The time to reach one percent strain was ten minutes and the time to reach three percent was one-half hour.) For the given strain rate and ice temperature the influence of ice's structure is clearly seen: columnar-grained, type S-1 ice, is appreciably more brittle in its behavior than the other three, and granular, type II ice, more ductile.

To adequately describe the deformation and strength behavior of ice many things must be considered. These include the type of ice, the temperature, the type of load (compression, tension, etc.), load direction with respect to ice structure, the strain rate, and the stress level. Furthermore, to analyze ice in an engineering context it is necessary to establish failure criteria in either ductile or brittle modes of behavior under a variety of boundary conditions. There is much that is not presently understood about ice behavior.

Use of published laboratory strength test data from ice specimens (or simulated ice specimens) must be done with care as the values determined by such testing may not be representative of conditions in the field. For example, uniaxial laboratory strength tests may not represent biaxial and triaxial conditions encountered in the field, because of influence of shear stress in the deformation and failure behavior of ice. Also, small samples generally do not adequately represent conditions in large, sometimes irregular, cracked sheets of ice containing various impurities.

Appended to this guide is an abbreviated summary of a comprehensive monograph entitled, "Deformation and Strength of Ice", by Lavrov (1969). For the reader wishing to delve more into ice strength and behavior it will be of interest. Also the previously mentioned viscoelastic ice model by Nevel (1976) will be very useful.

In subsequent sections of this guide we will introduce strength parameters and design criteria we believe applicable to ice engineering problems in small craft harbors.

PART II

ICE PRESSURES OF THERMAL ORIGIN

Drouin and Michel (1974), working in the Laval University Ice Mechanics Laboratory have prepared a comprehensive work entitled, "Pressures of Thermal Origin Exerted by Ice Sheets Upon Hydraulic Structures". This work, undertaken by Drouin for a doctoral thesis, we believe is the most comprehensive and recent treatment of thermally induced ice forces on structures.

Prior to this work, Michel (1970) prepared a monograph entitled, "Ice Pressure on Engineering Structures". With respect to thermal forces, the later work should be used as it corrects some previous assumptions now proven erroneous.

The material that follows in this section of the guide is largely based on Drouin and Michel (1974). It is recognized that it was originally prepared in the context of reservoir ice forces that affect dams, and that it is a laboratory study. Nevertheless, understanding this work and the ranges of thermally produced loads that can be predicted for dams, will give the harbor designer a feel for the problem he faces. With flexible harbor structures, the forces will be less. Also with real field ice, the forces will be less because of faults, cracks, discontinuities, etc. Laboratory samples test stronger than field ice.

Introduction

Thermally induced ice pressures on dams and other hydraulic structures are significant, and may even be the controlling design load for some structures. In the past these thermal thrusts have been calculated for the condition of ice failing by crushing at stresses like 400 psi (2756 kN/m²). This resulted in thermal loads now known to be too high, e.g. 50 kips per foot (730 kN/m). Now values of $\frac{1}{4}$ to $\frac{1}{2}$ these are normally used. Flexible structures experience even smaller loads than rigid structures, but ways to accurately determine ice loadings on flexible structures have not been developed. The loads would however be expected to be less than those predicted in this section.

During the past fifty years, engineers and scientists, e.g. Royen, Brown and Clarke, Rose, Monfore, Lofquist, and Lindgren researched and estimated thermally induced ice pressures. None of them took into account the crystallographic characteristics of the types of ice and actual strain rates that exist in nature in the regime of thermal expansion of ice sheets. Also the initial temperature of the ice was not considered in some of the previous work, and some research was based on too few tests. The composite picture of past studies is one with ambiguities and many deficiencies. Therefore we won't review this previous work, but instead will present the recent work of Drouin and Michel (1974) who have, however, comprehensively analyzed the work of those before them.

To estimate ice pressures, we must first consider the thermal properties of ice, heat transfer in ice sheets, types of naturally occurring air temperature rises, and the deformation characteristics of several types of ice strained slowly. With a laboratory testing program supported by theoretical rheological mathematical models, Drouin and Michel (1974) have developed estimates for thermally induced ice thrusts.

Thermal Properties of Ice

Conductivity, specific heat, diffusivity and expansion are all properties of ice important to estimating thermal forces. Like most properties of ice their values are still imprecise. Drouin and Michel (1974) estimate the most probable values of thermal properties from the works of others. Table 3 presents their estimated values.

Temperature °F (°C)	Thermal Conductivity Btu in hr ⁻¹ ft ⁻² F-1 (W m-1 °C-1)	Specific Heat Btu 1b ⁻¹ °F-1 (J kg-1°C-1)	Thermal Diffusivity ft ² hr ⁻¹ (cm ² hr ⁻¹)	Thermal Expansion F -1 (°C -1)
32°	15.6	.506	.0449	30.0×10^{-6}
(0°)	(2.25)	(2117)	(41.7)	(54.0 × 10 ⁻⁶)
14°	16.1	.487	.0481	29.0×10^{-6}
(-10°)	(2.32)	(2039)	(44.7)	(52.2 × 10^{-6})
-4°	16.8	.468	.0522	28.0×10^{-6}
(-20°)	(2.42)	(1961)	(48.5)	(50.4 × 10 ⁻⁶)
-22°	17.6	.450	.0570	27.0×10^{-6} (48.6×10^{-6})
(-30°)	(2.54)	(1883)	(53.0)	
-40°	18.6	.431	.0629	26.0×10^{-6} (46.8×10^{-6})
(-40°)	(2.68)	(1805)	(58.4)	

Table 3 Thermal Properties (after Drouin and Michel, 1974)

In the above table, the thermal diffusivity (equal to the conductivity divided by the factor density times the specific heat) was computed assuming an ice density of 57.2 pcf (916 kg/m 3). Ice containing air bubbles may easily have a lower density, say as low as 49 pcf (785 kg/m 3).

Heat transfer through a porous medium, like bubbly ice, will be less because the air cells offer greater resistance to heat transfer than does the solid ice. The value for thermal conductivity becomes less as the amount of entrapped air increases.

Notwithstanding the inherent variations in the thermal properties of ice, we can adequately proceed, within the range of design accuracies, to make engineering estimates of ice thrusts.

Heat Transfer in Ice Sheets

The studies and solutions by Drouin and Michel (1974) do not take into account any thermal boundary layers (air-ice and air-snow) or any solar radiation on the ice. However, they believe their pressures calculated on the basis of heat transfer solely by conduction are of the proper order of magnitude, and perhaps greater than the pressure calculated with the air-ice boundary layer and absorption of solar radiation taken into consideration.

Air-ice or air-snow boundary layers can contribute to increasing the initial temperature of the surface of an ice sheet. Absorption of solar radiation by the ice adds to the heat transferred by conduction. The rates of temperature increase at various levels in the ice would then be higher than those computed on the basis of conduction alone. Increasing the initial surface temperature results in a decrease in the stresses throughout the ice sheet. Increasing the heat transferred by convection (because of solar radiation) results in an increase in stresses in the upper portions of the ice sheet and in shorter times to reach maximum stress levels which thereafter diminish rapidly. However, an increase in the average rate of temperature increase of the surface of an ice sheet implies a decrease in the maximum pressure. Hence their pressures are thought to be conservative, i.e. greater than necessary for design of conditions actually present in the field.

In the foregoing explanation it should be noted that an increase in the rate of temperature increase in the surface of the ice sheet results in a decrease in the maximum pressure. Drouin and Michel (1974) found, contrary to many other published theories and results, that the pressure in an ice sheet was highest for small rates of temperature increase at the surface of the ice. From a physical standpoint, they explain this by the deeper penetration into the ice sheet of the temperature variations at the time when the stress at the surface attains its maximum. It is in the top zone of an ice sheet where the stressed condition related to the maximum pressure exerted by thick ice sheets develops. Attenuation of temperature variation in the interior of an ice sheet is very rapid.

The stress in an ice sheet is determined by the temperature distribution in the ice sheet as a function of the variation of air temperature. Drouin and Michel (1974) use a sinusoidal air temperature variation, rather than a linear variation or step function. The use of sinusoidal variation produces larger stresses and fits well with natural temperature variations (in Canada).

The sinusoidal variation used is one where the temperature increases at an increasing rate until it has achieved one half of its total rise. This occurs when one half of the time to achieve its total rise has passed. The temperature continues to rise but at a decreasing rate until the total temperature rise has occurred for the period.

Other factors that affect the temperature variation in an ice sheet include the presence of snow on the ice sheet, the thickness of the ice sheet, increase of thickness in the ice sheet as a function of time, solar radiation absorbed, and variable thermal properties of ice. (These vary primarily with temperature.)

Analytic methods are available to describe temperature distribution in a mass as a function of the variation in temperature of an ambient medium. To solve the equations developed by these analyses requires simplifying assumptions. For example, Drouin and Michel (1974) had to assume that the thickness of the ice remained constant during the period of ambient temperature increase. This would not be the case for thin ice but would be reasonable for sheets thicker than 16 inches (40 cm).

Climatological Data

In order to compute the pressures from thermal ice expansion we need to know the total temperature rise and the time required to obtain that rise. Also we should know if the temperature rise is linear or varies somehow, such as sinusoidally. Drouin and Michel (1974) present climatological data for Quebec City from 1944 to 1967, inclusive. Quebec City is at 46°48'N latitude and 71°23'W longitude. Climatological data from Quebec City will not be applicable at most other sites, but by presenting them here we will be able to subsequently demonstrate a design methodology. The designer of a small craft harbor at another location would of course investigate climatological conditions there.

Table 4 presents a statistical analysis of data for 23 winters at Quebec City. For that city it was calculated that 90 percent of the air temperature increases had durations of less than 20 hours. About two-thirds of the time the duration was less than 10 hours. The most frequent temperature spread was 17°F (9.5°C) with a duration of 7 hours, which is an average rate of air temperature increase of 2.4 F/hr (1.3 C/hr). Sinusoidal variation in temperature rise was generally present.

Period of	Duration of Temperature Increase (hrs)					
Recurrence (years)	<u>5</u>	<u>10</u>	<u>20</u>			
12.5	4.4 (2.4)	3.1 (1.7)	1.6 (0.9)			
25	4.7 (2.6)	3.3 (1.8)	1.7 (0.9)			
50	5.2 (2.9)	3.5 (1.9)	1.8 (1.0)			
100	5.6 (3.1)	3.7 (2.1)	1.9 (1.1)			

Table 4 Rates of Increase, F/hr (C/hr), for Linear Air Temperature Increases as a Function of the Period of Recurrence and Duration of the Increase (Quebec, December 1, 1944—March 31, 1967)

(after Drouin and Michel, 1974)

Rheological Aspects

When an ice specimen, mechanically fixed at its ends, is heated it will try to expand. However, since the ends are fixed a thermal stress will be produced. The strain from this thermal stress offsets that due to the temperature rise. Thus, even though the specimen undergoes no apparent change in length, it is subjected to mechanical strain equal to the thermal deformation that would have resulted if it had been unrestrained. Maintaining a constant length and measuring the forces (stresses) that occur at different temperatures is one way of carrying out laboratory tests on ice.

The laboratory procedure used by Drouin and Michel (1974) used a somewhat different approach, one that avoids errors caused by the metal parts of the testing machine expanding. The tests were performed, at different temperatures, on ice specimens at constant strain rates. The tests determined by means of the experimental curves obtained, the stresses and thrusts induced by the effects of temperature variation.

Ice in nature responds to thermal changes in a ductile manner. The strain rates are quite small and hence ice exhibits ductile behavior. The small strain rates accompanying thermal behavior can be shown with the following calculation.

Assume a temperature rise of 5 F/hr (2.8 C/hr) and a coefficient of thermal expansion of 28 x 10^{-6} F-1 (50.4 x 10^{-6} C-1)

Strain Rate = Coefficient of Expansion x Rate of Temperature Rise

$$= \frac{(28 \times 10^{-6} \text{ F}^{-1}) (5 \text{ F/hr})}{(3600 \text{ s/hr})}$$

$$= 3.89 \times 10^{-8} \text{ s}^{-1}$$

The thermal strain is also quite small.

Assume a rise in temperature from -30°F (-34°C) to 32°F (0°C) at the above rate. The strain rate is 3.89 x 10^{-8} s⁻¹ for a coefficient of thermal expansion of 28 x 10^{-6} F⁻¹ (50.4×10^{-6} C⁻¹)

Strain =
$$\frac{(62 \text{ F}) (3600 \text{ s/hr}) (3.89 \times 10^{-8} \text{ s}^{-1})}{(5 \text{ F/hr})}$$

= 0.00174

= 0.174%

The total strain is less than one quarter of one percent. However, if for this strain the ice were unrestrained, it would expand nearly 10 feet (3 m) in a mile (1.6 km)

To deform laboratory specimens by a factor of one percent at constant rates of 10^{-7} , 10^{-8} , and 10^{-9} per second, test durations of 28, 280, and 2800 hours are necessary. These long periods explain why few tests in the range of strain rates that accompany thermal expansion in the field are cited in the literature. At very low strain rates even small temperature variations in the ambient medium impair the validity of a test. Also, the deformation of a measuring element of a load cell, is not negligible in comparison with the actual deformation of the specimen. The laboratory testing program undertaken by Drouin and Michel (1974) overcame many difficulties of this nature.

Laboratory Test Results, Ice Thrusts

Two types of ice were tested, S1 columnar ice loaded perpendicular to the c-axis (and parallel to the basal plane) and T1 snowpack ice. Samples were generally columns 1 inch (2.54 cm) in diameter and 3 inches (7.62 cm) long. As they were compressed, at strain rates ranging between 10^{-7} , and 10^{-8} per second and at different temperatures ranging between freezing and about -20° F (-29° C), the deformations and loads were measured.

In the snowpack ice the maximum loads occurred after the stress had increased at a constant rate for some time. The stress, after having attained its maximum, usually remained constant or nearly so. The average deformation to reach maximum stress was 0.15 percent.

In the columnar ice the stress increased at a constant rate for a long time. Following that, the rate of increase of stress gradually diminished until it was nil. The maximum stress was reached when the specimen had been deformed between 0.08 and 0.20 percent. The stress, after reaching its maximum value, fell very rapidly and then stabilized at a near constant value.

At the low strain rates the sample behavior is plastic, i.e. the sample doesn't break or shatter, but yields as time proceeds.

The maximum thrusts for columnar and snowpack ice are given in Table 5. They have been computed from the laboratory test data and a sinusoidal varied ambient temperature rise from an initial linear steady temperature state in the ice (i.e. the initial ice temperature profile varies linearly from the ambient temperature at the surface to the melting point at the water-ice interface). With time the ice warms up and thereby exerts thermal stresses. These stresses reach maximums and the stress in the ice at various depths can be summed up to obtain the maximum ice thrust per unit length for a given ice thickness, total temperature rise, and duration of temperature rise.

Ice thrusts for ice thinner than indicated in the table will be less than the tabulated values. Also, the analysis of stresses indicates that the stresses in an ice sheet are particularly high in the first 8 to 12 inches (20 to 30 cm) of the thickness of a thermally stressed ice sheet. The bottom region of a thick ice sheet may consist of a different type of ice without appreciably altering the stressed conditions, the heat transfer being the same.

Ice Surface Temperature, F (C) and Classification of Ice:
S1 = Columnar Ice and T1 = Snowpack Ice

	S 1	= Columnar	ice an	d 1 =	Snowpack I	ce
Ice Thickness and Duration of Temperature Increase (hrs)	14°F (-10°C)	-4°F (-20°C)	-22°F	(-30°C)
	<u>S1</u>	<u>11</u>	<u>S1</u>	<u>T1</u>	<u>51</u>	<u>11</u>
20 in (50 cm)						
t = 5 hrs	5	4	11	8	19	13
	(73)	(58)	(160)	(117)	(277)	(190)
t = 10 hrs	6	5	14	11	23	16
	(88)	(73)	(204)	(160)	(336)	(233)
t = 20 hrs	9	7	18	13	27	18
	(131)	(102)	(263)	(190)	(394)	(263)
30 in (75 cm)						
t = 5 hrs	5	4	11	8	20	14
	(73)	(58)	(160)	(117)	(292)	(204)
t = 10 hrs	7	5	15	11	24	18
	(102)	(73)	(219)	(160)	(350)	(263)
t = 20 hrs	9	8	20	15	30	20
	(131)	(117)	(292)	(219)	(438)	(292)
40 in (100 cm)						
t = 5 hrs	5	4	11	8	20	15
	(73)	(58)	(160)	(117)	(292)	(219)
t = 10 hrs	7	5	15	12	25	19
	(102)	(73)	(219)	(175)	(365)	(277)
t = 20 hrs	9	8	20	16	32	22
	(131)	(117)	(292)	(233)	(467)	(321)

Table 5 Pressures of Thermal Origin by an Ice Sheet Restrained in One Direction kips/ft (kN/m)

(after Drouin and Michel, 1974)

Use of the table is illustrated by the following example. Given a 30 inch (75 cm) ice sheet with a surface temperature of $-4^{\circ}F$ ($-20^{\circ}C$). The thrust exerted by a rise in temperature to $32^{\circ}F$ ($0^{\circ}C$) over a period of 10 hours (i.e. at a rate of 3.6 F /hr or 2.0 C /hr) would be 11 to 15 kips per foot (160 to 219 kN/m) depending on whether the ice was granular or columnar (loaded perpendicular to the c-axis).

Except for complete biaxial restraint, most factors encountered in nature tend to decrease calculated ice thrusts. Among the factors which reduce stresses of thermal origin are the insulating qualities of a snow cover on the ice and cracks in the surface of the ice.

Biaxial Restraint

Determining rheological behavior of ice deformed in two directions is a subject of current research. Drouin and Michel (1974) performed some biaxial tests on snowpack ice and concluded that snowpack ice deforms in the manner of an almost perfectly plastic material when the maximum biaxial stress of thermal origin is obtained. This maximum biaxial stress is about twice the maximum stress previously obtained in the uniaxial tests. Also they concluded that maximum stresses for S1 and S2 columnar ice having pressures of thermal origin and biaxially restrained can be adequately estimated by using the uniaxial test results for S1 ice.

They give no ice thrusts for the biaxial condition as such, but suggest using uniaxial thrusts for a dam and reservoir with the following example. Assume an ice sheet is totally restrained in a direction perpendicular to a dam crest. Stresses of great magnitude develop in the first 8 inches (20 cm) of the ice thickness. Under natural conditions, the crystallographic orientation of the grains of columnar ice in the top layer of an ice sheet can be considered random. In such case, the stresses developed are considerably lower than if all the columns have, for instance, a vertical crystallographic orientation. Experimentally the biaxial restraint of S1 or S2 columnar ice having some grains randomly oriented gives substantially the same maximum stress values as are obtained from S1 ice specimens restrained in only a single direction (deformed perpendicular to the c-axis). Also a reservoir with rocky vertical walls around the entire perimeter would be unusual. It would correspond to an ice sheet in a swimming pool (or a small-craft harbor completely surrounded by vertical sheet piling walls).

Within the present state of knowledge, Drouin and Michel (1974) recommend using thrusts in reservoirs on laboratory testing of ice under uniaxial loading.

Effects of Cracks and Snow on Ice Thrusts

Snow covered ice or cracked ice or both significantly reduce the thrust of an ice sheet.

The effect of cracks (either from thermal contractions or water level changes) is to mitigate the effect of a temperature rise or spread.

Metge (1976) has reported on a five year study and field observation on thermal cracks at Kingston, Ontario. The frequency of cracking tends to decrease as an ice cover becomes thicker. He categorizes thermal cracks into three groups. The first are dry cracks which absorb a significant amount of thermal ice movement. They are common and consist of cracks that extend from the top of the sheet down one-half to two-thirds of the ice thickness. These dry cracks open and close according to the ice temperature. The second and third groups are respectively, wet cracks which are relatively rare, and wide wet cracks which are just too wide to easily refreeze into sound ice. Wet cracks are narrow enough to refreeze rapidly and in so doing add material to the ice sheet. Since wet cracks are not prevalent, we are mainly concerned with the beneficial effects of the dry thermal cracks. The following example illustrates their effect on thermal thrusts.

Assume a coefficient of thermal expansion of 28 x 10^{-6} F-l (50.4 x 10^{-6} C-l) and a total summation of crack widths of 0.02 inches per yard (0.56 mm/m)

Strain (%) = (Coef. of Expansion) (Temp. Rise) (100)

 $\frac{(0.02 \text{ in/yd})(100)}{(36 \text{ in/yd})} = (28 \times 10^{-6} \text{ F}^{-1})(100) \text{ (Temp. Rise F)}$

Temp. Rise = 20 F (11 C) = Equivalent Temperature Spread

In this example, the first 20 F $(11\ \text{C})$ of a temperature rise would be taken up by the cracks. The remainder of the rise would produce thermal thrusts.

Table 6 gives temperature spreads for cracked ice.

Summation of Crack Widths per Unit of Length of an Ice Sheet		Total Deformation or Strain of the Ice	Equivalent Temperature Spread		
<u>in/yd</u>	mm/m	<u>%</u>	<u>F</u>	<u>C</u>	
0.009	0.25	0.025	9	5	
0.018	0.5	0.05	18	10	
0.036	1.0	0.10	36	20	
0.054	1.5	0.15	54	30	
0.072	2.0	0.20	72	40	

Table 6 Summation of Crack Widths Per Unit of Length of an Ice Sheet Expressed as Percent Thermal Deformation of the Ice and Equivalent Temperature Spreads (Coef. of Expansion = 28×10^{-6} F⁻¹ or 50.4×10^{-6} C⁻¹)

(after Drouin and Michel, 1974)

Snow accumulated on an ice sheet has a drastic effect on the temperature variations in the ice, and as a result the pressures of thermal origin are much less. A uniform fresh snow cover can drown an ice sheet causing no thrust until the ice reforms on the top. This can occur when a snow fall roughly equal to the ice thickness occurs.

Snow, a complex material which undergoes a continuous process of metamorphosis, can be converted to an equivalent thermal thickness of ice by proportioning through the coefficients of thermal conductivity for snow and for ice. However, the thermal conductivity of snow varies with density (as well as other things). Mellor (1964) gives snow conductivities as a function of density. Drouin and Michel (1974) report the average density of snow 20 to 40 inches (50 to 100 cm) thick on the ground to be 25 to 28 pcf (400 to 450 kg/m 3). Snow at this density will barely show a man's foot print. For snow layers of small thicknesses the density is much less—lower than 12 pcf (200 kg/m 3).

Table 7 gives equivalent ice thicknesses for snow densities.

	Ratio of Thermal						
Snow Density	Equivalent Thicknesses of Ice to the						
$pef(kg/m^3)$	Actual Thicknesses of Snow						
6.2 (100)	24.0						
, ,							
9.4 (150)	17.1						
12.5 (200)	13.2						
15.6 (250)	10.6						
18.7 (300)	8.7						
21.8 (350)	7.2						
25.0 (400)	6.1						
28.1 (450)	5.0						

Table 7 Ratios of Thermally Equivalent
Thickness of Ice to Actual Thicknesses of Snow
(Drouin and Michel, 1974)

To illustrate the use of the above table, assume we have 6 inches (15 cm) of snow weighing 21.8 pcf (350 kg/m 3) on top of a two foot (0.6 m) thick ice sheet. The ratio of equivalent thickness of ice to actual thickness of snow is 7.2; therefore, thermally the snow simulates an additional layer of 43.2 (110 cm) thick. The fictitious sheet is now 67 inches (170 cm) thick. The temperature variations to be taken into account in calculating the stresses of thermal origin are those which prevail 43 inches (110 cm) below the surface of the fictitious sheet. At this depth into an ice sheet the attentuation of temperature is complete and thus the pressure of thermal origin is very small or nil.

Ice Thrust on Individual Piers

Petrunichev (1954), cited by Korzhavin (1971), considers static horizontal pressure of ice on individual piers. The question is raised as to whether one should determine the pressure from an ice field by the width equalling half the sum of the adjacent spans or be limited to considering only the pier width.

The deformation of the part of the ice field lying between the piers would not occur separately. It exerts a partial effect on the amount of pressure on a pier. It would be increased. Petrunichev (1954) suggests that the total force on a pier should be equal to the thrust (force per unit length) times the sum of the width of the pier and one-third the half sum of the spans contiquous with the pier.

The suggestion for somehow increasing forces for 'bridging' action would be applicable for small craft harbors. For example, consider the case of a marina pier constructed with horizontal elements spanned across intermittent cribs. The total lateral force would exceed that computed from ice contact with the crib alone. In part, the ice would 'bridge' over and produce a larger horizontal force.

Closure

In this section we have reviewed the thermal ice pressures proposed by Drouin and Michel (1974) and presented some estimates of thrusts. Flexible marina structures and conditions existing in nature should result in forces smaller than those given.

The forces were computed for assumed temperature rises equal to the difference in the cold ice sheet surface temperatures and the melting point. Temperature rise is a function of the site's climatology. The highest pressure occurs for the smaller rates of temperature rise. Also the attenuation of temperature rise occurs very rapidly and it is the top zone of the ice sheet that experiences the thermal stresses. (Query, can one mitigate these stresses in this area somehow. If the ice were weakened or broken, or if made thin or warmer, or if separated from the structure with a pliable material, the force on the structure would be reduced.)

Strain rates in the regime of thermal expansion of ice sheets are very low (less than $10^{-7}~\rm s^{-1}$) and the ice behaves in a ductile manner. Also forces produced from uniaxial testing can adequately represent some level of biaxial conditions that might be present at a particular site.

Each site where thermal pressures will exist should be separately studied. The climatology, characteristics of the ice sheet, bank configuration. and type of structure are all important.

PART III

ICE SUPPRESSION AND WEAKENING

Ashton (1974) has prepared a monograph on air bubbler systems to suppress ice. Analytical methods were used to develop a procedure for predicting the effectiveness of using bubbler systems to suppress ice formation under various field conditions. We believe air bubbling is an effective method to suppress ice. The material contained in this section is based largely on the procedures suggested by Ashton but simplified somewhat to fit our concern, namely, melting ice in a small-craft harbor (rather than in a lake shipping channel or a deep-water port). The resulting suppression predictions agree well with our field observations and experiences of air bubbler manufacturers.

Those wishing to examine air bubbler systems in more detail than what is presented here can refer to the original Ashton (1974) monograph, Ashton (1975) and Ashton (1977). The 1975 work is a laboratory experimental report on heat transfer coefficients associated with flow induced by a line source bubbler system. The 1977 work is a refinement of the monograph and a numerical simulation for a system. US CRREL has developed a computer program in BASIC language that simulates analyses of diffuser lines, plumes and heat transfer, ice melting, and thermal reserves. The simulation agrees satisfactorily with one example case for a winter in Duluth, Minnesota.

We will begin with a description of how a compressed air bubbler deicing system works. This will be followed with a review of the suggested design method and practical considerations to achieve a successful installation. Also included are a few notes on other ways to suppress or weaken ice.

Principles of Bubbler System Operation

Figure 3 is a cross sectional view along the axis of an air diffuser pipe on a lake bottom. It is also representative of the action of a point source diffuser.

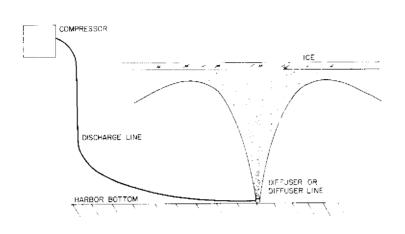


Figure 3 Bubbler System Schematic (after Ashton, 1974)

Air is compressed by a compressor either located indoors or suitably protected and mounted on a marina pier. In an overall ice suppression system, several compressors may be used and interconnected through a discharge manifold piping system. Air travels through the discharge line to the diffuser or diffuser line on the lake bottom. The air has been warmed during compression but normally is cooled down to the ambient lake temperature by the time it reaches the diffuser. Here the air is discharged through an orifice. It must have enough pressure at the diffuser to overcome the depth of water hydrostatic pressure.

The discharging air is cooled as it expands going through the orifice. If it expands too quickly an icing condition can occur at the orifice. With proper air pressures this does not happen. The momentum of the air jet sets the bubbles in motion. The momentum quickly dissipates and bubble buoyancy takes over. As the bubbles rise they entrain water into the rising plume from the lateral direction. Many bubble sizes are present, but the tinier ones are more efficient as they move more water for the same volume of air. In fact, this principle explains why large air bubbles, like those produced by belcher devices, aren't as effective as a continuous stream of tiny bubbles. The ice melting that occurs on the underside of the sheet is the result of both temperature and volume of water being moved upwards from the warmer bottom water by the bubble plume.

At some point near the surface the buoyant plume spreads as it either impinges on the ice cover or encounters open water. If a free water surface is present, the bubbles will escape directly to the atmosphere. This results in a certain amount of heat being wasted. If a cover exists the bubbles will move laterally along the underside of the ice. As they do, melting primarily by convection occurs. This heat loss results in cooling of the flow whose velocity has also decreased as it spread out. Rapid decay in the heat transfer rate is the result.

Finally, the plume imposes a net circulation on the water and this allows more warm water to be drawn into the area from distant lateral directions. A bubbler system would not work in a swimming pool for example, because the amount of warm water is limited by the pool's volume and eventually it would all be cooled down. The pool would freeze to the bottom. Design of bubblers in sea water is complicated because salinity may vary with depth and theoretically the maximum density occurs at the freezing point. (There may be no "warm" dense bottom water.) Also rivers with currents tend to be isothermal and even super-cooled, and hence do not have enough heat. What heat there may be is already impinging on the river cover.

In a fresh water or brackish water marina, however, a bubbler system can work if lateral warm water recharge is present and it is designed and maintained properly.

Ice Suppression Bubbler Systems Design

The following design procedure is a trial-and-error procedure based upon Ashton (1974). For a given site and conditions the quantity of air required \mathbb{Q}_a is estimated for a tolerated ice equilibrium thickness \mathbf{n}_e . This thickness is a steady state thickness of an ice cover that will persist with a continual flow of heat to the cover's lower surface and a cooler atmosphere forming ice. At the equilibrium thickness, the ice is melting as fast as it forms.

The design of an ice suppression system requires the selection of this thickness. To reduce it to zero or near zero requires disproportionate amounts of heat, i.e. the air flows required become large. The ice itself is an insulator and a free water surface is heat wasteful. (However, the designer may wish to choose to obtain free water in order to better monitor the bubbler's operation throughout the winter.) In selecting a thickness to be tolerated, factors like the following must be considered: first cost and operating costs, resistance available to ice uplift through embedment of piles being protected, magnitude of lateral forces from thicknesses of ice, availability of manpower to chop ice during severe cold periods, temperature extremes existing at the site, and the amount of damage to be tolerated.

The quantity of air required Q_a is estimated from experience and local conditions. Table 8 gives heat transfer coefficients as a function of water depth H and air flow rates Q_a . (Q_a is in terms of volume per time per length of diffuser which reduces to area per time.) The depth of water

used is from the underside of the ice (nominally the top surface for suppressed ice) to the diffuser level (on the bottom generally).

Water Depth		Air Flo	w Rate,	Q _a , ft ² mi	in-1 (m2 s	₅ -1)	
H, ft (m)	.01 (1.55)*	.02 (3.10)	.03 (4.64)	.04 (6.19)	.05 (7.74)	.06 (9.29)	.07 (10.8)
6 (1.8)	151 (857)		180 (1022)		195 (1107)	201 (1141)	206 (1170)
8 (2.4)	142 (806)	158 (897)	169 (960)		184 (1045)	189 (1073)	
10 (3.0)	134 (761)		159 (903)			178 (1011)	
12 (3.7)	127 (721)		151 (857)	158 (897)		169 (960)	
14 (4.3)	121 (687)	135 (767)	144 (818)			162 (920)	166 (943)
16 (4.9)	116 (659)	130 (738)	138 (784)		150 (852)		159 (903)

*Multiply values by 10^{-5}

Table 8 HEAT TRANSFER COEFFICIENTS,
$$h_b$$

Btu hr^{-1} ft⁻² F^{-1} (W m^{-2} C^{-1})
(after Ashton, 1974)

The heat transfer rate $q_{\boldsymbol{W}}$ is obtained from

$$q_w = h_b (T_w - T_m)$$

where $\mathbf{T}_{\mathbf{W}}$ is the water temperature

and T_{m} is the melting point temperature of ice

The water temperature can be measured at a given site or conservative estimates made for its value. We have made water temperature measurements in many western Great Lakes small craft harbors and found them to be isothermal and only slightly above the freezing temperature ($\frac{1}{2}$ F, 0.3 C, or cooler). Others have also reported very cold water. It is reasonable to assume that the water temperature is constant throughout its depth because being shallow and agitated by rising bubbles are conditions tending to destroy any thermal stratification.

Figure 4 gives the equilibrium thicknesses n_e for the calculated heat transfer rate q_{Ψ} as functions of the ambient air temperatures T_a . It has been assumed that no snow cover is present and that the average wind speed is 10 mph (4.5 m s $^{-1}$). If snow is present the equilibrium thicknesses become smaller for a given q_{Ψ} . Conversely, if windier conditions prevail the equilibrium thicknesses increase.

The average daily temperature during the period of the winter under consideration can, for most sites, be used for Γ_a . Ice once formed grows rather slowly during cold spells. During the day, warmer temperatures counteract cooler evening temperatures.

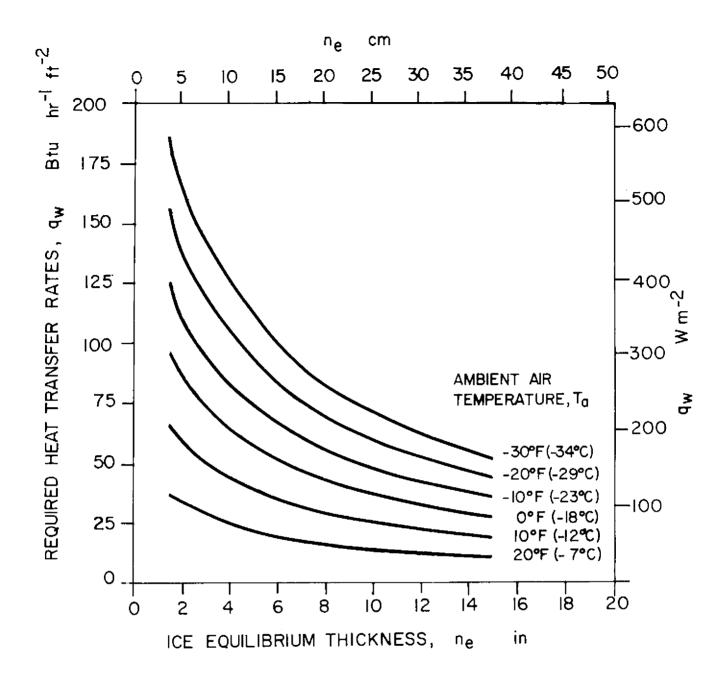


Figure 4. REQUIRED HEAT TRANSFER RATES TO ATTAIN EQUILIBRIUM THICKNESS AS FUNCTION OF AMBIENT AIR TEMPERATURE (after Ashton, 1974)

The use of Table $8\,$ and Figure $4\,$ is illustrated by the following example.

Assume Q_a = 0.06 ft² min⁻¹ (9.29 x 10⁻⁵ m² s⁻¹) of air

H = 10 ft (3 m) water depth T_w = 32.5°F (0.3°C) water temperature T_a = -10°F (-23°C) ambient air temperature

From Table 8 at Q_a = 0.06 and H = 10 find the heat transfer coefficient h_h

$$h_b = 178 \text{ Btu hr}^{-1} \text{ ft}^{-2} \text{ F}^{-1}$$
(1011 W m⁻² C⁻¹)

Find q_{W} , the heat transfer rate from

$$q_w = h_b (T_w - T_m)$$

= 178 Btu hr-1 ft-2 F-1 (32.5°F - 32.0°F)
= 89 Btu hr-1 ft-2 (281 W m-2)

From Figure 4 at $T_a = -10^{\circ} F$ and $q_w = 89$ find $n_e = 4$ inches (10 cm)

Should the ambient temperature (approximately the average daily temperature) fall below -10°F (-23°C) say to -30°F (-34°C) for a few days, then the ice equilibrium thickness would increase to 7 inches (18 cm). Whether 7 inches (18 cm) of ice would be tolerable would depend on the site and what is being protected.

Now assume instead of Q_a = 0.06 ft² min⁻¹ (9.29 x 10⁻⁵ m² s⁻¹) we use one-half as much air. From Table 8 at Q_a = 0.03 ft² min⁻¹ (4.64 x 10⁻⁵ m² s⁻¹) h_b = 159 Btu hr⁻¹ ft⁻² F⁻¹ (903 W m⁻² C⁻¹). For ½ F (0.3 C), q_w = 80 Btu hr⁻¹ ft⁻² (252 W m⁻²). from Figure 4 , at an air temperature of -10°F (-23°C) we find n_e = 5 inches (13 cm) and at -30°F (-34°C), n_e = 9 inches (23 cm). Changing the quantity of air does not have a linear effect on the ice thickness. If however, the water temperature is cooler (or warmer) than estimated, large changes in n_e occur.

Assume the water temperature is $32\frac{1}{4}^{6}F$ (0.1°C) instead of $32\frac{1}{2}^{6}F$ (0.3°C). For $Q_{a}=0.06$ ft² min⁻¹ (9.29 x 10^{-5} m² s⁻¹) $h_{b}=178$ Btu hr^{-1} ft⁻² F⁻¹ (1011 W m² C⁻¹). The required heat transfer rate is $(\frac{1}{4})$ (178) = 45 Btu hr^{-1} ft⁻² (142 W m⁻²). Corresponding equilibrium thicknesses at $-10^{6}F$ ($-23^{6}C$) and $-30^{6}F$ ($-34^{6}C$) would be 11 inches (28 cm) and more than 15 inches (38 cm). These ice conditions would probably be intolerable.

The above analysis is for a line source diffuser system. A solution to a point source diffuser is not presently available. It is estimated that a series of point sources (such as would exist in a bubbler system layout protecting individual piles spaced throughout a marina) would approximate the line source condition. This of course would be a function of how far apart the point sources become. But in general, adding up the air discharges at each point source and dividing by the distribution diffuser length would be a reasonable approximation for \mathbb{Q}_a in a typical small craft harbor.

We can estimate the time required to melt out an ice cover if one is already in place. For thicker ice there is little tendency for it to thicken. The simplest analysis occurs when we assume there is no conduction through the cover and that it is melted by the warm water impinging on its lower surface.

For these conditions the melting rate can be calculated as follows:

Melting Rate =
$$\frac{q_{\psi}}{p_{;} ?}$$

 p_{\odot} is the mass density of ice and equal to 57.2 lbs ft $^{-3}$ (916 kg $\text{m}^{-3})$

 \nearrow is the heat fusion of ice and equal to 143.7 Btu $1b^{-1}$ (3.34 x 10^5 J kg⁻¹)

Find q_{w} to melt one inch (2.5 cm) in a day.

$$\begin{array}{l} 1 \text{ in da}^{-1} = \frac{ (q_{\text{W}} \text{ Btu hr}^{-1} \text{ ft}^{-2}) \ (12 \text{ in ft}^{-1}) \ (24 \text{ hr da}^{-1}) }{ (57.2 \text{ lbs ft}^{-3}) \ (143.7 \text{ Btu lb}^{-1}) } \\ \\ q_{\text{W}} = 28.5 \text{ Btu hr}^{-1} \text{ ft}^{-2} \ (89.8 \text{ W m}^{-2}) \\ \end{array}$$

Therefore with 100 Btu hr-1 ft-2 (315 W m-2) we can melt $3\frac{1}{2}$ inches (9 cm) of ice per day.

We can also make an estimate of the exhaustion of the thermal reserve in a closed body of water. The following computations show the absolute necessity of having a warm water recharge at any small craft harbor using air diffusers to suppress the ice.

We can compute the volume of warm water necessary to melt a given volume of ice as follows. Assume we have ice at the freezing point and warm water at $33\,^\circ\text{F}$ (0.55 $^\circ\text{C}$). On a cubic foot (0.0283 m³) basis then:

Ratio of Volumes = <u>Heat Required to Melt Ice</u> Heat Content of Water

$$= \frac{p! \lambda}{p_W c_p (T_W - T_m)}$$

where p; and ≥ are defined above

 p_{ψ} is the mass density of water and equal to 62.4 lbs ft $^{-3}$ (1000 kg $\text{m}^{-3})$

 c_p is the specific heat capacity of water and equal to 1.0 Btu $1b^{-1}~~{\rm F}^{-1}~(4187~{\rm J~kg}^{-1}~{\rm C}^{-1})$

 $(\text{T}_{\text{W}} - \text{T}_{\text{m}})$ is the difference in the temperature of the water and the melting point of ice

=
$$\frac{(57.2 \text{ lbs ft}^{-3}) (143.7 \text{ Btu lb}^{-1})}{(62.4 \text{ lbs ft}^{-3}) (1.0 \text{ Btu lb}^{-1} \text{ F}^{-1}) (1^{\circ} \text{ F})}$$

= 132

It takes 132 cubic feet $(3.74~{\rm m}^3)$ of water at 33 F $(0.55^{\circ}{\rm C})$ to melt 1 cubic foot $(0.0283~{\rm m}^3)$ of ice at the freezing point.

Assume we have an air diffuser system installed in a closed body of water that is at a temperature of $33^{\circ}F$ (0.55°C). Also assume a depth of water of 10 feet (3 m), diffuser lines spaced 30 feet (9 m) on centers, and that we want to use the diffuser lines to melt ice in a band 3 feet (1 m) wide.

Suppressed Thickness =
$$\frac{(30 \text{ ft}) (10 \text{ ft}) (1 \text{ ft})}{(132) (3 \text{ ft}) (1 \text{ ft})}$$

= 0.76 ft or 9 in (23 cm)

Under these closed conditions the thermal reserve would be exhausted in a few days. If an ice cover grows rapidly a large amount of heat can be sealed in. However it would still not be enough to operate a suppression system throughout a winter. There must be a source of warm water. Also operating these systems in shallow water, say six feet (2 m) or less, is not recommended because of lack of enough warm water to melt an ice cover.

The detailed design and layout of an ice suppression bubbler system is outside the scope of this ice engineering guide. A few comments on the design of systems that are presently installed and are working are in order however.

Having determined the required amount of air (a function of water depth, water temperature, water body thermal reserves, atmospheric conditions, and thickness of ice to be tolerated) an air manifold distribution system together with branch diffuser lines would be designed in accordance with standard air flow design procedures. It is important that manifold and diffuser lines are large enough to reduce pressure losses and to make sure the air supply is not exhausted before the last diffuser is reached. Balancing air flows and pressures is quite important for a successful compression system.

The total air pressure required will be a function of the line and orfice losses and the hydrostatic head to be overcome at the diffuser. Where the harbor bottom is sloping, it is best to distribute the air in an uphill direction from the deep end; otherwise if the air is first emitted in shallow water, the deeper diffusers may be starved.

Spacing of orifices would normally be between 1/2 and 1/3 the water depth along a line diffuser. Pre-slit vinyl weighted bubbler hoses are available from component manufacturers. Fixed opening orifices are usually a nominal 3/64 inch (1 mm) size. The orifice diameter (or head differential) can be determined from a standard discharge equation given below.

$$Q_{o} = sQ_{a} = C_{d} \frac{\pi d^{2}}{4} \sqrt{2 \triangle p/p_{a}}$$

where $Q_{\mathbf{O}}$ = discharge from a single orifice

s = spacing of orifice

 $\mathbf{Q}_{\mathbf{a}}$ = air discharge per unit length

 C_d = orifice loss coefficient (from fluid mechanics texts, etc.)

d = orifice diameter

 Δp = pressure difference across orifice

 p_a = mass density of air

Usually an air bubbler system is a low pressure, high volume design. This suggests the use of rotary blowers instead of piston compressors used previously. A disadvantage of a piston compressor is the presence of finely atomized oil in the air diffuser lines which has passed by the pistons. This causes clogging of the lines and orifices. Noise is a disadvantage of the rotary compressor system.

Standby power or other alternative procedures should be available at a marina in the event that power is lost during a storm or for other reasons.

A well designed air bubbler ice suppression system will rapidly prove inadequate if little attention is paid to its pre-season cleaning and balancing. It is essential to begin an icing season with clean, balanced and operable system components.

Other Ways to Suppress or Weaken Ice

At this time it appears that compressed air ice suppression systems are best for the small craft harbor. A few other methods are given below and they may have some applicability.

Velocity or propeller systems drive a propeller that churns warm water to the surface. A lot of water is moved with only a few percent of the volume contributing to melting ice where suppression is desired. Methods to analyze melting ice with these systems have not been published. The velocity system's chief applicability appears to be in keeping limited areas open, e.g. around a boat left in the ice.

US CRREL personnel have observed a field experiment where $57^{\circ}F$ ($14^{\circ}C$) warm water was emitted from a diffuser 19 feet (6 m) below the ice surface. The warm water rose until it cooled to $39^{\circ}F$ ($4^{\circ}C$); at which time it descended as it was denser than the surrounding water. It did not reach the surface to melt the cover. Warm water diffusers might work with shallow water, higher pressures for momentum and insulated manifold and diffuser lines.

Insulating ice is very effective in reducing its thickness and retarding its growth. A few trials have been made in river ice. They effectively reduced a normal cover of 12 inches (30 cm) to zero. Problems involved are installing, maintaining and removing the insulation material.

Snow is a good insulator and snow thrown by a snow-blower is reworked and tends to set up. In this form, a snow cover might provide satisfactory insulation during severe periods of a winter.

Dark substances will absorb heat and cause melting. However coal dusting and other techniques are one-time temporary measures. Such procedures as well as chemical applications for melting are not acceptable today.

Electrical resistance heating is more expensive than other suppression methods but is feasible and perhaps applicable. Frazil ice, which is quite sticky can be electrically melted from intake structures. Lake ice could also be melted with electrical heat. As the area of contact between the structure to be protected and the ice diminishes, the more favorable resistance heating becomes. (Query, could pilings whose diameters were made small at the ice line, an area not usually requiring a large pile section modulus, be incorporated in a dock system design with electric heating to suppress the ice.)

finally, ice can be weakened and cracked by impacting, chopping, wave generation schemes and other mechanical means. These do not appear feasible for a small craft harbor however.

PART IV

BEARING CAPACITY OF ICE

This section of the guide contains information on the capacity of ice to support loads, especially construction loads resulting from using ice as a work platform for building small craft harbors. Winter construction has proven very successful in the Great Lakes. The ice easily handles men and small equipment with few accidents. Heavier loads have been moved and supported by ice also, including driving piling with light pile drivers and excavation dredging with regular construction cranes and draglines. Predicting safe allowable loads necessarily must be conservative. Ice strength must be assumed to be weaker than strengths which will be assigned in the next section for estimating the forces and uplifts due to ice.

This section begins with a discussion of ice parameters and bearing capacity theory. A figure is presented for maximum safe loads. The section concludes with some comments on thickening and strengthening ice.

Flexural Rigidity Length

Nevel (1976) introduced a term, flexural rigidity length, for a characteristic feature of all ice plate problems. The term in fact has been called characteristic length, or sometimes action radius. To derive the term, let us first consider the differential equation for a beam.

$$EI \frac{d^4w}{dx^4} = q$$

where w = deflection

x = distance

q = load

E = modulus of elasticity

I = moment of inertia

The term EI is a measure of how stiff the beam is. If now the beam rests on water the deflection increases the water pressure which reduces the load.

$$EI \frac{d^4w}{dx^4} = q - kw$$

where k = weight density of water

$$EI \frac{d^4w}{dx^4} + kw = q$$

$$\frac{EI}{k} \frac{d^4 w}{dx^4} + w = \frac{q}{k}$$

The term $\frac{EI}{k}$ on a unit width basis has the dimensions of the length to the fourth power, so the fourth root of $\frac{EI}{k}$ is a length, i.e.

$$\frac{EI}{k} = \frac{FL^{-2}L^3}{FL^{-3}} = L^4$$

where F and L represent force and length units

$$\frac{EI}{k} = \frac{(E) (1/12) (h^3)}{k} = \frac{Eh^3}{12k}$$

where h is the plate thickness

Now for a plate instead of a beam the term $(1 - v^2)$ is introduced through strain compatibility relationships, where v is Poisson's Ratio. The flexural rigidity length l for a plate is defined as follows:

$$1 = \left[\frac{Eh^3}{12 (1 - v^2)} \right]^{\frac{1}{4}}$$

The value of 1 is rather insensitive to E and v, and sensitive to h, the plate thickness. This means that we need not know E and v precisely, but do need to know the ice thickness to make good engineering estimates. The flexural rigidity lengths for various E's and h's for a Poisson's Ratio of 1/3 are listed in Table 9. The values will be used in solving bearing capacity and other plate problems.

Ice Thickness, h inches (cm)		<u>Modul</u>	us of Elasti kips/in ² (MN/m ²)	icity, E	
	500	750	1000	1250	1500
	<u>(3445)</u>	(5168)	<u>(6890)</u>	(8613)	(10335)
6	10.8	11.9	12.8	13.6	14.2
(15)	(3.3)	(3.6)	(3.9)	(4.1)	(4.3)
12	18.1	20.1	21.6	22.8	23.9
(30)	(5.5)	(6.1)	(6.6)	(6.9)	(7.3)
18	24.6	27.2	29.2	30.9	32.3
(46)	(7.5)	(8.3)	(8.9)	(9.4)	(9.8)
24	30.5	33.8	36.3	38.3	40.1
(61)	(9.3)	(10.3)	(11.1)	(11.7)	(12.2)
30	36.1	39.9	42.9	45.3	47.5
(76)	(11.0)	(12.2)	(13.1)	(13.8)	(14.5)
36	41.3	45.7		52.0	54.4
(91)	(12.6)	(13.9)		(15.8)	(16.6)
42	46.4	52.4	55.2	58.4	61.1
(107)	(14.1)	(16.0)	(16.8)	(17.8)	(18.6)

TABLE 9 FLEXURAL RIGIDITY (CHARACTERISTIC)
LENGTHS, 1, Feet (m)
(Poisson's Ratio = 1/3)

Bearing Capacity of Ice Sheets

Kerr (1976) has presented a critical survey of the literature on bearing capacity of floating ice plates subjected to static or quasi-static loads. The work contains discussions of general questions, analytical solutions, and field and laboratory results.

When using an ice sheet for supporting loads one question of obvious interest is what is the maximum load or the "breakthrough" load. Also the time at which this breakthrough occurs is of interest. Simple proven methods have not been developed to tell us the breakthrough load and when it will occur. It is a viscoelastic problem.

Field and laboratory observations indicate that indeed an ice sheet continues to carry a load after the first crack occurs. To solve analytically any plate problem requires limiting the problem, choosing an idealized model and formulating the proper equations from mechanics (statics and dynamics). Having formulated the problem, exact or approximate solutions are sought that agree with field and laboratory observations. However, if we limit the plate problem, for ease of analysis and conservatism, to an uncracked plate we then have a different problem after the first crack appears. More complex and different problem formulations are then necessary. We will however use the "first crack" analysis to determine our safe loads for construction on the ice.

If we place a uniform load on a circular area and increase the load, the ice will crack. There will be one or more radial cracks that extend out from the load. If we continue to add load a circumferential crack will form somewhat concentrically around the load but at a distance out from the load at the ends of the radial cracks. This forms wedges which usually number from four to eight or more. When this circumferential crack forms breakthrough can be expected shortly thereafter. On thick ice, even more load may be added due to side interaction between the wedges. However this possible support cannot be counted upon. In fact, it is reasonable to assume that when radial cracking begins the ice is sustaining a load that is no longer a safe load.

Nevel (1976) has formulated a mathematical creep model for ice which includes primary, secondary and tertiary creep. His equations show that at the load, the deflection increases with time while the stresses decrease, or relax. This means that the maximum tensile stress (which is usually the critical stress in an ice plate bearing problem) occurs at the moment the plate is loaded. A usual failure criterion is to limit the maximum tensile stress. Using this criterion means the sheet should fail at once or not at all as the stresses relax with increasing deflection. This however is contrary to observations on ice sheets under sustained loads. A possible explanation is that the tensile strength is somehow affected by the creep process.

Frankenstein (1966) describes the results of a number of large scale breakthrough field tests, for both concentrated loads and distributed loads placed on an Arctic lake cover. In the distributed load tests a

15-ft (4.6-m) diameter aluminum tank, with a height adjustable to 20 ft (6.1 m), was placed directly on the ice surface. Lake water was pumped into the tank at a more or less constant rate to load the ice. The concentrated load tests were conducted in the same manner as the distributed tests except that the tank was placed on a platform balanced on a 24 inch (61 cm) diameter wooden block.

For distributed load tests the first circumferential crack did not produce breakthrough. Additional cracking of the wedges occurred parallel to the first circumferential but closer to the load. Some side wedge interaction probably occurred before final breakthrough on the closer circumferential cracks, numbering one or two in addition to the first crack connecting the far ends of the radials.

The crack phenomenon for the concentrated tests differed greatly from that of the distributed tests. The distributed tests always yielded a circumferential crack but in some of the concentrated tests, no circumferential cracks were visible either during loading or after failure. The failure hole diameter for the concentrated tests was very close to the diameter of the bearing block.

Tables 10 and 11 present some of Frankenstein's load test data. Temperatures taken in the ice ranged between 30°F and 15°F (-1°C and -9°C). The load contact pressures were computed on a nominal 15-ft (4.6-m) tank diameter and a 2-ft (0.6-m) block diameter. The punching shears for the concentrated load tests data were computed using the nominal block diameter, the ice thickness, and the collapse load.

Ice Thickness inches (cm)	Maximum Loads pounds (kN)	Time of Failure minutes	Maximum Deflection feet (m)	Ratio of Max. Load and Square of Ice Thickness lbs/in ² (N/cm ²)	Load Contact Pressure 1bs/sf (kN/m ²)	Pressure per Ice Thickness 1bs/sf/in (kN/m ² /cm)
6.3	23955	11.7	1.06	604	136	22
(16.0)	(107)		(0.32)	(416)	(6.5)	(0.41)
8.2	28550	12.2	U.77	424	162	20
(20.8)	(127)		(0.23)	(292)	(7.8)	(0.38)
9.5	36980	14.6	0.87	410	209	22
(24.1)	(164)		(0.27)	(283)	(10.0)	(0.41)
12.8	44330	30.5	1.56	271	251	20
(32.5)	(197)		(0.48)	(187)	(12.0)	(0.38)
12.4	48303	21.6	0.98	314	273	22
(31.5)	(215)		(0.30)	(216)	(13.1)	(0.4 ₁)
15.6	77945	30.9	1.26	320	441	28
(39.6)	(347)		(0.38)	(221)	(21.1)	(0.53)

Table 10 Distributed Load Test Data (Frankenstein 1966)

Ice Thickness inches (cm)	Maximum Load pounds (kN)	Time of Failure minutes	Maximum Deflection feet (m)	Ratio of Max. Load and Square of Ice Thickness lbs/in ² (N/cm ²)	Load Contact Pressure 1bs/sf (kN/m ²)	Punching shear lbs/in ² (kN/m ²)
11.4	30,968	9.1	0.76	238	9,860	36
(29.0)	(138)		(0.23)	(164)	(472)	(248)
11.0	28,970	8.6	0.45	240	9,220	35
(27.9)	(129)		(0.14)	(165)	(441)	(241)
13.5	34,810	9.7	0.65	191	11,080	34
(34.3)	(155)		(0.20)	(132)	(531)	(234)
15.4	51,339	16.8	0.71	216	16,340	44
(39.1)	(228)		(0.22)	(149)	(782)	(303)
17.8	61,603	21.2	0.76	194	19 , 610	46
(45.2)	(274)		(0.23)	(134)	(939)	(317)

Table 11 Concentrated Load Test Data (Frankenstein 1966)

Frankenstein's distributed load tests indicate a fairly narrow range of collapse load contact pressures per unit thickness of ice, i.e. 20 to 28 lbs/sf/in thickness (0.38 to 0.53 kN/m²/cm thickness). The concentrated load tests also give a narrow range for punching shear, i.e. 34 to 46 lbs/sq in (236 to 317 kN/m²). The ratios of collapse loads to the squares of the ice thicknesses varied in the range of 320 to 604 lbs/in² (221 to 416 N/cm²) for distributed loads and 191 to 240 lbs/in² (132 to 165 N/cm²) for concentrated loads.

In the next section we will consider safe loadings for ice sheets that could be developed during wintertime construction of a small craft harbor. Quick intense loads causing breakthrough and long term viscoelastic loadings are not dealt with here. (Refer to Kerr, 1976 and Nevel, 1976.)

Safe Loads on Ice Sheets

Nevel (1977) has developed a pocket calculator program to predict safe loads based upon an elastic first crack analysis. His analysis and the program are included in Appendix III. Circular and rectangular loads or combinations of loads can be handled.

For our purposes we have solved his general equations for the simple case of one concentrated load. The computed maximum "safe load" can be used for estimating single construction loads on ice sheets; provided the ice is sound and free of discontinuities. The load cannot be placed near free edges of the sheet or in the vicinity of walls or shores where cracks behave as a free edge condition. To be out of the direct influence of these edge conditions the load should be away a distance equal to two or more flexural rigidity lengths. The expression for the maximum safe load P with Poisson's Ratio assumed to be 1/3 is:

$$P = 1.396 \ \nabla h^2 \frac{(a/1)}{\text{kei}!(a/1)}$$

where of is the allowable stress

h is the ice thickness

a is the radius of the load distribution

l is the flexural rigidity length

kei' is the first derivative of the Kelvin function kei

The expression for the maximum deflection \boldsymbol{w} which will occur directly under the load is

$$w = \frac{(P) [(a/1) \ker' (a/1) + 1]}{(a/1)2 k 12 \pi}$$

where ker' is the first derivative of the Kelvin function ker

k is the weight density of water

Kelvin functions and their derivatives are tabulated in mathematical handbooks. They are difficult to work with for small values of argument (a/l) due to the sensitivity of the functions' values with changes in (a/l). Included in Appendix IV are tables of Kelvin functions and derivatives from Nevel (1959) for arguments ranging between 0.000 and 0.500. Also the appended pocket calculator program has a Kelvin function subroutine that can be used for these or other arguments.

For small values of (a/1) the following approximation for the maximum deflection w is useful.

$$w = \frac{P}{8k1^2}$$

Figure 5 is a plot of the maximum safe load for ice with a flexural strength σ of 100 psi (689 kN/m²) and a modulus of elasticity E of 750 ksi (5168 MN/m²). The maximum safe load is linearly proportional to the flexural strength. For this type of ice loading problem, design flexural strengths would range between 50 and 100 psi (345 and 689 kN/m²) for fair to good ice. Loadings beyond this stress level are not recommended for usual construction procedures. Of course care is necessary in placing loads on any thin ice or placing very large loads regardless of the ice thickness or the analysis. Figure 5 simply gives recommended maximum values that should not be exceeded and that should be used with caution.

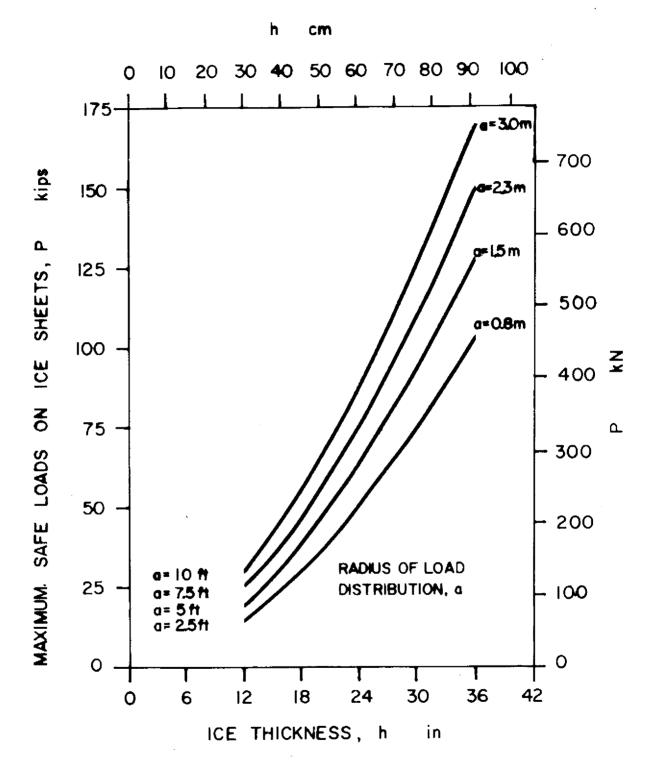


Figure 5. MAXIMUM SAFE LOADS ON ICE SHEETS (after Nevel, 1977)

To illustrate the use of Figure 5 and the equations let us assume we want to place a 20 kip (89 kN) load on a 5 ft by 8 ft (1.5 m by 2.4 m) area. We have 15 inches (38 cm) of sound ice. Through area ratios we will convert the rectangular area to an equivalent circular area.

$$a = \left[\frac{5 \text{ ft x 8 ft}}{11}\right]^{\frac{1}{2}} = 3.6 \text{ ft (1.1 m)}$$

From Figure 5 at 15 inches (38 cm) and 3.6 feet (1.1 m) we find the maximum safe load to be 23 kips (102 kN) which is greater than the specified 20 kips (89 kN).

If we have weaker ice, because of warm temperatures or previous damage from construction we would want to reduce the load. The flexural strength may be only 50 psi (345 kN/m^2) and for this case the safe load would be halved. In this example we should also consider the contact pressure and the deflection.

contact pressure =
$$\frac{20,000 \text{ lbs}}{5 \text{ ft x 8 ft}}$$
 = 500 psf (23.9 kN/m²)

This pressure is high for this relatively small loading area and it would be well to use a larger area for the 20 kip (89 kN) load. But assuming the given loading condition we can compute the deflection of the ice sheet from the formula

$$w = \frac{(P) \left[(a/1) \text{ker'} (a/1) + 1 \right]}{(a/1)^2 \text{k} 1^2 \text{TT}}$$
where $a/1 = 3.6/27.0 = 0.1333$
 $(1 \text{ interpolated from Table 9})$

$$k = 62.4 \text{ lbs/ft}^3 (1000 \text{ kg/m}^3)$$

$$v = \frac{(20,000 \text{ lbs}) \left[(0.1333) \left(-7.4501 \right) + 1 \right]}{(0.133)^2 (62.4 \text{ lbs ft}^{-3}) (27.0^2 \text{ ft}^2) \text{TT}}$$

$$= .054 \text{ ft or } 0.65 \text{ in } (1.6 \text{ cm})$$
or
$$v = \frac{P}{8 \text{k} 1^2} = \frac{(20,000 \text{ lbs})}{(8) (62.4 \text{ lbs ft}^3) (27.0^2 \text{ ft}^2)} = 0.55 \text{ ft } (1.7 \text{ cm}) \text{ approx.}$$

One criterion for allowable deflection is that the deflection be less than 0.08 times the ice thickness, i.e. the ice sheet will not be submerged. For our 15 inch (38 cm) sheet the allowable deflection criterion is equal to 1.2 inches (3 cm).

If the freeboard of the ice is exceeded water can flood the surface through cracks and other openings. This warms the ice and weakens it and also adds further loads. Loadings on ice less than 18 inches (46 cm) in thickness should be checked for possible submergence. Ice thicker than this will not submerge under the recommended maximum loads.

Figure 5 should be of assistance in estimating construction loads on small craft harbor ice covers. Special precautions are necessary however if moving loads across frozen bodies of water is planned. Moving loads create dynamic effects in the ice cover and water. Unfortunately the natural critical velocity is easily achieved on many lakes. At this velocity, the deflections and stresses are amplified. It is necessary to go at a slow speed or very fast. The technical literature should be reviewed by those planning moving loads across frozen lakes.

Thickening and Strengthening Ice Covers

In using ice for bearing capacity, actual thicknesses should be measured. Dates when a natural ice sheet will be thick enough to support a specified load are of interest and can be estimated from past observations and rules of thumb. One rule is that the thickness is directly proportional to the square root of the Freezing Index (previously defined in Part I). Using ∞ as a proportionality factor we can write the following rule of thumb expression:

 $h = \propto \sqrt{FI}$

where h is ice thickness in inches

FI is Freezing Index, Fahrenheit

Assume we have had 45 days of weather where it has generally been below freezing. The Freezing Index would be calculated by summing algebraically the plus and minus departures of the daily mean temperatures from the freezing temperature. The Freezing Index in our example will be assumed to be 750 Fahrenheit degree-days and our locality factor equal to 0.8.

 $h = 0.8 \sqrt{750}$

= 22 inches (56 cm)

The locality factor ∞ can be historically established by several years measurements of ambient temperatures and ice thicknesses. It can be used for estimating when sufficient thicknesses of ice might naturally become available. Ice can also be intentionally thickened and strengthened to aid nature.

Hoffman (1967) describes surface flooding techniques for improving natural ice areas. These techniques have been developed by the U.S. Naval Civil Engineering Laboratory. Among their recommendations are the following items.

The maximum mean daily temperature at which ice flooding can be performed satisfactorily is considered to be about $+15^{\circ}F$ ($-9^{\circ}C$). At temperatures much above this point, freezing rates are slow and long periods are required between applications.

The depth of water applied to any point should not be greater than that which will freeze through in 24 hours. At temperatures from $0^{\circ}F$ to $-10^{\circ}F$ ($-18^{\circ}C$ to $-23^{\circ}C$) with little wind, this is about 4 inches ($10^{\circ}C$ cm) of water.

A cooling period equal to the freezing period should be allowed before an ice area is reflooded. Such cooling is necessary for restoration of ice temperature and recovery of ice strength and resistance to creep.

Additional water should not be applied until all areas of the previous flood have frozen solid. The premature reflooding of an unfrozen area is very undesirable since the freezing time increases exponentially with depth.

Air bubbles which form in the flooded ice surface should be broken before reflooding.

Ouff (1958) describes in general terms the thinking of logging operators on the strengthening of ice sheets. Ice landings are constructed and main-tained by flooding on top of existing ice and rolling any snow that falls.

Where snow is lying on an ice sheet it is, if at all possible, compacted by rolling before flooding. Slush always seems to form after a snowfall. The weight of the snow cover depresses the ice sheet and free water comes up through cracks or the insulating blanket of snow allows a current, probably a thermal in lakes, to weaken the ice sheet from below. Rolling immediately after a snowfall, or during a storm if it is a severe one, improves frost penetration and lessens the slush problem.

Rose and Silversides (1958) describe the merits of surface flooding for increasing the speed at which ice thickens. Starting with a layer of ice 12 inches (30 cm) thick and -20°F (-29°C) weather, if 5 inches (13 cm) of water are added on the surface it will take approximately 15 hours for it to freeze and give a total thickness of 17 inches (43 cm) of ice. If the ice surface is bared and ice is to be formed on the under-surface of the 12 inch (30 cm) ice sheet, it will take 60 hours to add an additional 5 inches (13 cm). This is 4 times as long.

Obstrom and DenHartog (1976) have reported on a series of cantilever beam tests designed to determine the efficacy of adding reinforcement to an ice cover. Tests were run using 1 inch (2.5 cm) diameter tree branches, 3/16 inch (5 mm) diameter wire rope and 9/16 inch (14 mm) half-round wood dowels as reinforcement. A definite advantage was noted from using reinforcement, even when poorly placed. The reinforced ice carries a load

even after it cracks. Thus, after the initial cracks there is time to remove people and equipment before final breakthrough. Disadvantages to the reinforcement are that the darker types absorb radiation and thereby cause weakening of the ice cover. Also, in many cases, the time and effort required to place reinforcement may exceed those required to achieve equal strength by additional thickening of the ice sheet.

The reinforcing process consists of laying reinforcement material on the ice, then flooding the area and allowing the reinforcement to freeze into the ice. For optimum strength, the reinforcement should be added to the side of the ice that carries the tensile forces. Initial cracking of the ice sheet is caused by tensile stresses near the bottom surface of the ice, but final breakthrough is caused by tensile stresses near the top surface.

PART V

VERTICAL ICE FORCES AND OTHER FORCES

One of the most significant and damaging ice forces results from changes in water levels. These changes cause the ice sheet to move up and down tearing and pulling structures built in small craft harbors.

In Part IV we considered the capacity of an ice sheet to support construction loads. Necessarily we used a conservative value for the strength of the ice. In estimating the uplift forces from ice we will assume stronger ice as it will probably exist sometime during the winter. Computing uplift loads is complicated because we do not know which of several failure criteria best represents the uplift force. We have selected a "first crack" analysis to predict the minimum uplift for which we should design. The maximum uplift force for design still remains a matter of engineering judgment and site circumstances.

This part begins with a discussion of lake level changes. Ranges of uplift loads are then presented together with some ways to attenuate these loads. The section concludes with some comments on buckling of ice sheets and other loads.

Water Level Fluctuations

Water levels, and hence ice levels, in the Great Lakes vary seasonally. The small craft harbor designer must anticipate and estimate these seasonal variations when planning fixed harbor structures, such as docks and piers, launching ramps and bulkheads. To a lesser extent, water level variations affect floating structures. They are of concern, however, e.g. in connection with anchorage and gangways.

To illustrate the importance of water levels on design, consider a fixed height dock specified to be a desired freeboard above normal summer water levels. Should levels prove otherwise, due to larger natural variations or inaccurate estimates, such a dock can end up well above the water or too close to the water. It becomes barely usable by the intended boats. Also, during wind storms over water, severe water level changes can occur. At Buffalo, New York the lake has risen to as much as 8 feet (2.4 m) above the normal water level.

Problems can result with winter ice levels also. Individual pilings supporting a fixed dock may be protected (deiced) with a compressed air bubbler system. Normally the number of pilings would be kept to a minimum to reduce the amount of air required. By doing this, dock members tend to have greater depths because of their longer spans. Should the ice become higher than estimated, the lower elements of the dock members (which are not deiced) may come in contact with the ice sheet or be shoved by broken ice pieces that are frequently found on top of an ice cover.

Local experience can aid the designer in estimating water and ice levels. If severe water level changes can be anticipated, a floating or adjustable height design would be called for. The following publications will also help the designer.

Monthly Bulletin of Lake Levels for the Great Lakes, Department of the Army, Detroit District, Corps of Engineers, P. O. Box 1027, Detroit, Michigan 48231. (Bulletin lists maximum, minimum, and expected levels.)

Hydrograph of Monthly Mean Levels of the Great Lakes, U. S. Department of Commerce, NOAA, National Ocean Survey, C3314, 6001 Executive Boulevard, Rockville, Maryland 20852. (Lists monthly mean lake levels from 1860 to date.)

In addition to seasonal lake level variations, there is an unusual water level phenomenon known as a seiche (pronounced sash). It is a standing wave oscillation of an enclosed or semi-enclosed water body that continues, pendulum fashion, after the cessation of the originating force. It is a short-term rise and fall of the water level and is caused by either persistent, strong winds piling up the water at one end of a basin, or changes in barometric pressure over the lake, and sometimes a combination of both. Although seiches are unnoticeable on small lakes, they are quite noticeable on the Great Lakes and in bays and harbors. The period of a seiche is a few minutes in a bay or harbor and about ten hours for a Great Lake.

Hodek and Doud (1975) measured an almost constant fluctuation of the winter water level in the Ontonagon, Michigan harbor. The amplitude had an observed maximum of 0.8 feet (0.24 m). The major period varied from 5 minutes to more than 10 hours. Higher frequency water oscillations also were observed and a change in the water level of 3 inches (8 cm) in 10 minutes was observed.

Seiches are easily observed in a marina in the summertime by noting a drop in the water surface with respect to some fixed object like a sloping launching ramp or a piling. In fact, severe seiches have been observed where the water level fluctuated three feet (1 m) and more in a few hours.

As a result of seiche action, an ice plate in a harbor moves up and down exerting large forces on structures penetrating the ice sheet. In the next section we will estimate the minimum level of these forces.

Minimum Ice Sheet Uplift Loads

In a manner similar to that used in computing safe loads on ice sheets, we can estimate minimum uplift loads to be expected. The analysis is an elastic one and is compatible with the rapidity with which water levels change under seiche action. The analysis is also a thin plate analysis which is not strictly correct for ice sheets thicker than the embedded pile diameter. For this case the stresses computed are too high in the vicinity of the load. However this is acceptable because it is conservative.

We will assume a round pile frozen into (rather than setting upon) an ice plate and determine the force when the plate first cracks. This will be the minimum force we can expect when a pile or structure is frozen into the ice.

The differential equation formulating this problem has been solved. We have evaluated the solution for the following boundary conditions. At a distance from the center of the concentrated load equal to the radius of the load (distance a), the slope of the deflected ice plate is zero; and at distance a the total load is carried by the stress on the failure surface which is the thickness of the ice plate times 2 Tra. The crack that forms at failure is a circumferential crack located a distance a out from the center of the load. This crack is consistent with what is seen in the field when an ice sheet pulls upward on a strong well-embedded piling and eventually cracks. The distance a is not the radius of the piling, but is larger. This is due to the failure occurring by cracking in the ice rather than slippage at the ice-pile interface, and because immediately adjacent to the pile the ice is thicker than the rest of the sheet. (An ice collar forms in this area.) This thickening phenomenon results from heat transfer through metal piling (and to some extent through wood piling) increasing the ice thickness in this area. Also as the winter goes on and the crack has formed and refrozen many times, the piling experiences a kind of dipping action in the cold water, (somewhat like dipping a candle wick in wax to form a candle).

The equation for the minimum ice sheet uplift load, based upon this first circumferential crack is as follows:

$$P = \frac{(1/3) \, \text{Th}^2 \, (a/1)}{\left[\frac{\text{kei } (a/1) \, \text{kei'} \, (a/1) + \text{ker } (a/1) \, \text{ker'} \, (a/1)}{\text{kei'} \, (a/1) \, \text{kei'} \, (a/1) + \text{ker'} \, (a/1) \, \text{ker'} \, (a/1)}\right]}$$

where T is the allowable stress

h is the ice thickness

a is the radius of the load distribution

l is the flexural rigidity length

kei and ker are Kelvin functions

kei' and ker' are first derivatives

Figure 6 is a plot of the minimum ice sheet uplift load for ice with a flexural strength σ of 200 psi (1378 kN/m²) and a modulus of elasticity E of 750 ksi (5168 MN/m²). The minimum uplift load is linearly proportioned to the flexural strength, which could be greater than the assumed value of 200 psi (1378 kN/m²). Also, as we will show later, making different assumptions for the failure criterion will result in larger loads.

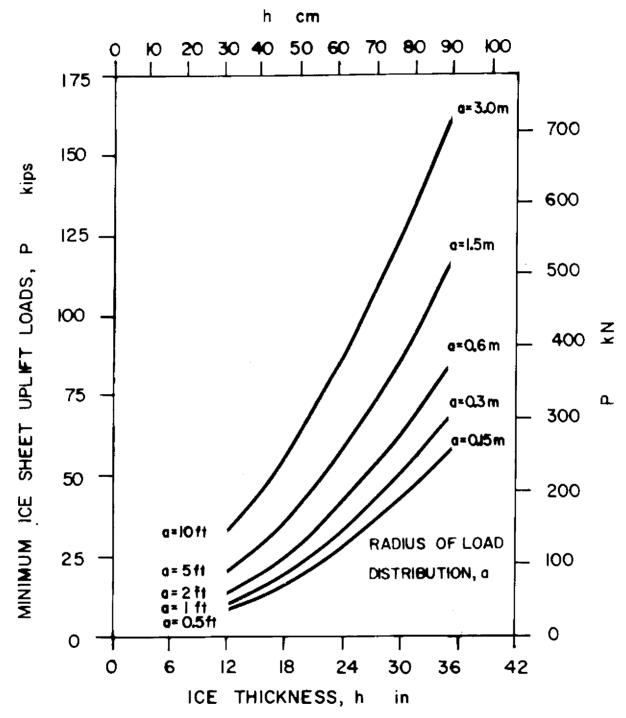


Figure 6. MINIMUM ICE SHEET UPLIFT LOADS

To illustrate the use of Figure 6, let us compute the minimum uplift for a 12 inch (30 cm) round steel pile in 24 inches (61 cm) of ice. As stated above the radius a will be greater than one half pile diameter. For the usual steel pile shapes, the radius of the ice collar will be about six inches (15 cm) greater than the pile radius. (For wood, a value of about half this much has been observed.) From Figure 6 at h = 24 inches (61 cm) and a = 1 foot (0.3 m), we find the minimum uplift load is 33 kips (147 kN). This estimated load is the least that we should design for with a 12 inch (30 cm) pile in 24 inches (61 cm) of sound cold ice (σ assumed = 200 psi, 1378 kN/m²). We do not know precisely what the maximum load could be, but it could be several times as much. The ice does not fail on the ice-pile interface, but fails in the strength of the ice sheet itself.

The uplift on a square construction could be reasonably approximated by an inscribed circle of radius a equal to $\frac{1}{2}$ of the length of the side. This would be conservative since the angles of a square would tend to cause stress-concentrations in the ice.

Nevel (1972) has formulated a more severe failure criterion for the ice pile uplift problem. It is a criterion which gives a near, if not the maximum, upper bound to the problem. It is assumed that radial cracks form and then a circumferential crack forms connecting the ends of the radials. This cracking pattern forms a series of truncated wedges whose tips are supporting the load and whose bases are failure planes when the circumferential crack develops. If we neglect any side interaction between the wedges, and assume that six wedges have formed the Nevel (1972) solution for the maximum load becomes:

P = 1.154
$$\sigma$$
 h² [1.05 + 2.00 (a/1) + 0.50 (a/1)³]

where the bracketed term is an approximation to an exact solution.

For our example we can compute the maximum load P as follows:

```
a = 1.0 feet (30 cm)

1 = 33.8 ft (10.3 m) at E = 750 ksi (5168 kN/m<sup>2</sup>)

a/1 = .0296

[ ] = 1.109

P = (1.154) (200 lbs in<sup>-2</sup>) (24<sup>2</sup> in<sup>2</sup>) (1.109)

P = 147,400 lbs or 147 kips (654 kN)
```

This truncated wedges load estimate is $4\frac{1}{2}$ times the estimate based upon the first crack analysis. For smaller values of a the loads are 6 times as great and for larger values of a. 3 times as great.

The proper values to use for design are matters of engineering judgment and experience. The values should lie within the ranges presented and an uplift of not less than the minimums indicated on Figure 6 should be used for design with appropriate safety factors. If for example a minimum load is estimated to be 25 kips (lll kN), then the ultimate resistance of the soil foundation system preventing uplift, or the dead weight of the structure, should be greater than 25 kips (lll kN) by perhaps a factor of 2 or 3. In selecting a design criteria for a particular small craft harbor, other factors will enter in; such as whether some damage will be tolerable, the ease with which pullout resistance can be developed, or the location of the pile and its relationships with other piles or dock members.

Ice can also exert down-drag loads when the water level drops and leaves the ice hanging on the dock piles. The ice can easily span across between piles and develop the full dead weight of the ice sheet. Large drops in levels have occurred and it is believed that the ice sheet is capable of losing all its buoyancy. The down load situation is aggravated by a tendency to have thicker ice under shady dock areas than in the more exposed aisle or fairway areas.

There is limited field data on forces on full size piles but several direct observations have been made. Hodek and Doud (1975) and Hodek (1976) have measured vertical loads of 18 kips (80 kN) compression and $11\frac{1}{2}$ kips (51 kN) tension or uplift on a 15 inch (38 cm) instrumented sleeve placed around a dock pile. The thickness of the ice was about 16 inches (41 cm). The loads were those actually felt by force transducers on the instrumented pile sleeve.

Muschell (1976) jacked against an ice cover to pull frozen-in piles. Loads up to 25 kips (lll kN) were recorded. The loading rate was 0.5 psi per second (3.4 kN/m 2 /s). Ice adhesion to steel was computed by dividing the pullout load by the actual ice thickness and area of ice contact to steel on HP 8 x 36 and HP 10 x 42 piles. The ice adhesion values ranged between 12 and 24 psi (83 and 165 kN/m 2).

In January 1977, Muschell (unpublished) recorded a value of 68 psi (469 kN/m^2) for ice adhesion to an 8-inch (20-cm) pipe pile in a comparable field test.

The lifting force per unit of length of circumference of a structure frozen into an ice sheet becomes less as the size of the structure increases. The lifting force per unit of length will approach the values for long straight walls. Lofquist (1951) and Michel (1970) have estimated the vertical forces exerted by ice rigidly attached to walls. The estimates are based upon an elastic analysis and rapid rises in water levels.

$$p = \frac{\sqrt{2} \sqrt{h^2}}{6 \cdot 1}$$

$$\triangle = \frac{p}{\sqrt{2} + 1}$$

where p is the lifting force per unit length

△ is the rise in water level

Assume we have 24 inches (61 cm) of ice with a strength of 200 psi (1378 kN/m²) and a modulus E = 750 ksi (5168 MN/m²). The uplift force per unit length and the water level rise are computed as follows:

$$p = \frac{(\sqrt{2}) (200 \text{ lb in}^{-2}) (242 \text{ in}^2)}{(6) (33.8 \text{ ft})}$$

= 803 lbs/ft (11.7 kN/m)

$$\Delta = \frac{(803 \text{ lbs ft}^{-1})}{(\sqrt{2}) (62.4 \text{ lbs ft}^{-3}) (33.8 \text{ ft})}$$

= .27 ft or 3½ inches (8½ cm)

Assuming a different strength value for the ice will change the uplift load in a linear manner. However the loads computed at \mathcal{T} equal 200 psi (1378 kN/m²) are believed to be adequate since the method of analysis is conservative. There normally is cracking parallel to a wall or long crib structure and this reduces the uplift from that assumed for rigid attachment. Walls and long cribs seem to experience little damage due to ice uplift. Occasionally, however, the top of a crib will be pulled off because it is inadequately secured to the rest of the crib.

The next section presents some suggestions for reducing and eliminating uplift forces.

Reducing Ice Sheet Uplift Loads

In Part III we discussed ways to suppress ice and reduce its thickness. This is the most positive way at this time to reduce ice uplift on small craft harbor structures. A number of other techniques have been tried and are being tested now. The projections are optimistic that reliable and economical ways will be developed to reduce ice uplift forces.

Ice adheres tenaciously to most construction materials. There is no comprehensive treatise on ice adhesion; only technical and research papers.

Representative adhesion values for ice to various materials and coatings determined by Freiberger and Lacks (1961) are presented in Table 12. The results are based on experiments for artificial ice loaded at the rate of 5 psi/s ($34 \text{ kN/m}^2/\text{s}$), and depend strongly on the rate.

Ice	Adhesion	Range	of	Values
-----	----------	-------	----	--------

	psi	$\frac{kN/m^2}{}$
Metals	85-120	586-827
Woods	45-80	310-551
Plastics	25-40	172-276
Glass	20-150	138-1034
Paints	80-100	551-689
Resin Films	30-130	207-896

Table 12 Ice Adhesion (after Freiberger and Lacks, 1961)

Jellinek (1957) performed shear experiments with snow-ice sandwiched between two stainless steel plates. The results of his work on small samples indicated the adhesive strength to be independent of thickness of the ice layer and cross-sectional area. He found the adhesive strength to be a linear function of the temperature until it becomes larger than the cohesive strength of ice at about $9^{\circ}F$ ($-13^{\circ}C$), where a very sharp transition from adhesive to cohesive breaks takes place. Representative values are as follows: $28^{\circ}F$ ($-2^{\circ}C$)--30.4 psi (209 kN/m²), $7^{\circ}F$ ($-14^{\circ}C$)--236 psi (1626 kN/m²), and $-4^{\circ}F$ ($-20^{\circ}C$)--233 psi (1605 kN/m²).

Michel (1970) has observed that there is very little available data on strength of adhesion of ice to structures and, there is a lot of scatter in the experimental results because of insufficient control of the uniformity of the structure of the ice and other relevant factors. Michel notes that

in most cases the strength of adhesion of ice to construction materials is as high as the shear strength of the river or lake ice itself, which is usually taken to be from 80 to 150 psi (551 to 1034 kN/m²).

And as previously stated, Muschell (1976 and unpublished) found values ranging up to 68 psi (469 kN/m^2).

For our purposes we can say that slippage between ice and the piling material will probably not naturally occur at loads less than the calculated minimums. Therefore in order to reduce uplift effects we must cause the ice to slip at loads less than these minimums.

Field and laboratory experiments have shown that wrapping a pile with various smooth materials doesn't work very well because the materials become unattached after a number of cycles of uplift. Materials that have been tried include tetrafluorethylene polymer wrap (teflon-like) and thin polyethylene sheets. Both materials were torn loose and worked their way up and out of the ice. A three ply wrap with polyethylene sheets was unsuccessful too. The three ply system was tested to see whether the outer layers of polyethylene, which would be assumed stuck to the ice and to the pile, would slide on the inner sheet sandwiched between the two. All layers worked their way out of the ice.

Laboratory studies and field tests are in progress now (1977) to evaluate protective jackets and coatings that could be installed or applied on existing piles or specified for new piles. US CRREL is undertaking this work and reports on their results are expected in 1978. One system of jacketing uses existing fiberglass products and methods for repairing damaged underwater piling. Preliminary laboratory test results show about a 4 fold reduction in the maximum force recorded on a test pile frozen into an ice sheet and displaced up and down in a simulated cyclic seiche action. Work is also underway to evaluate navigation lock wall epoxy and co-polymer coating systems described by Frankenstein et al (1976). Large reductions in adhesion values have been demonstrated but work on coating life and methods of application are needed.

Ice uplift loads can be reduced by causing the ice to crack so that large loads are not transmitted to structures. Group action is frequently effective in accomplishing this cracking. For example, several piles lashed together to act as a single dolphin, are effective in resisting uplift. A series of piles spaced 5 feet (1.5 m) on centers under a pier supporting a boat hoist machine have resisted uplift because of the group action. Also a T-head pier with pairs of pilings spaced about 25 feet (8 m) on centers has been found to act as a unit. A crack forms following the outline of the pier and completely encircling it. This crack demonstrates the group action of the entire piling system.

The piles toward the center of a marina and the piles at the ends of finger piers are more prone to pulling than those closer to shore. This is apparently due to the plate action and boundary conditions on the plate subjected to hydrostatic uplift. The ice sheet may also be frozen fast to the shoreline. This phenomenon suggests designing the outer piles for large

uplift forces and thereby cause the ice to fail about them. If this happens the piles closer in should feel little if any uplift force.

Another system that has been successfully used to eliminate uplift forces on marina pilings is the sleeved pile system (U.S. Patent No. 3.543.523). It allows vertical motion of a pile. It consists of a round hollow steel pile driven (with a follower) so that the top of the pile is 2 feet (0.6 m) below the bottom surface of the ice. A slightly larger diameter hollow round steel pile piece (the sleeve) fits over the top of the driven pile. The driven pile has a cap plate, and welded inside the sleeve about 4 ft (1.2 m) from its lower end is a round bearing plate. When the sleeve is dropped over the pile, the round plate bears on the cap plate. On top of the sleeve is another cap plate for bearing of dock beams. As ice freezes to the sleeve, the sleeve moves up and down over the lower, stationary pile (analogous to a stationary piston in a moving cylinder). The sleeve works well, allowing vertical motion. However horizontal ice forces may still be acting on the piles. Also dockage systems using sleeved piles are more complicated and costly to build because differential vertical motions must be provided for in the dock members and supported utilities that may be on the docks.

Other Ice Forces and Actions

In Part II we reviewed thermally induced thrusts, and in this part have considered uplift. Ice exerts other forces. For example, ice is very abrasive and causes damage to structures, particularly wood structures whose life in ice may be only twenty years, Striegl (1952). Cladding wood structures is necessary to achieve longer service lives. Unprotected concrete is subject to abrasion and freeze-thaw spalling.

Sheltered small craft harbors can experience damage from chunks of ice being driven into the harbor and against the docks during Spring breakup, and during periods when ice goes out of a harbor during Winter storms.

In addition to thermally induced horizontal forces, wind and currents can produce dynamic forces. Specific design recommendations are not now available but in general, marina dockage should be located in protected areas away from dynamic ice forces. However, those wishing information on dynamic effects should refer to the Neill (1976) comprehensive review and assessment of dynamic ice forces on piers and piles. Analytical approaches, design formulas, and full-scale and small-scale laboratory investigations are considered.

Ice under horizontal loadings may buckle before it crushes. Sodhi and Hamza (1977), using the finite element method, have performed a buckling analysis of a semi-infinite plate on an elastic foundation. The analysis determines the effective pressure which will cause the plate to buckle. If the crushing strength of ice is greater than this effective pressure, the plate will buckle; if less, a crushing failure will occur before buckling. Nevel (unpublished) has developed a useful expression based upon the Sodhi and Hamza (1977) finite element analysis.

$$P = k 1^{3} \left[b/1 + \frac{3.32}{1.0 + 0.25 b/1} \right]$$

where P is the buckling load

b is the structure width

l is the flexural rigidity length

k is the weight density of water

To illustrate, let us assume we have 6 inches (15 cm) of ice pushing against a wall 100 feet (30 m) long. The flexural rigidly length 1, assuming E equal 750 ksi (5168 MN/m^2), is 11.9 feet (3.6 m). The ratio b/l is 100 feet/11.9 feet or 8.40.

P =
$$(62.4 \text{ lbs ft}^{-3})$$
 (11.9^3 ft^3) $\left[8.40 + \frac{3.32}{1.0 + 0.25 (8.40)} \right]$
= 996,000 lbs or 996 kips (4430 kN)

We next compute the effective pressure as follows:

Effective Pressure =
$$\frac{P}{b h}$$

= $\frac{996,000 \text{ lbs}}{(100 \text{ ft}) (0.5 \text{ ft}) (144 \text{ in}^2 \text{ ft}^{-2})}$
= 138 lbs in⁻² (951 kN/m²)

If the crushing strength of ice is greater than this pressure, buckling will occur. The maximum crushing strength may be on the order of 300 psi (2067 kN/m^2) . In general, for small structures the ice crushes before buckling, except in the presence of thick ice where crushing always takes place.

There is little or no experimental data, expecially from full scale structures, to assist in the design of piles for horizontal focus. The American Association of State Highway and Transportation Officials (AASHTO) recommends that highway bridges be designed with the following provisions:

All piers and other portions of structures which are subject to the force of flowing water, floating ice, or drift shall be designed to resist the maximum stresses induced thereby.

The pressure of ice on piers shall be calculated at 400 psi (2756 kN/m²). The thickness of ice and height at which it applies shall be determined by investigation at the site of the structure.

The shortcoming of the AASHTO provision for 400 psi (2756 kN/m^2) ice pressure is that it is representative of an impact loading for strong ice failing in compression. The effects of pile shape and orientation, e.g. sloping piles where ice would fail in bending, are not considered. The maximum measured effective ice pressure on a bridge pier to date has been less than 200 psi (1378 kN/m^2) .

In a marina where Spring ice movements occur the effective ice pressure is probably less than 200 psi (1378 kN/m 2). Spring ice in a moving field has very low internal strength.

There is a tendency of ice to squeeze objects placed in it. How this happens is not understood and no methods of analyses are known. When a pontoon is placed in water (ice) it will either be squeezed up and finally rest on the ice, or be squeezed down and drawn further into the ice. The result of this action may be damage to the dock framing members connecting the pontoon units together; or sometimes rupturing of the pontoon material itself.

One hypothesis of the squeezing action is that it is the result of refreezing water. A small amount of ice adjacent to the pontoon melts during the day's warming sun. At night an ice cover forms over the melt water. This encloses the melt water which subsequently freezes in a partially confined space. This refreezing exerts a squeezing force which either ruptures or moves the pontoon.

A field study to investigate the action of pontoon units frozen in lake ice is scheduled for the winter of 1977–1978.

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APPENDIX 11--DEFORMATION AND STRENCTH OF ICE, LAVROV

This appendix contains an abbreviated summary of a monograph prepared by Lavrov (1969). It is used on experimental studies on structurally simulated ice vhose small-scale crystal structure is similar to that of an ice cover. This very important requirement is not fuffilled in small smalles but from a natural ice cover. The laws of deformation and strength presented are for farces acting on ice for chart periods, i.e. from a few seconds to a few minutus. Problems with very long-term strength and creep of ice are not dealt with.

Three peculiarities distinguish ice from other materials. First, it destats at a teaperature within a few degrees (a difference of tens of Celaius degrees being rare) of its melting point. Its the therefore usually contains some liquid water between crystals. The therefore usually contains fing the sliding resistance between erystals. Secondly, ice does not reactive, succembing inpurities and forms no celld solutions with them. The impurities are meanwhalt and form many pares, cavities, etc. in the ice. The third parallely is the comperstively large size of ice crystals. Their transverse discussions very between about a tenth of an inch (a few ma) and an inch (several cm). In metals these discussions come out to small tractions of a few handrallus of an inch (a mm). These peculiarities cusse ice to be a material which in its tobailly leacts to static loads.

The Matio of Transverse and Londitudinal Strains (Poisson's Matio)

Poisson's ratio plays as important role in the laws governing the behavior of ice under load. Poisson's ratio for ice is usually determined dynamically from the difference between the propagation velocities of the transverse and longitudinal olastic avees. Its average value at temperatures below 32°F (0°C) is 0.53. Experimental values range from 0 to more than 2.5. Poisson's ratio depends upon stress, in some cases exceeds the theoretically possible maximum, and is different for tension and for compression. The results obtained can be explained only on the besis of the actual behavior of the ice under loud, i.e., by allowing for its plasticity.

Schavior of Ice Under Load

Whatever the type of static load acting on ice (compression, tension, bending, shear) slip plays a decisive role in the overall deformation pattern. Slip in ice begins simultaneously with the load application. It is most intensive at the beginning, and then decreases according to a power law. The Widely held view that at the beginning the deformation is purely elactic, and that the above-mentioned processes take place only after some time, is erronmous. The clastic limit of ice is small, less than 7 psi (50kN/m²).

The fraction of the clastic deformation in the total deformation at the beginning of a short-time load action is less than later, i.e., ice is not a linearly clastic material and does not obey Hocke's law. This basic peculiarity of ice distinguishes it from other materials, in particular metals.

The linearity of some parts of the deformation diagrams for compression and bending do not indicate that ice obeys Hooke's law in its usual meaning, since Young's modulus E in the expression "stress equals E times strain" meat be replaced by another proportionality factor. This is the strain modulus which is not constant but depends on the duration of the load application.

The failure mechanism of ice varies with the test conditions (temperature, duration of load action, and direction of force), and is in the final analysis determined by the extent to which slip takes place both inside and between the ergetals.

None of the existing models simulating the mechanical properties of a body is fully applicable to ice, because one of its constituent elements, the "dashpot" (considered as plunger moving in a liquid according to Howton's law of viscosity) must be replaced by a more complicated structure, by scans of which not only the permanent set but also the chastic lag is allowed for. Studies carried out during recent years have shown that ice does not obey the law of viscosity since the relationship between the stress and the strain rate is not linear for it. It is a present impossible to express the stress-strain relationship analytically by allowing for all the laws entitioned above, which govern the behavior of ice under load.

The improndence of the Properties of Ice on the Conditions of its Formation and Grow, h

It is well known that the quality of ice (including the mechanical properties of the crystals) depends on the conditions of its growth. The mechanical strength of the crystals vary directly with the time taken for the ice to grow. A number of defects in the crystal structure is small when the locatives shouly, since the build-up of new lattice layers proceeds in small staps; three new layers are connected more regularly to the already finitived planes. Purdictions, injurities are oble to diffuse from the growing faces into the solution.

Other conditions being equal, the upper layers grow faster than the lower ones, so that the quality of the upper layers is worse.

of an ice cover formed at low temperatures is less than that of an ice cover of equal thickness, formed during milder weather. The quality of ice depends on the rate of its growth, so the strength

crystallization from the lower boundary of the ice, so that its growth rate snow on ice retards the removal of the latent heat of The thickness and strength of natural ice cover there ore The presence of decreases.

The Modulus of Elasticity and the Shear Modulus

instatons of the crystal axes; the crystals are separated by intercrystalline layers containing impurities and liquid phase. In addition, air bubbles are distributed ever the bulk of the ice. The moduli of elacticity for the ice cover are therefore smaller than the elastic constants of monocrystals, and Matural icc, particularly when its crystals are small, has various orvary with the atructure of the ice.

The elastic properties of ice are most accurately determined by means of dynamical methods, i.e., by measuring the propagation velocity of elastic (secund) waves in ice. Ultrusonic waves and the method of resonance vibrations yield the local values of the moduli of elasticity, while selemic nethods yield averaged values for an area,

Ceneral shortcoming of all dynamical methods of determining the elastic properties is the approximateness of the theory, which results from considering the ice as an isotropic body,

deformining the roduli of closticity, it appears that the latter differ in-skinthicantly for fresh-water ice of different origin and from different Allowing for the peculiarities of the various dynamical methods of

Young's modulus E and the shear modulus G vary with temperature and method by which determined. Approximate average values are 1200 ksi (8300 MI/m') and 400 ksi (2860 MI/m') respectively. The Strain Moduli (Uniaxial Compression, Tension, Bending, and Slip and Shear)

In studying the behavior of ice under static load it is useful to intro-duce the concept of the strain modulus. The reason is that the stress-strain diagrams have considerable linear parts,

from the total, i.e., the elastic, the recovered, and the remanent, strain. It therefore depends not only on the temperature and structure of the ice but also on the magnitude of the load and the time during which it acts. The In contrast to the rodulus of elasticity the strain medulus is computed burain modulus may therefore have many values for each kind of deformation. If only the properties of the samples are considered, all values of the strain modulus should be assumed to be equally probable, taking into account the dependence of the strain modulus on the above-mentioned factors. The scatter of its values indicates only that the test conditions or the structure

the various sources on the mechanical properties of ice should, therefore, be grouped according to their similar features. This, however, is not always possible due to the lack of information on the test conditions. The strain modulus must in any case be selected according to the test conditions. The less time the load acts on the ice, the larger should be the value of the The data found in of the ice in some tests differed from those in others. strain modulus.

When information on the properties of the ice cover is required, and the question arises of which value of the strain modulus should be given preference, there is, strictly speaking, no possibility of solving this problem by using the results of tests of camples having an arbitrary structure. Far more accurate data on the mechanical properties will then be provided by samples of structurally simulated ice.

The results of strain moduli in unlaxial compression, tension, bending, and slip and shear are presented in Table II-1.

	Strain Modulus Range	lus Range
	kui	MN/n;2
Uniaxial Compression	71- 768	9625 -067
Tension	17- 71	118- 490
Bending	135-1707	932-11768
Slip and Shear	6- 36	41- 245
	Table II. Strain Moduli Structurally Simulated Ice	

(after Lavrov, 1969)

The strain moduli in compression and tension increase under repeated load applications. The large differences in the strain modulus in bending are due not only to differences in the ice structure and temperature but also to the influence of the sample dimensions. This last factor was not considered important until recently, although it plays an important role in bending tests. The clear dependence of the strain modulus in bending on the ratio of the distance between supports to the sample height is due to the proximity of the ice temperature to the melting point and to the appearance of stress concentrations and microcracks at the corners of cavities. This causes irreversible and slowly reversible deformations which take place in the ice practically as soon as the load is applied.

The strain modulus in bending decreases when the load is increased and increases when a load is applied again unless the stress is very high.

The strain modulus in shear depends on the temperature, the concentration of impurities, the stress system, the dimensions of the shear surface, and the duration of the load application.

Ultimate Compressive Strength

The ultimate compressive strength of ice is, like that of other materials, usually established from the results of tests of samples (cubes) subjected to forces acting on one side. The sample is under these conditions deformed in the direction of compression but is free to expand in the transverse direction. The quality of work done in preparing the sample has a considerable influence on the test results.

Carefully prepared samples of structurally simulated fresh-water ice gave the following range of results: ullimate compressive strength with the direction of force perpendicular to the ice surface 730-1220 pc; (5031-8416 KW/M2) and for the direction of the force parallel to the ice surface 156-64 pp; (1079-4236 KW/M2).

Fonsile Strength

The tensile strength of ice has been little investigated, like its shear strength. This may be due to the lack of examples in which ice is under real conditions subjected to tension.

Lavrov reports on the research of several other people. They found that tensile strengths (el) in the range of 7) to 199 psi (490 to 1373 $kB/m^2)$.

Lavrov's tests on structurally chaulated ice show lower strengths than natural ice. This may be due to the different dimensions of the crystals and to differences in the development of the intercrystalline layers.

The Pifferences in the Properties of ice Under Tension and Compression

As noted, the tensile strongth and strain roddius in tension are considerably less than the corresponding magnitudes in compression. Lavrov carried out special twestigations on homogeneous structurally simulated ite. He found the compressive arteagh of this type of ite to be 9 times higher than its tensile strongth. This is due to the structure of the ice. Part of the vertical cross-sectional area of ice consists of intercrystalline layers and cavities filled with impurities and air. During tension the force is comheracted by the cohesive attength of the material connecting adjacent grains. Only part of the cross-section thus takes up the force. The crystal layers be dening the cavities can nove freely in opposite directions, thus altering the shape of the cavities and increasing their size.

During compression the cavities become filled and the untile cross-section of the sample thereafter takes up the force.

On the basis of the experimentally established differences in the ultimate strength (1:9) it may be assumed that in fresh-water ice not subjected to any external force about 90% of the cross-sectional area is taken up by intercrystalline layers and cavities, with only 10% of it formed by direct bonds between adjacent crystals, with only 10% of it formed by direct bonds between adjacent crystals, the ice (including the artio is not always observed since the quality of the ice (including the extent to which intercrystalline layers have developed in it) depends on the conditions of its formation and its growth rate. Furthermore, the microclructure of the ice canges with the lapse of time, due to recrystallization and other factors. There are thus cases in practice where the differences in the tensile and compressive strength of ice are less than 1:9, on the other hand, in thaving ice these differences are greater than 1:9.

Shear Strength

The results of Lavrov's tests carried out on structurally simulated ice gave a shear strength range of 68 to 422 psi (471 to 2913 kN/m²). The shear strength decends on the sample size (scale factor), the duration of the load action, and the temperature when these occurs along a circular surface in the center of an icc sheet. The direction of the force during the tests was purpondicular to the ice surface (along the vertical axes of the crystallites),

Bending Strength

Lavrov determined an average bending strength of 316 psi $(2177~{\rm kK/m}^2)$ for structurally simulated ice. The direction of the force was perpendicular to the ice surface and the duration of the load action was 4 seconds.

The bending strength was computed by means of the usual formulas of the strength of naturials, which are suitable for determining the stresses only when the elastic strains are mull, in the case of failure the strains are comparatively large and irreversible. Furthermore, the behavior of the ice at the beginning of the load action considerably differs from that assumed in the theory of clasticity. Hence, the bending strength, determined from the formula for a freely supported beam, cannot be considered as the true maximum tensile stress, but is an arbitrary magnitude. The computed building strongth therefore exceeds the computed tensile strength.

Lavrov determined an average bending strongth of 161 psi (1108 kN/m²) for structurally simulated ice in cantilever bending. The tensile strength and bending strongth are of similar magnitude when cantilever beams are tented.

The Freezing Strongth of Ico

Lavrov reports the research results of several others. In one case the average value of the normal freezing strength of ice to different surfaces at $2T^*F$ (-3°C) was about 284 psi (1961 kK/m²). Surfaces included were ice to copper, ice to iron, and ice to glass. Tests on ice to polystyrene gave values about one-tenth as much,

Researchers have concluded that the strength with which water and other liquids freeze to solid surfaces does not depend upon whether the latter are

wetted by the liquid considered or not. They state further: "The strength with which ice freezes to plastics and varmishes is approximately one-tenth of that which it freezes to glass or metals".

Other tests have given freezing strengths of fresh-water ice lower values, in the range of 71 to 142 psi (490 to 980 $k{\rm N/m}^2)$.

The Friction Coefficient

The friction coefficient was determined as the ratio of the towing force to the weight of the sample together with an additional load. The static friction coefficient was determined by starting the sample, while the dynamic friction coefficient was determined during steady motion of the sample,

An ice surface almost always carries a thin liquid film consisting of water molecules. The frictional force thus is the sum of two components, namely dry friction and boundary friction with water lubrication. This is supported by the increase in the friction coefficient when the tumperature of the ice surface is lowered.

It was found by experiment that steel skids are belter than wooden ones at temperatures down to $-22^{\circ}F$ ($-30^{\circ}C$), since the coefficient of boundary friction of steel is less than that of wood. Skids made of a material which is a poor host conductor are better at temperatures below $-31^{\circ}F$ ($-35^{\circ}C$), since part of the heat developed by friction is lost to the surroundings via heat-conducting skids; this interferes with the formation of a watery film, so that friction increases.

Researchers found that the static friction coefficient of ice on non-resty etcel is 0.15 to 0.25 and 0.30 to 0.35 on painted red lead.

Tests on Lake Ladoga were performed in an area cleared of snow. The ice surface was smooth and without cracks, but was uniformly covered by several hundreths to a tenth of an inch (1 to 3 mm) prominences. The lower surface of the ice was smooth.

The average friction coefficients of ice on steel and steel on ice were at 36°F (2°C) as follows: static friction coefficient ranged from 0.48 to 0.91 and the dynamic friction coefficient ranged from 0.08 to 0.15. The field results differed considerably from those obtained in the laboratory, perhaps due to scale effects.

Summary of Strongth Properties Reported by Lavrov

A summary of strength properties reported by Lavrov is presented in Table 11-2. The summarized values are qualified on one or more variables such as strain rate, temperature or sample size. They are applicable only for forces acting on ice for short periods, i.e., from a few seconds to a few minutes.

Property	Average Value or Kange of Values
Poisson's Ratio	0.33
Elastic Limit	less than 7 psi (50 kN/m^2)
Young's Modulus	1200 ksi (8300 MN/m²)
Shear Modulus	400 ksi (2800 MW/m²)
Uniaxial Compression Strain Modulus	71-768 ksi (490-5296 MN/m²)
Tension Strain Modulus	$17-71 \text{ ksi } (118-490 \text{ MM/m}^2)$
Bending Strain Modulus	135-1707 ksi (932-11768 NW/m ²)
Slip and Shear Strain Modulus	6-36 ks1 (41-245 MN/m ²)
Ultimate Compressive Strongth (Force perpendicular to surface)	730-1220 psi (5031-8414 kN/m²)
Ultimate Compressive Strength (Force parallel to surface)	156-614 psi (1079-4236 kH/m ²)
Tensile Strength	71-199 psi (490 to 1373 kN/m²)
Shear Strength	68-422 psi (471 to 2913 kN/m²)
Freely Supported Beam Bending Strength (Force perpendicular to surface)	316 psi (2177 kN/m²)
Cantilever Beam Bending Strength (Force purpendicular to surface and applied from above)	161 psi (1108 kN/m²)
Freezing Strength to Metals	71-284 psi (490-1961 kN ² /m ²)
Freezing Strength to Plastics and Varnishes	about one-tenth that of metals
i i i i i i i i i i i i i i i i i i i	:

Table II-2 Strength Properties (after Lavrov, 1969)

APPENDIX III

SAFE ICE LOADS COMPUTED WITH A POCKET CALCULATOR

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INTRODUCTION

When a load is placed on a floating ice sheet, the fee short is bent downward creating stresses within the ice. The downward deflection increases the water pressure on the bottom of the ice which in turn supports the load. If the maximum tensile strens within the ice exceeds the tensile strength of the ice, the ice will enack. Althou,n the ice sheet can still carry nore load before breakthrough occurs, the first crack is generally used to predict safe bearing capacity. In some cases flooding of the ice sheet may create operational difficulties, and a limiting defluction is used rather than first canek.

Previously the computations for the deflection and stresses of a floating ice sheet have been obtained by computers or programmable desktop calculators have become powerful enough to perform these computations. Although they are slower, they do provide a means for computation in the field which has been non-existent before.

The purpose of this report is to provide a program to calculate the deflection and stresses for a floating ice sheet on the Hewlett-Packard model 67 pocket calculator. CRREL does not necessarily endorse the HP-67 since there are other programmable pocket calculators that may also perform the same computations.

The user of the program must select appropriate values for the mechanical properties of the ice in order to compute reliable deflection and stresses. Engineering judgement must be used in selecting the allowable ice strength and when dealing with non-ideal situations.

REQUIRED INFORMATION

A rectangular coordinate system must be defined on the ice sheet for locating the loads relative to each other. For each load, the coordinates of the center of the load, the magnitude of the load, the distribution of the load must be known. The load may he assumed to be uniformily distributed over a circular area or a rectangular area whose dimensions and orientation are known. The load from a preumatic tire can be assumed to be uniformly distributed over a circular area whose radius can be determined from the air pressure in the tire.

Young's modulus, Poisson's ratio, and the ice thickness must be specified for the ice sheet, as well us the density of supporting water.

The coordinates where the deflection and stresses are to be corputed must be specified. For onfo bearing capacity prediction, one should choose the coordinates that will give the largest atress within the ice sheet. When the loads are sufficiently far apart, the largest stress will occur directly under the center of the huaviest load. When two leads of equal magnitude are extremely close together the largest stress may occur between the loads. A few trial computations may be necessary to find where the stresses are the largest.

MEGRETICAL APPROXIMATIONS AND LIMITATIONS

The celculator program uses the elastic thin plate theory to represent a floating ice sheet of uniform thickness, and considers a

which frequently is slowly convergent. This program uses a new solution expressed by a power series which is more rapidly convergent. In order rectangular area. This is a good approximation because the deflection load applied uniformly to a circular or rectangular area. Previously the rectangular area solution has been expressed by a Fourier series to further reduce the space and time requirements, the deflection is calculated from four concentrated loads, one at each corner of the does not depend significantly upon the distribution of the load.

pendent on the load's distribution as the calculation point moves farther stituted for the rectangular area. For distances greater thun five times requires much more time than the circular area solution. Further timethe flexural rigidity length 4, the deflection and stresses are nearly away from the load. Hence, at a sufficiently large distance from the savings can be obtained by noting that the stresses becomes less de-Even with these simplifications, the rectangular area solution load, an equivalent circular area or concentrated load may be subzero for elther the eirquiar or rectangular area solution and the calculation may be emitted.

used. For this theory solutions exist for both rectangular and circular load distributions, but they are expressed inconveniently for numerical computation. In practice, loads distributed over a circular area (e.g. The clastic thin plute theory is adequate except for calculating a preumatic tire) tend to be more concentrated than loads distributed this situation, the three-dimensional theory of elasticity should be the stresses in the vicinity of a relatively concentrated load. For

method proposed by Westergand when the stresses are evaluated directly The method replaces the radius a radius a is use than 1.742 times the ice thickness h. Using this new of the load distribution circle with $\sqrt{1.6}$ a² + h² -0.675 h, when the over a rectangular area (e.g. a crawler track). The program uses a radius in the thin plate theory gives the same stress as the threeunder the center of a circular area. dimensional theory of elasticity.

EQUATIONS

The most convenient form of the solution for a load uniformly distributed over a circular area has been given by Max Wyman. The symbols are defined as follows:

- E is Young's modulus of the ice,
- v is Poisson's rutio of the ice,
- h is the thickness of the ice,
- , the flexural rigidity length, k is the veight density of the vater, $\frac{1/4}{1} \text{ is } \left\{ Eh^3/12k(1-v^2) \right\} \text{ , the flexural}$
- P is the total load,
- r and 8 are polar coordinates from the center of the load,
- R is r/k, the non-dimentional radial coordinate,
- a is the radius of the load's distribution,
 - A is a/k, the non-dimensional load radius,
- Is the deflection of the ice sheet,
- o is the radial stress at the bottom of the ice sheet.
- is the tangential stress at the bottom of the ice sheet,

When R > A, the equations are:

$$\frac{(\sigma_{r} - \sigma_{\theta})}{2P(1 - v)} \frac{h^{2} \pi}{3} = \frac{2}{R} \left[-D \text{ ket'R} + C \text{ ker'R} \right] + \frac{(\sigma_{r} + \sigma_{\theta})h^{2} \pi}{2P(1 + v) \ 3} ,$$

$$\frac{(\sigma_x + \sigma_\theta)^2 \pi}{2P(1+\nu)^3} = C \text{ ket } R + D \text{ ker } R,$$

where
$$C = \frac{ber^{1/4}}{A}$$
 and $D = \frac{be1^{1/4}}{A}$, when $A = 0$, $C = 0$ and $D = 1/2$.

When A > R, the equations are:

$$\frac{(\sigma_x - \sigma_0)h^2 \pi}{2F(1-v)3} = \frac{2}{R} \ [-D \text{ beith } + C \text{ bertR}] + \frac{(\sigma_x + \sigma_0)h^2 \pi}{2P(1+v)3} \ ,$$

$$(\sigma_{r} + \sigma_{\theta})_{1}^{2} \pi$$

$$\frac{2P(1+v)_{3}}{2P(1+v)_{3}} = C \text{ bei } R + D \text{ ber } R,$$

where
$$C = \frac{ker^{1/A}}{A}$$
 and $D = \frac{ket^{1/A}}{A}$. When $R = 0$, $\sigma_{r} = 0$.

The stresses must be determined relative to the x, y coordinate system in order to add to the stresses from other leads. The equations for this are

$$\frac{\alpha - \alpha}{2} = \frac{\alpha - \alpha_{\theta}}{2} \cos 2\theta$$

Consider a load uniformly distributed over a rectangular area as shown in Figure 1. The symbols are defined by:

x and y are coordinates where the stresses are evaluated, x_0 and y_c are coordinates where the stresses are evaluated in a x_p and y_c are coordinates where the stresses are evaluated in a coordinate system whose origin is at the center of the load and is parallel to the sides of the rectangle,

0 is the angle of the $x_{_{\rm I}}$ axis measured from the x axis in a counterclockwise direction,

X is $x_{\rm r}/\epsilon$ and Y is $y_{\rm r}/\epsilon$, non-dimensional coordinates,

a is one kalf the length of the load in the $\mathbf{x}_{\mathbf{r}}$ direction,

b is one half the length of the load in the $\mathbf{y}_{\mathrm{F}}^{-}$ direction,

A is a/t and b is b/t, non-dimensional longth of the lond, and

E, v, h, k, 2, P and w are defined as before. The equations to determine the $\mathbf{x}_{\mathbf{r}}$ and $\mathbf{y}_{\mathbf{r}}$ coordinates are:

$$x_{F} = (x_{o} - x) \cos \theta + (y_{o} - y) \sin \theta$$

 $y_{F} = -(x_{o} - x) \sin \theta + (y_{o} - y) \cos \theta$

The equations to determine the stresses at the bottom of the ice sheet relative to the X, Y coordinate system are:

$$\frac{(\sigma_{\chi} \tau_{\sigma_{\chi}})^{2}}{2P(1\pm\sigma)} = \frac{3}{\chi\pi AB} \left[I(X+A, Y+B) - I(X+A, Y-B) + I(X-A, Y-B) - I(X-A, Y+B) \right]$$

$$\frac{3}{4\pi AB} \left[I(Y+B,X+A) - I(Y-B,X+A) + I(Y+B,X-A) - I(Y+B,X-A) \right]$$

$$\frac{\sigma_{XY}}{P(1-v)} = \frac{3}{\ln AB} \left[\ker \sqrt{(X+A)^2 + (Y+B)^2} - \ker \sqrt{(f_1+A)^2 + (Y-B)^2} \right.$$
+ \kei \sqrt{(X-A)^2 + (Y-B)^2} - \kei \sqrt{(X-A)^2 + (Y+B)^2} \right]

there

$$I(x,y) = (\frac{x}{2})^2 \binom{x}{x} \sum_{n=1,3}^{\infty} \frac{\frac{n+1}{2}}{n! \frac{(x/2)}{(x+1)!}} \left[-\frac{\pi}{L} \right] \sum_{k=0,1}^{\infty} \frac{n! \binom{k/x}{k/x}}{k! \frac{(k-k)!}{(k-k)!}} \frac{1}{2^{k+1}}$$

$$+ (\frac{x}{2})^2 \begin{pmatrix} x \\ y \end{pmatrix} \begin{pmatrix} x \\ z \end{pmatrix} = \frac{n/2}{x} \begin{pmatrix} x/2 \\ z \end{pmatrix} = \frac{2n}{(x/2)!} \left[-\sqrt{\ln \frac{x^2 + y^2}{2}} + \phi(n) + \frac{5}{n+1} \right] = \frac{n}{n+1} \frac{n! (x/x)}{k! (n-k)!} \frac{1}{2k+1}$$

$$+ (\frac{x}{2}) \left(\frac{x}{x} \right) \left(\frac{x}{x} \right) \left(\frac{(x-1)}{x} \right) \left(\frac{(x/2)}{(n+1)!} \right) \left(\frac{x}{x} \right) \left(\frac{x}{x} \right) \left(\frac{x}{(n-k)!} \right) \left(\frac{x}{2k+1} \right) \left(\frac{\arctan(x/2)}{(y/x)} \right)$$

$$\sum_{r=0,1}^{K} \frac{(-1)^{r} (v/x)^{2r}}{2r+1}$$

where γ is Eiler's number and $\phi(n)=\frac{1}{3}\cdot\frac{1}{2}\cdot\frac{1}{3}\cdot\cdot\cdot\cdot\frac{1}{n}$. The equations to determine the stresses relative to the x, y coordinate system are:

$$\frac{\sigma_X + \sigma_Y}{2} = \frac{\sigma_X + \sigma_Y}{2}$$

$$\frac{\sigma_X - \sigma_Y}{2} = (\frac{\sigma_X - \sigma_Y}{2}) \cos 2\theta - \sigma_{XY} \cdot \sin 2\theta$$

$$\frac{\sigma_{XY} - \sigma_Y}{2} = (\frac{\sigma_X - \sigma_Y}{2}) \sin 2\theta + \sigma_{XY} \cos 2\theta$$

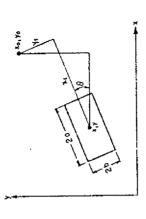


Figure 1. Rectangular Load distribution.

GENTERAL PROCEDURE

Turn on the calculator and set the other switch to run. The program to compute the deflection and the stresses for a load uniformly distributed over a circular area is divided into two parts which are for a load uniformly distributed over a rectangular area is divided into two parts which are labeled [3] and [4]. The results of either program $\mathbb Q$ or [4] are:

	(inches)	(ps1)	(ps1)	(psi)
	>	$(\sigma_{\mathbf{x}}^{+\sigma_{\mathbf{y}}})/2$	$(\sigma_{\mathbf{x}}^{-\sigma_{\mathbf{y}}})/2$	2× 0
Identification #	deflection	stress sum	stress difference $(\alpha_x - \alpha_y)/2$	sheering stress oxv
) ()	REG 6	REG 7	REC 8	REG 9

Program 5 adds the deflection and stresses from either program (2) or [4] to the respective sums which are in registers I through 4. From these new sums it determines and stores the maximum stress in register A and the angle from the x axis of the crack in register B. Between computation of each load, the sums are saved on a magnetic card referred to as the sum card. Side 2 of the program 5 magnetic card is not used, and can used for the sum card.

INPUT DATA

For a load uniformly distributed over a rectangular area (Fig. 1),
 the following data should be keyed into memory. Neither the length
 a or b of the rectangle may be equal to zero.

DEFINITION OF 100 Per control of the second	coordinate of the load's center	coordinate of the load's center	magnitude of the load	rectangle half-length along $\mathbf{x}_{\mathbf{r}}$	rectangle half-length along $\mathbf{y}_{\mathbf{r}}$
DATA (units) ID#	x (inches)	y (inches)	P (pounds)	a (inches)	b (inches)
REGISTER O	ਜ	eų.	m	-3	٠

the angle of $x_{\bf r}$ from x	Young's modulus of ice	Poisson's ratio of ice	thickness of ice	weight density of water	coordinate of computations	coordinate of computations
9 (degrees)	E (ps1)	2	h (inches)	k (1b/in ³)	x _o (frehes)	yo (inches)
9	-	æ	٥	¥	д	O

- 2. For a load uniformly distributed over a circular area, the input is the same except for the distribution of the load in registers 4, 5 and 6. Registers 5 and 5 are not used, while register 4 contains the radius a (inches) of the circular area. This radius may be equal to zero.
- 3. To store data into a register, key the data into the display register, press 570 followed by the register symbol. For example to store 1 for the 1D#, press 1, 570, 0. To display data from a register in order to check the input, press PCL followed by the register symbol.
- the data for each load. For each load the ice and water data as well as the place of computation do not change. The operator may find it convenient to record the input data for the first load on a magnetic card. Then for succeeding loads, this card may be read into memory and only data which are different need be keyed into memory.

ទ

When the computation is to be repeated for a different ice thickness or other parameter change, the operator may find it convenient to record putations, each card can be read into memory, and only the appropriate the input for each load on a magnetic card. Then for succeeding comparameter need be keyed.

secondary registers should be clear since this allows memory to be recorded on one side of a magnetic card rather than on two sides. If the input data is to be recorded on a magnetic card, the Consult the HP manual for ways to clear these registers. To record the memory registers on a magnetic card, press f, W/DATA and insert the card. To read the registers from a magnetic card into memory, just insert the card.

5. The units of the input data are pounds, inches and degrees. A different force and length unit may be used for input which will give the deflection and stress expressed in these same units.

Programs () and (3)

area. Program (2) is a continuation of program (1) and assumes that the results of program () are still in memory. The operating instructions Program (1) is for a load uniformly distributed over a circular

- 1. Read program (1) by inserting side 1 of the (1) magnetic card (side 2 is not used).
- 2. Store input data in memory.
- Press A to start the program. After about 5 seconds, the calculator stops and displays 2, the next required program.

4. Read progrem (2) by inscrting side 1 of the (2) magnetic cand followed by side 2.

5. Press A to start the program. After about 30 seconds the

program stops and displays 5, the next required program.

6. Proceed to program 5.

area. Program 4 is a continuation of program 3 and assumes that the results of program [3] are still in memory. The operating instructions Program [3] is for a load uniformly distributed over a rectangular

1. Read program [3] by inserting side 1 of the [3] magnetic card followed by side 2.

2. Store input data in memory.

Press A to start the program. After about 60 seconds, the

calculator stops and displays 4, the next required program.

4. Read program . by inserting side 1 of the 4 magnetic card

followed by side 2.

5. Press A to start the program. After about 300 seconds, the calculator stops and displays 5, the next required program.

6. Proceed to program 5.

Program 5 is a continuation of program (2) or [4]. Program 5 assumes that the results of program (2) or (4) are still in memory. The operating instructions are:

2

- Read program 5 into memory by inserting side 1 of the 5 magnetic
 rd. Side 2 is not used by program 5.
- 3. For a succeeding load (any load except the beginning load of a computation), press A to start the program. The program pauses for about one second and displays the ID#. During this pause, insert the sum card, and the sums will be read. If the sum card is not inserted during the pause, the program will repeat until it is inserted.

For the BECINIIIG load of a computation ONLY, press 5 to start thu program. This zeros the sum and eliminates the reading of the sum card,

- When the program stops, CRD is displayed. Insert the sum tard, and memory will be recorded.
- 4 . The angle of the crack is stored in register \bar{b} and the maximum stress is stored in register A as well as being displayed. If this was not the last load of a computation, process any additional load by either program (1) or [3].

Subroutines

The subroutines Kelvin (z) and I(x,y) may be of use in other programs. The following is a brief description of these subroutines that "will allow their use for this purpose.

The Kelvin (z) subroutine is stored in steps 102 through 213 of program (Z). This subroutine uses registers A, B, I, and O through 9. The subroutine also uses Euler's number $\gamma = 0.5772156649$ which is assumed to be in register E. Therefore, be sure that register E contains γ . To execute the subroutine, place the argument z in the x register and

ress E. The subroutine will compute ker (z), ber (z), ker'(z), ber'(z) kei (z), bei (z), kei'(z), bei'(z) in registers I through 8 respectively. The argument z will be in register A. The method of computation is from Don Nevel, Tables of Kelvin functions and their derivatives, CRREL Technical Report 67, June 1959. The series is truncated then the last term becomes less than 10⁻⁶. The exponent 6 is stored in step 159 and may be changed to obtain other accuracies.

The subroutine I (x,y) is stored in steps 61 through 220 of program $\frac{1}{4}$. This subroutine uses registers A, B, I, and 10 through 19. The subroutine also uses Buler's number $\gamma = 0.577215649$ which is assumed to be in register E. Therefore, be sure register E contains γ . To execute the subroutine, place the argument y in register x, and press E. The subroutine will give I (x,y) in register x. The series is truncated when the last term becomes less than 10^{-6} . The exponent 6 is stored in step 135 and may be changed to obtain other accuracies.

Examples

The following numerical examples are provided for the purpose of checking and debugging the programs. For the Kelvin (z) subroutine with the accuracy set at 10^{-6} , the results for z=1 are:

	Ė	ģ	5	-0
	-4.949946366 -01	2.495660400 -01	3.523699135 -01	4.973965115 -01
	10	5	~	ø.
,	REG 5	REG 6	REG 7	REG 8
	щ	124	т,	44
! !				
•	то <u>-</u>	-0-	5	102
	2.867062088 -01	9.843817812 -01	-6.946038910 -01	-6.244575218 -02
	7062	3817	6038	4575
,	2.8	9.84	6.9t	6.2t
			ŧ	1
	Н	CI	ന	- ≠
	REG 1	REG 2	REC 3	REG L

7

For programs () and () there are various options which should be checked.

For RPA, A may or may not be zero. For APR, R may or may not be zero.

For APR and Rep, Westergaard's substitution may or may not be made. The following assumption of cover all cases.

The next two options can also be used to check program 5. The input

For A>R ar	For APR and ReO, Westergaard's substitution may or may not be made.	substitution may or	may not be made. The	data are:		
following	following exceptes cover all cases.	,300		PEG	$R > A \neq 0$	A>R≠O
The	The input data for progress (1) are:	n 🛈 are:		o	1	2
£83	3×A=0	A> E = 0	Westergaard A>R=0	r	O	63
-	300	100	0	€,	0	63
, N	100	100	. 0	m	10000	10000
, m	00001	10000	10000	7	20	20
י בי	o	50	5	۲	100	got
ţ-	106	106	106	æ	1/3	1/3
- α	27 -	1/3	2/2	σ	1.0	10
	· · ·	C C		¥	62.4/1728	62.4/1728
ъ.	or .			kr.	. 40	70
₫	62.4/1728	62,4/1725	62.4/1723	ţ	Ç	10
p)	170	100	0	,	2	2
Ü	170 .	100	o	The output d	The output data from program (2) are:	
The	The output data from program (2) are:	gram (2) are:		REC	R>A#0	A>R≠O
7	N>A=0	A>R=O	Westergaard A>E=0	9	5.974546370 -01	6.748975099 -01
9	5.582837411 -01	6.761493320 -01	6.790823603 -01	1	6,221160230 +01	1.857763777 +02
! —	6.216941289 +01	1.935510530 +02 .	2.807218953 +02	8	O	o
ന	o	o	Q	Ø.	-1.501987756 +01	-1.943671630 +00
6	-1.534340994 +Ol	0	O			

are:
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program
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results
The

<u> 1(x,z)</u>		(x/2) ² (y/x) ² 4	и и и и и и и и и и и и и и и и и и и
15# [2] 15# x x X+A	3F(1+v)/4π/Εη X - A Y - B πkg ² /P	20 3P(1-v)/47ABh ² Y R	
MEMORY REGISTER USE EXCEPT FOR HIPUT AND GUTPUT MELVIN (2)		2/2 - 2 PT Y Index n	Kerra Kerra Kerra Kerra Kerra Kerra Kerra Kerra Kerra Kerra
MEMOI EXCEPT FY O and O 1D# R R	Λ P/πκ ² 3P(1+v)/π _h ² 3F(1-v)/πh ²	ט א א ט	
BEG O T S	w = v v F o v	∢ в о а я н о г	2.1 2.1 2.1 5.1 6.1 6.1
$\Sigma(\alpha_{x} + \alpha_{y})/2$ $\Sigma(\alpha_{x} - \beta)/2$ $\Sigma(\alpha_{x} - \beta)/2$ $\Sigma(\alpha_{x} - \beta)/2$	maximum stress crack angle (degrees) he accuracy set at 10"6, the and 🗓 is:	Outrut	6.00887%0%7 -01 6.566463308 +01 7.01%7%10%1 +00 -1.337575308 +01
1.272352147 +00 2.479879800 +02 0 0 - 1.696354919 +01	A 2.649515292 +62 maximum stress B 4.500000000 +01 crack angle (c For the I (x,y) subroutine with the accuracy set t is I (1,.5) = $6.407935361 - 62$. The example to check program [3] and [4] is:	Input Out 0 0 10000 30 5	ر ج د د
ol.	A 2.649515292 462 maximum or 2.649515292 462 maximum or 3.649515292 462 maximum or 3.649515292 462 maximum or 3.649515292 462 maximum or 3.649515392 462 max	합 - 여 20 - 크 10	

> > a

5 STO F 5 STO E 5 STO

PROGRAM (1)

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29 - x₀-x 30 + P polar 31 RCI t 33 STO 1 R 34 R + 0 35 2 2 20 36 X 20 2 0 36 X 20 2 0 37 STO 2 2 0 38 NCL 3 P P/m 40 + P/m 40 + P/m 41 STO X 6 P/mkc² 42 RCL 9 h² 43 3 3 3 3 4 4 1 P/m 44 STO T 3P/mh² 45 X 3F/mh² 46 STO T 3P/mh² 47 STO X 8 3vP/mh² 48 SCL 8 3vP/mh² 50 - 3(1-v)P/mh² 51 STO X 8 3vP/mh² 52 RCL 1 R

PROGRAM (I)

		g - d ₀	3F(1-v)/πh		rect	$\{\sigma_{\mathbf{x}}^{-\sigma_{\mathbf{v}}}\}/2$, XX	, A	֓֞֜֞֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓	(0,+0,)/2	·			o		ť	•	A		KELVIH (z)	5	2/2	2/5	0	E	ker z	ker'z	ber'z	kei z	bei z
		ACL:		×																						STO 1			STO 5	sTO 6
\sim	Ċ	60 6	Ď.	81	88	89	8	16	95	93	ąę	95	8	6	8	. 0	* :	100	101	102	103	104	105	106	107	108	109	110	111	112
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			HCL D	× .	RCL B	+	STO B	RCL A	×	+	RCL A	0 * X	*	STO A	RCL (1)	327. C	· ·	× 1	RCE O	;	.	F1 0	GTO 8	FCL 4	χ _ε	1/X	÷	1.31. 8	STO X 6	RCL 2
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PROGRAM (2)

PROCRAM (2)

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STI CSB 3 CSB 4

126 127 128 129 130 131 132 133 134 135

sro - 5

125

STO B STO 2 RCL A X # 0 CTO 2

11.3 11.5 11.5 11.6 11.8 11.9 12.0 12.1 12.2 12.3 12.3 PROGRAM 3

FROGRAM 3

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HEN
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PROGRAM 3
    PROGRAM [3]
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	(y/x) ²	(y/x) ²		1+(y/x) ²	(x/2) ²	R ² /4	R/2	LN(R/2)	>-	Y+1.N(R/2)	Ç I	o +	Þ	'n	9+4.5	o ผ	ວເລ	0	ជ					'n	ជ	7	n+1	.5/(n+1)
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	I(x,y)	$(\sigma_{\chi}^{+}\sigma_{\gamma})/2$	$(\sigma_{\chi}^{-\sigma_{\chi}})/2$	X+X	X-3	I(x,y)	$(\sigma_{\chi}^{+}\sigma_{\chi})/2$	$(\sigma_{\chi} - \sigma_{\chi})/2$	X+X	X+B	I(x,y)	$(\sigma_{\chi}^{+}\sigma_{\gamma})/2$	(ax-ay)/2	3P(1+v)/4 #ABh	$(\sigma_{\mathbf{x}}^{+}\sigma_{\mathbf{y}})/2$	3P(1-v)/4mABh ²	$(a_{X}-a_{Y})/2$	20	, XX	28	$(\sigma_{\chi}^{-\sigma_{\chi}})/2$	rect	(ox-ax)/2		ž	•	rect	Š
	GSB E I(x,y)	٦	89	7	5	ы	STO - 7 $(\sigma_{\chi}^{+}\sigma_{\chi})/2$	∞+		RCL 2 Y+B		STO + T $(\sigma_{\chi}^{+}\sigma_{\chi})/2$				RCL D 3P(1-v)/4 mABh ²	ac.				RCL 8 $(\sigma_{X}^{-\sigma_{Y}})/2$							STO + 9 ° xx
	CSB E	ST0 - 7	STC - 8	RCL 1	RCL 5	GSB E	STO - 7	STO + 8	RCL 1	RCL 2	GSB E		STO - 8	RCL 3	STO X 7	RCL D	STO X 8	RCL C	RCL 9	RCL C	BCL 8	‡ £;	STO 8	+. գո;	6 ous	÷	± +	STO + 9
PROGRAM L	. 29 CSB E	30 STO - 7	31 STC - 8	32 RCL 1	33 RCL 5	34 GSB E	35 STO - T	36 STO + 8	37 RCL 1	38 RCL 2	39 GSB E	40 STO + T	41 STO - 8	42 RCL 3	43 STO X 7	hb RCL D	8 X O. X S 54	ויפ אכד כ	47 RCL 9	48 RCL C	49 RCL 8	50 R+	51 STO 8	52 m +	53 STO 9	54 R +	55 R ←	6 + OIS 95
PROGRAM L	. 29 CSB E	30 STO - 7	31 STC - 8	32 RCL 1	33 RCL 5	34 GSB E	35 STO - T	36 STO + 8	37 RCL 1	38 RCL 2	39 GSB E	STO + T	41 STO - 8	42 RCL 3	43 STO X 7	hb RCL D	8 X O. X S 54	ויפ אכד כ	47 RCL 9	48 RCL C	49 RCL 8	50 R+	51 STO 8	52 m +	53 STO 9	54 R +	55 R ←	6 + OIS 95
PROGRAM L	. 29 GSB E	Y-8 30 STO - 7	X-A 31 STC - 8	I(x,y) 32 RCL 1	$(\sigma_{\chi}^{+}\sigma_{\chi}^{-})/2$ 33 RCL 5	$(\sigma_{\rm X} - \sigma_{\rm Y})/2$ 34 GSB E	Y+B 35 STO - 7	X-A 36 STO + 8	I(x,y) 37 RCL 1	$(\sigma_X^+ \sigma_Y^-)/2$ 38 RCL 2	$(\sigma_{X} - \sigma_{Y})/2$ 39 GSB E	40 STO + T	Y-B 41 STO - 8	I(x,y) 42 RCL 3	$(\sigma_X + \sigma_Y)/2$ by STO X 7	$(\sigma_{\chi}^{-}\sigma_{\chi})/2$ hb rcl D	X+A 45 STO X 8	Y+B 46 RCL C	I(x,y) 47 RCL 9	$(a_X^+a_Y^-)/2$ hBCL C	(α _X -α _Y)/2 49 RCL 8	Y+B 50 R+	X+A 51 STO 8	I(x,y) 52 R +	$(\sigma_X^{+}\sigma_Y^{-})/2$ 53 sro 9	$(\sigma_{\chi} - \sigma_{\chi})/2$ 54 R +	Y-B 55 R ←	X+A 56 STO + 9

	LI LI	4 н	ρ, ,	•	$(y/x)^2$	$-(y/x)^2$	o,k	• p.	· н	ħ	h	Çŧ	2r	2r+1	P_/2r+1)	' 143	i,	ж	۱۰ ۲		ដ	ים י	P, 2r	ť				
-	197 STC 7		199 STO 8																		217 RCL 7				-			
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	113 RCL 9	- 711	115 X	116 STC + 1	117 SF O	118 GSE 4	119 STO + 1	120 LBL 1	121 RCL 2	122 CHS	123 STO 2	124 CF 0	125 CSB 3	126 π	127 X	128 4	+ 621	130 STO - 1	131 RCL 1		133 ABS	134 FEX	135 6	136 CHS	137 X < Y	138 oro 2	139 RCL 1	340 P + 3

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7echnical-Report 67

Tables of Kelvin Functions and their Derivatives

by D. E. Nevel

U. S. ARMY SNOW ICE AND PERMAFROST RESEARCH ESTABLISHMENT

Corps of Engineers Wilmette, Illinois

TABLES OF KELVIN FUNCTIONS AND THEIR DERIVATIVES

bу

D. E. Nevel

The bearing capacity of a floating ice sheet can be computed by considering it a plate on an elastic foundation. The differential equation describing this is

$$\nabla_{\chi}^{4} w + w = \frac{q}{k}$$
 where the operator

$$\label{eq:potential} \nabla_{\chi}^{\,\,2} \,=\, \frac{d^2}{d\chi^2} \,+\, \frac{1}{\chi^2}\,\,\frac{d^2}{d\,\theta^2} \,+\, \frac{1}{\chi}\,\,\frac{d}{d\chi} \ .$$

Assuming that the boundary conditions are symmetrical with respect to the origin, the differential equation becomes

$$\frac{d^4 w}{d\chi^4} + \frac{1}{\chi} \frac{d^3 w}{d\chi^3} - \frac{2}{\chi^2} \frac{d^2 w}{d\chi^2} + \frac{2}{\chi^3} \frac{dw}{d\chi} + w = \frac{q}{k}$$

where

q = uniform distributed load

k = foundation modulus

w = deflection

x, θ = polar coordinates

$$\chi = x/k$$

$$\ell = \sqrt[4]{\frac{Eh^3}{12k(1-\mu^2)}} = \text{characteristic length}$$

E = modulus of elasticity

h = thickness of the plate

μ = Poisson's ratio.

The solution of this differential equation is given in terms of Kelvin functions as (Wyman, 1950):

$$w = Aber(\chi) + Bbei(\chi) + Cker(\chi) + Dkei(\chi) + \frac{q}{k}$$

These Kelvin functions are also important in other physical problems, for instance in the fields of electrical engineering, heat flow, and shell analysis. Because of the wide application of Kelvin functions, there is a need for a detailed accurate table presently unavailable.

The formulas for these functions and their first derivatives are:

$$ber(\chi) = \sum_{n=0, 2, 4, 6}^{\infty} \frac{(-1)^{\frac{n}{2}} (\frac{\chi}{2})^{2n}}{n! \ n!} \qquad ber_{\phi}(\chi) = \sum_{n=2, 4, 6}^{\infty} \frac{(-1)^{\frac{n}{2}} (\frac{\chi}{2})^{2n}}{n! \ n!} \phi(n)$$

$$\begin{split} \operatorname{ber'}(\chi) &= \sum_{n=2,4,6}^{\infty} \frac{(-1)^{\frac{n}{2}} (\frac{\chi}{2})^{2n-1}}{(n-1)! \ n!} & \operatorname{ber'}_{\varphi}(\chi) &= \sum_{n=2,4,6}^{\infty} \frac{(-1)^{\frac{n}{2}} (\frac{\chi}{2})^{2n-1}}{(n-1)! \ n!} \varphi(n) \\ \operatorname{bei}(\chi) &= \sum_{n=1,3,5}^{\infty} \frac{(-1)^{\frac{n-1}{2}} (\frac{\chi}{2})^{2n}}{n! \ n!} & \operatorname{bei}_{\varphi}(\chi) &= \sum_{n=1,3,5}^{\infty} \frac{(-1)^{\frac{n-1}{2}} (\frac{\chi}{2})^{2n}}{n! \ n!} \varphi(n) \\ \operatorname{bei'}(\chi) &= \sum_{n=1,3,5}^{\infty} \frac{(-1)^{\frac{n-1}{2}} (\frac{\chi}{2})^{2n-1}}{(n-1)! \ n!} & \operatorname{bei}_{\varphi}^{\mathsf{i}}(\chi) &= \sum_{n=1,3,5}^{\infty} \frac{(-1)^{\frac{n-1}{2}} (\frac{\chi}{2})^{2n-1}}{(n-1)! \ n!} \varphi(n) \\ \operatorname{ker}(\chi) &= \operatorname{ber}_{\varphi}(\chi) - \operatorname{ber}(\chi) \ln \frac{\chi^{\chi}}{2} + \frac{\pi}{4} \operatorname{bei}(\chi) \\ \operatorname{kei}(\chi) &= \operatorname{bei}_{\varphi}(\chi) - \operatorname{bei}(\chi) \ln \frac{\chi^{\chi}}{2} - \frac{\pi}{4} \operatorname{ber}(\chi) \end{split}$$

 $kei'(\chi) = bei_{\phi}'(\chi) - bei'(\chi) ln \frac{\gamma \chi}{2} - \frac{\pi}{4} ber'(\chi) - \frac{bei(\chi)}{\gamma}$

where $\phi(n) = \frac{1}{1} + \frac{1}{2} + \frac{1}{3} + \dots + \frac{1}{n}$

and $\ln y = 0.577215664901533$ (Euler's constant)

Other important derivatives and integrals may be expressed in terms of the functions and first derivatives.

$$ber''(\chi) = -bei(\chi) - \frac{ber'(\chi)}{\chi}$$

$$ber'''(\chi) = -bei'(\chi) + \frac{bei(\chi)}{\chi} + \frac{2ber'(\chi)}{\chi^2}$$

$$\int \chi ber(\chi) d\chi = \chi bei'(\chi)$$

$$bei'''(\chi) = ber(\chi) - \frac{bei'(\chi)}{\chi}$$

$$bei'''(\chi) = ber'(\chi) - \frac{ber(\chi)}{\chi} + \frac{2bei'(\chi)}{\chi^2}$$

$$\int \chi bei(\chi) d\chi = -\chi ber'(\chi)$$

$$ker'''(\chi) = -kei(\chi) - \frac{ker'(\chi)}{\chi}$$

$$ker''''(\chi) = -kei'(\chi) + \frac{kei(\chi)}{\chi} + \frac{2ker'(\chi)}{\chi^2}$$

$$\int \chi ker(\chi) d\chi = \chi kei'(\chi)$$

TABLES OF KELVIN FUNCTIONS AND THEIR DERIVATIVES

$$kei''(\chi) = ker(\chi) - \frac{kei'(\chi)}{\chi}$$

$$kei'''(\chi) = ker'(\chi) - \frac{ker(\chi)}{\chi} + \frac{2 \cdot kei'(\chi)}{\chi^2}$$

$$\int \chi \cdot kei(\chi) \, d\chi = -\chi \cdot ker'(\chi)$$

Table I gives the Kelvin functions and their first derivatives. Table II gives ber $_{\varphi}(\chi)$ and bei $_{\varphi}(\chi)$ and their first derivatives. These functions were computed by means of a Bendix G-15 D electronic computer.

The tabulated value consists of a sign, followed by two digits representing a floating point exponent p as used in Bendix G-15 D Intercom 1000 D (Bendix Computer Division, 1958), followed by a twelve digit number. The correct value of the function is

Example:
$$ker'(0.050) = -52.199803991221$$
 in floating point or -19.9803991221 in decimal notation.

The computed functions are accurate to the last digit; except that $\ker(\chi)$, $\ker'(\chi)$, $\ker'(\chi)$, $\ker'(\chi)$, and $\ker'(\chi)$ are occasionally in error only in the last digit.

The values of the functions are plotted in Figures 1-3.

The appendix gives the method used in computing the functions.

The solution of the symmetrical plate problem is

$$w = A ber(\chi) + B bei(\chi) + C ker(\chi) + D kei(\chi) + \frac{q}{k}$$

radial slope =
$$\frac{1}{\ell} \frac{\mathrm{d}w}{\mathrm{d}\chi}$$

$${\tt radial\ moment\ =\ -\ kf^2} \bigg[\frac{d^2\ w}{d\chi^2} + \frac{\mu}{\chi}\,\frac{dw}{d\chi}\,\bigg]$$

radial shear =
$$-k\ell$$
 - A bei'(χ) + B ber'(χ) - C kei'(χ) + D ker'(χ)

tangential slope = 0

tangential moment =
$$-k\ell^2 \left[\frac{1}{\chi} \frac{dw}{d\chi} + \mu \frac{d^2w}{d\chi^2} \right]$$

tangential shear = 0

displaced vol of water = $2\pi \ell^2 \int \chi w d\chi$.

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1000 programming system for the Bendix G-15 computer.

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Wyman, M. (1950) Deflections of an infinite plate, Canadian Journal of Research,

A28.293-302.

* 5630 Arbor Vitae Street, Los Angeles 45, California.

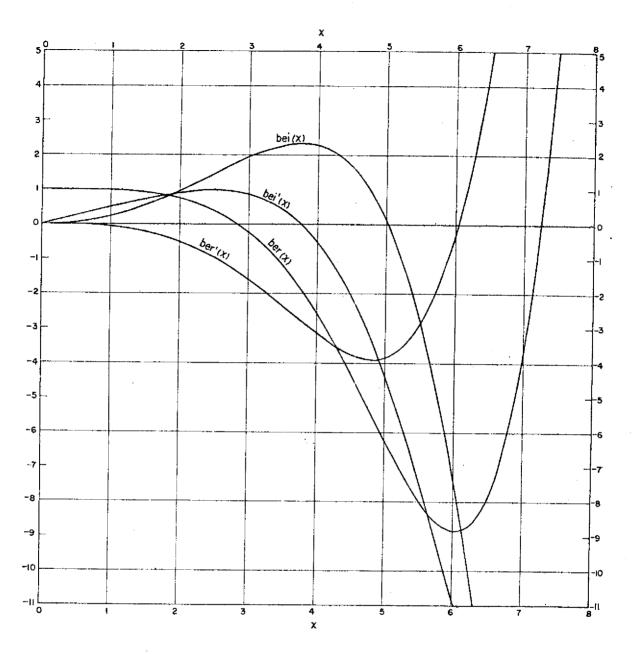


Figure 1.

TABLES OF KELVIN FUNCTIONS AND THEIR DERIVATIVES

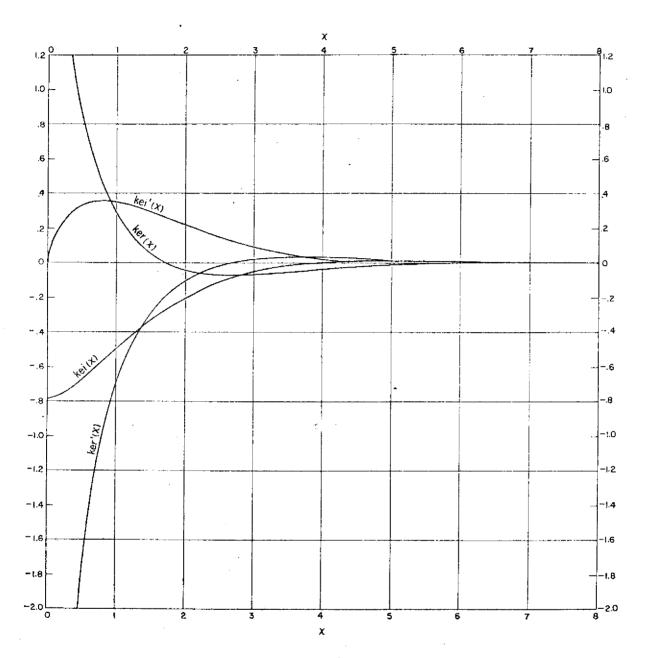


Figure 2.

×	000.	. 001 . 002 . 005 . 005 . 006 . 008	010.	. 011 . 012 . 013 . 014 . 015 . 016 . 019	,020	. 023 . 023 . 024 . 025 . 027 . 028	.030	. 031 . 032 . 035 . 035 . 037 . 038	.040	. 041 . 042 . 044 . 045 . 045 . 045 . 049	.050
, kei'(x)	00,000000000000000000000000000000000000	48.376184344641 48.683054000679 48.963761308432 49.122747880086 49.171957925790 49.217770061440 49.239691157616	49.261055575940	49.281919187028 49.30232656881 49.322319011807 49.341925409425 49.361174401096 49.360089856250 49.417001215798 49.435031909003	49,452799378625	49.470316802859 49.487596103614 49.504648118083 49.53108039734 49.554535359279 49.586818293086 49.602688658954	49.618386662194	49.633918056178 49.649288223084 49.664502208714 49.697564753100 49.709231108282 49.73887098230 49.738386045438	49.766992870694	49.781107335728 49.795099958041 49.808973645010 49.8363731168887 49.84908195196 49.86590842839 49.889665007169	49.902977267648
kei(x)	-50,785398163397	-50.785396157476 -50.78539633858 -50.785382581979 -50.785371613825 -50.785342076036 -50.785323710604 -50.785323710604	-50.785255135732	-50.785227983024 -50.785198167058 -50.78513431390 -50.785134316082 -50.785089158218 -50.785062092317 -50.784982363566 -50.784982363591	-50.784895365982	-56.784849208141 -50.784801310558 -50.784751696495 -50.784700388180 -50.78454746890 -50.784592773037 -50.784578625328 -50.784478625328	-50.784358093341	-50.784295476740 -50.784231315104 -50.784165624301 -50.784098419710 -50.7840929716249 -50.783959528405 -50.783814755487 -50.783814755487	-50.783664209054	-50.78356803016 -50.783507991649 -50.7834207992 -50.78334200792 -50.78317929456 -50.783092266519 -50.783006266468 -50.783006266468	-50.782828296864
$\ker^{-1}(\chi)$	8	-53.997999607301 -53.499999214606 -53.33332155248 -53.24998429231 -53.169698036557 -53.166664310560 -53.124996858606 -53.124996858606	-52.999960733824	-52.909047717080 -52.83286215695 -52.769179726189 -52.714230746078 -52.6660773544 -52.624937182230 -52.55684889343 -52.556241199494	-52.499921486573	-52,476108039650 -52,45459095238 -52,43692326980 -52,3990187493 -52,38451333951 -52,370264402855 -52,357032969197	-52.333215605833	-52.32245898559 -52.312374434761 -52.302900819633 -52.285376967473 -52.2763654875 -52.27012511926 -52.263008828082	-52.249843103762	-52,243741628833 -52,237930514524 -52,232389503167 -52,222100178801 -52,222045762050 -52,21210933189 -52,208145142079 -52,208145142079	-52,199803991221
ker(x)	8	51.70236899099 51.633054039948 51.592507627311 51.563739557599 51.54125379088 51.52319343938 51.494425781892 51.482647812097	51,472112133563	\$1.462581527874 \$1.45388841723 \$1.445877061759 \$1.433466794600 \$1.431568076759 \$1.423114332700 \$1.419053018817 \$1.419053018817	51.402803304708	51, 39792509041 51, 393273935458 51, 388829642406 51, 384574603392 51, 384574603392 51, 375799302404 51, 372799302404 51, 36565603763	51.362266606273	51.35898820862 51.355815187090 51.352739296467 51.34975314775 51.346857914484 51.341303754065 51.38638400446 51.338638400446	51.333512130640	51, 311044457913 51, 328636330425 51, 32987701587 51, 321742161046 51, 3139546054531 51, 317397257273 51, 31323378718	51.311215421270
bei'(x)	000000000000000	47.5000000000000048.100000000000000000000000	48.49999999974	48.54999999958 48.599999999935 48.649999999860 48.749999999802 48.749999999902 48.8499999999998 48.89999999998	48.99999999167	49.10499999886 49.109999999886 49.1149999999;93 49.124999999746 49.1249999999691 49.134999999681 49.1349999999626	49,14999999367	49.15499999254 49.159999999126 49.164999998817 49.174999998632 49.174999998632 49.18499998194 49.18499998194	49.199999997333	49.20499996983 49.20999996597 49.2149999995705 49.22499999636 49.234999996636 49.23999999636 49.23999999636	49.249999991862
$bei(\chi)$	000000000000000	44.2499999987 44.9999999987 45.22499999999 45.63499999999 45.8999999999 46.122499999999 46.12499999999	46.24999999996	46.30249999992 46.32599999987 46.42249999987 46.56249999950 46.63999999995 46.72249999995 46.80999999895 46.80999999999999999999999999999999999999	46.99999999722	47.110249999953 47.12099999951 47.13224999936 47.1439999991 47.156249999894 47.182249999835 47.1959999999986 47.19599999999999999999999999999999999999	47.22499999684	47.24024999615 47.25599999534 47.28599999329 47.306249999202 47.30624999905 47.34249998886 47.3424999886 47.380249998693	47.399999998222	47, 420249997938 47,440999997618 47,462249997256 47,483999998581 47,506249996396 47,552899999888 47,5528499995322 47,600249994692	47,624999993218
$ber'(\chi)$	-00.00000000000	-40.624999912935 -41.499999999737 -42.16674999697 -42.781249999415 -43.781249999415 -43.134999999999 -43.1349999999999999999999999999999999999	-43.62499999943	-43.831874999985 -44.107999999997 -44.177312499992 -44.210937499986 -44.2599999996 -44.307062499976 -44.4286874999950	-44.49999999917	-44,578812499901 -44,665499999861 -44,760437499812 -44,86399999940 -45,10984999956 -45,123018749943 -45,123018749943	-45.168749999881	-45.186193749849 -45.204799999813 -45.224606249768 -45.245469999715 -45.29159999974 -45.316581249485 -45.316581249485	-45.399999999110	-45,430756248942 -45,46304998748 -45,496918748525 -45,59531249998267 -45,608349997635 -45,608349997635 -45,69119996815 -45,69119996815	-45.781249995761
ber(x)	51,1000000000000	51.100000000000 50.9999999999999 50.9999999999	50.99999999844	50.99999999771 50.99999999676 50.99999999554 50.99999999209 50.999999998978 50.999999988695 50.999999988695 50.999999988695	50.99999997500	50.99999996961 50.99999996340 50.9999999981816 50.999999991896 50.999999991896 50.999999991696 50.999999991696	50.999999987344	50.99999985570 50.99999983616 50.99999991120 50.999999973756 50.999999973756 50.999999973756 50.999999973756	50.999999960000	50,99999955847 50,99999951380 50,999999941436 50,999999930040 50,999999930040 50,99999993755 50,999999917056 50,999999917056	50.999999902344
×	.000	000 000 000 000 000 000 000 000 000 00	,010	.012 .013 .013 .015 .015 .017	.020	.021 .022 .023 .024 .025 .025 .027	.030	.032 .032 .034 .035 .037 .038	.040	041 042 043 044 045 045 047	.050

×	.050	051 052 053 055 055 056 056 056	.060	.061 .063 .065 .065 .066 .068	.070	073 073 073 075 075 076	080.	. 083 . 083 . 085 . 085 . 088 . 088	060'	.091 .092 .093 .095 .095 .097	. 100
kei'(x)	49.902977267648	49.91598668440 49.928904173643 49.941722671910 49.95446960724 49.957078850346 49.9761298721 50.100443639363 50.101671478777	50,102690860541	50,104101926210 50,105304812709 50,10649952551 50,107686574053 50,10037155485 50,111201032768 50,1112357506758	50.114648521872	50,115783293586 50,116911043851 50,118031870476 50,119145868992 50,1203313258 50,121333750663 50,12244781224 50,123535402683	50.125691502375	50.126760172459 50.127822693298 50.128879140465 50.129929587721 50.130974107079 50.1310426641762 50.134072792904 50.135094287891	50,136110190856	50,137120564517 50,138125470217 50,139124965973 50,1401107973336 50,142601594716 50,143070035773 50,14404335494 50,145011591770	50.145974811429
kei(χ)	-50.782828296864	-50.782737347693 -50.782645102193 -50.78251570058 -50.78236083744 -50.782263488053 -50.782164762721 -50.782164762721	-50,781861596442	-50.78155099361 -50.781653395316 -50.78164098650 -50.78132121868 -50.78122669305 -50.781122550075 -50.78112050075 -50.781000270180	-50,780773259334	-50.780658042837 -50.780541695087 -50.760424223057 -50.780185923544 -50.780485252542 -50.779943228248 -50.779943228248	-50.779571004996	-50.77944778643 -50.779317466701 -50.779189135281 -50.779059730420 -50.778929276081 -50.77879724158 -50.778531694782 -50.778531694782	-50.778261508070	-50.77812489234 -50.77793725864996 -50.777709020611 -50.777568406627 -50.7775680469 -50.777284225164 -50.777284225164 -50.777284255164	-50.776850646537
ker'(χ)	-52.199803991221	-52.195878514818 -52.197103868649 -52.186471515201 -52.18649755959 -52.181602640968 -52.175515247705 -52.175215247705	-52.166431611445	-52.163695470357 -52.161647466818 -52.158483403849 -52.155999346779 -52.15599346373 -52.16125679386 -52.1448991387863 -52.1448991387863 -52.1448991387863	-52.142583116435	-52,140567151351 -52,138607078035 -52,13670599547 -52,13039852209 -52,131281577907 -52,129568872958 -52,129568872958 -52,126273249358	-52.124687086133	-52.123139992430 -52.121630538917 -52.120157365148 -52.117314735154 -52.115442866974 -52.113292405494 -52.113292405494	-52.110739401531	-52.109534526077 -52.108336195127 -52.107163552448 -52.10492057229 -52.103794726831 -52.103794726831 -52.102713975543	-51,996095939455
ker(x)	\$1,311215421270	51.309237138168 51.307297348284 -51.305394566556 51.303527475085 51.20564692102 51.20584525482 51.2058127272743 51.20583127272743 51.29583127272743	51,293004819381	51.2891354259242 51.289730616213 51.288133030133 51.285612788844 51.28348866623 51.281987422842 51.280508537175	51,277615206144	51.276199502373 51.274603676826 51.273427182177 51.2706949440 51.270736465773 51.26946323328 51.266817014947 51.265817014947	51,264291414628	51.263052311186 51.26828489119 51.260619579463 51.2584250235410 51.258245033482 51.25707882139 51.255926121318 51.25546671932	51,252546343031	51.251.44495880 51.25035564032 51.24278086354 51.24821221092 51.24715769515 51.24611429082 51.246081780618 51.24608938018	51.242047398104
bei'(χ)	49,249999991862	49.25499991015 49.259999900099 49.264999988043 49.27499998894 49.27499998894 49.28499998431 49.28999982907	49.299999979750	49.304999976055 49.3149999976142 49.3149999974155 49.324999969784 49.324999969784 49.324999964840 49.334999964840	49,349999956232	49.35499953015 49.35999966612 49.364999938213 49.374999938202 49.384999928511 49.384999928511 49.384999928511	49,399999914667	49,404999903453 49,409999903453 49,419999891091 49,424999881091 49,439998870203 49,439998870203 49,444999885482	49,449999846227	49.45499837492 49.45999828364 49.464999318831 49.4749997788495 49.479999787663 49.48999764604 49.494999752346	49.499999739583
bei(x)	47.62499993218	47.650249992363 47.67599991419 47.702249990380 47.728999989238 47.78399998614 47.812249987986 47.84099983477 47.870249981692	47.899999979750	47.930249977639 47.96099975347 47.992249972863 48.10239997017 48.108899996413 48.112224996674 48.115599995709	48,122499994894	48.126024994440 43.12959993953 48.13124993432 48.13689992873 48.144399991636 48.14822499991636 48.1502498990226	48.159999988622	48.164024987742 48.168099986805 48.172224985810 48.175399984753 48.180524983631 48.189224981179 48.19359979844 48.198024978430	48.202499976934	48.207024975353 48.211599973683 48.216224971919 48.225624968095 48.220399966026 48.2303999961552 48.245024963847	48,24999956597
ber'(x)	-45.781249995761	-45.829068745131 -45.87879994422 -45.930481243627 -45.107984374174 -46.1075999963 -46.115745623939 -46.121944998802 -46.128361873650	-46.13499998481	-46.141863123295 -46.156279372863 -46.156279372863 -46.16383997614 -46.17640622340 -46.17684997040 -46.187976877712 -46.196519996352 -46.205318120960	-46.214374995532	-46.223694358 -46.233279994558 -46.243135619006 -46.263671867758 -46.26371867758 -46.274359992054 -46.283333186293 -46.296594990470	-46.319999988622	-46.332150612588 -46.34604986475 -46.357366860278 -46.37643993990 -46.376439913990 -46.41564354533 -46.425919977828	-46.455624974051	-46.470981846964 -46.486679969735 -46.507723092356 -46.519114964818 -46.529599537113 -46.5295995231 -46.58244952901 -46.606436824432	-46.62499945746
ber (x)	50.999999902344	50, 99999894294 50, 99999885756 50, 99999876711 50, 99999857021 50, 99999857021 50, 99999885065 50, 99999883180 50, 99999883180	50.99999797500	50,99999783659 50,99999753860 50,99999753860 50,99999737856 50,99999773320 56,9999996651320 50,99999665135	50.999999624844	50.99999602942 50.999999580096 50.99999556277 50.99999505051860 50.999999478716 50.999999421640 50.999999421640	50,999999360000	50,99999327395 50,99999223560 50,99999222076 50,99999322076 50,99999184365 50,999999164868 50,999999062976	50.999998974844	50.99998828816 50.99998880636 50.99998831169 50.999998780080 50.999998672334 50.99998616730 50.99998616730 50.999988616730	50.99998437500
×	.050	055 055 055 055 055 055 055 055	090.	063 064 065 066 068 068	.070	071 072 073 074 075 076 078	.080	081 083 083 085 085 087 088 088	060.	.091 .092 .093 .095 .095 .098 .098	.100

×	.100	.101 .102 .103 .104 .105 .106 .107	011.	. 111 . 113 . 114 . 115 . 116 . 117	. 120	121 122 123 123 124 125 125	6. E.I.	131 132 134 135 136 136 138	.140	141 142 143 144 145 146 147	.150
$kei'(\chi)$	50.145974811429	50.146933060267 50.14786388082 50.14883484369 50.169778475002 50.150717328958 50.151651451666 50.152580888281 50.15425883072	50.155341521551	50.156252649665 50.157159305654 50.158061530027 50.158959362581 50.158959362581 50.160742007969 50.161626897003 50.161626897003	50,164256273192	50.165124421288 50.16598847426 50.166848460763 50.1677044909 50.16855382935 50.169404382387 50.170248450298	50.172757377699	50.173586029961 50.174410903581 50.175232027761 50.176049431268 50.17673189216 50.177673189216 50.177673189216 50.177673189216	50.180877276744	50.181669406466 50.18248031106 50.183243175923 50.1834803125380 50.185577978820 50.185577978820 50.185117562646	50.188643804715
$kei(\chi)$	-50.776850646537	-50.776704192189 -50.776556782056 -50.776408421037 -50.776259113977 -50.776108865679 -50.775957680896 -50.775852526657 -50.775852526657	-50.775343670425	-50.775187872965 -50.77503166617 -50.77487355831 -50.774715045021 -50.77435568857 -50.77439540774 -50.77439155966 -50.774034155966	-50,773745321795	-50.773580631105 -50.773415074318 -50.77324865514 -50.77391348007 -50.772913248007 -50.772744267295 -50.772574440552	-50.772059923140	-50.771886751120 -50.771712752339 -50.771531930562 -50.771362289524 -50.77188832930 -50.77188332930 -50.7700856460 -50.770651606665	-50.770291444418	-50.770110170784 -50.769928106774 -50.76974525881 -50.769561621574 -50.76937207293 -50.76912016458 -50.768819318678 -50.76863181849	-50.768443555102
ker'(x)	-51.996095939455	-51.986156310226 -51.976410828763 -51.966853840656 -51.957479908976 -51.948283803917 -51.93260493025 -51.930405131980 -51.921713055880 -51.91737771018	-51,904800947093	-51.896572409855 -51.886490134136 -51.880550237253 -51.87748972761 -51.85584801113 -51.85548001113 -51.842859754781 -51.835699825847	-51.828658600341	-51.821733135558 -51.84920585280 -51.806218195848 -51.801623302434 -51.795133325486 -51.7867458209017 -51.776268307106	-51.764172459478	-51,758262179242 -51,752440881170 -51,746706558490 -51,741057264334 -51,735491109520 -51,73400936991 -51,719273410717	-51.708845082680	-51,703741065634 -51,698708411861 -51,693745624560 -51,68851248513 -51,6745644 -51,67564629678 -51,6754629678 -51,66357339108	-51.660845025929
ker(x)	51,242047398104	51.241056288400 51.246075020770 51.239103403913 51.237188384622 51.23624462688 51.235309807855 51.235393962178	51.232557351151	51.231656676826 51.230764157579 51.229879649100 51.229003010895 51.228134106150 51.2227272801606 51.226418967432 51.225572477114 51.22572477114	. 51,223901037896	51.223075851555 51.22257533987 51.221445973665 51.220641061765 51.219842692089 51.21826516722 51.21826516722 51.217485812006	51,215945433363	51.215184223544 51.214428879345 51.213679312791 51.212935437885 51.212937170549 51.2104442855 51.210737131521 51.210015200759	51,208587131925		51.201744157255
bei'(x)	49.499999739583	49,504999726299 49,509999712479 49,514999688106 49,52999968735 49,529999651504 49,539999617362 49,54499959317	49,549999580596	49.554999561183 49.559999541057 49.5649995498390 49.579999452036 49.584999429050 49.58999946230	49.599999352000	49.60499924546 49.60999295170 49.614999266848 49.614999218555 49.624999172870 49.62499913626 49.634999105215	49.649999033091	49.654998993326 49.65998895390 49.664998874896 49.664998874896 49.6799888389 49.689988964184 49.68998696641	49,699998599417	49.704998548676 49.714998442784 49.714998342764 49.724998372441 49.724998372441 49.734998120828 49.734998120828	49.749998022461
$bei(\chi)$	48,249999956597	48.255024953927 48.260099951121 48.265224948175 48.270399945082 48.275624941836 48.280899933432 48.286224934864 48.286224934864	48.302499923109	48.308024918819 48.313599914331 48.3124999638 48.324899904732 48.336324899607 48.342224888665 48.342224888665	48.359999870400	48.3602486388 48.37209985688 48.378224849704 48.384399842222 48.390624834432 48.3968888334 48.403224817887 48.40524799988	48,422499790503	48.429024780646 48.435599770406 48.442224759770 48.448899748727 48.45524737264 48.462399725368 48.469224733027 48.476099700227	48.489999673197	48.497024658939 48.504099644166 48.511224628863 48.51839613016 48.55624596610 48.540224562053 48.540224562053 48.547599543871 48.5550245265	48.562499505615
ber'(x)	-46.624999945746	-46.643938066833 -46.663254937680 -46.782954308275 -46.703039928606 -46.744384915426 -46.76551787881 -46.7855187881 -46.809393025822	-46.831874894275	-46.854769262361 -46.878079880063 -46.901810497363 -46.92596486443 -46.92556486443 -46.975559846668 -47.100100796217 -47.10268948718	-47,107999980560	-47.110722541897 -47.113490478175 -47.1161304164392 -47.122070286630 -47.125023472646 -47.128023908590 -47.131071969458	-47.137312465957	-47.140505651581 -47.147039772562 -47.157039772562 -47.150381457912 -47.150715953313 -47.160709513356 -47.16425446289	-47,171499942810	-47.17520125238 -47.178955456839 -47.182762371159 -47.186523930343 -47.19538989385 -47.195538667027 -47.202611915616 -47.205746724043	-47.21093/40/303
ber(x)	50.999998437500	50.99998374056 50.999988308700 50.999998172096 50.99998100772 50.99998100772 50.99998100772 50.99999810077380 50.9999977380 50.999997794404	50.999997712344	50.99997628015 50.99997451375 50.999997452385 50.99997361000 50.999997170876 50.999997267178 50.9999972651 50.99996866658	50.99996760000	50.99996650643 50.99996538540 50.99996423647 50.99996185303 50.9999956061760 50.9999958035240 50.999995803695	50,999995537344	50.99995398439 50.99995256317 50.99995110927 50.999994610147 50.999994654657 50.999994654657 50.999994654657 50.99999465465	50.99999399750I	50,99993824155 50,99993647081 50,99993466226 50,99993281537 50,99999229060 50,99992703925 50,999922703925 50,99992270388	
×	.100	101 102 104 106 106 108	.110	1112	.120	.121 .122 .124 .125 .125 .126 .127	.130	.131 .132 .134 .135 .136 .137 .137	.140	141 143 143 145 145 146 147 148 149	27.

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kei'(χ)	50,188643804715	50.189401979800 50.190156887498 50.190908549877 50.191656988723 50.1931442225540 50.19314421560 50.193683177743 50.194618934791 50.195351573142	50,196081112983	50,196807574253 50,197530976644 50,198251339610 50,19888682370 50,20039438293 50,201102778154 50,20110808227714	50.203210362232	50.203907082769 50.204600928870 50.205291917816 50.20566539292 50.20666539292 50.208027640878 50.208704596513 50.208704596513	50.210050251326	50.210718982278 50.211385003153 50.212048329412 50.212708976346 50.213366959087 50.214674991703 50.214674991703 50.215325671045	50.216617428312	50.217259734792 50.217899478628 50.218536673734 50.219803472707 50.220433103704 50.221060240234 50.221684895525 50.222307082674	50.222926814650
kei(x)	-50.768443555102	-50.768254531937 -50.76806475223 -50.767874219243 -50.767490906333 -50.767490906333 -50.767908132815 -50.767910367806 -50.766910367506	-50,766519665393	-50.766323220794 -50.766126031264 -50.7657295853 -50.76553022489 -50.765330124538 -50.763129433731 -50.76492797963	-50.764522959434	-50.764319400471 -50.76411514626 -50.763910199566 -50.76370456337 -50.763498240373 -50.763291233488 -50.763291233488 -50.76329123488	-50.762456422455	-50.762246037612 -50.7623498394 -50.761823268504 -50.76181988628 -50.76119418593 -50.760869807732 -50.760859807732 -50.760859807732	-50.760322863257	-50.760105924461 -50.75988344641 -50.759670126353 -50.759451272139 -50.75921784526 -50.759011666029 -50.759011666029 -50.759595954550166	-50,758124933029
ker'(x)	-51.660845025929	-51,656391988323 -51,651997056703 -51,647659092058 -51,6437698465 -51,63149654626 -51,63476047967 -51,626785926675 -51,626785926675	-51.618798718015	-51.614878843595 -51.61006909208 -51.607182032834 -51.503403353968 -51.59570032969 -51.595981250426 -51.592336206557 -51.583734120624 -51.585174230368	-51,581655791468	-51.578178077013 -51.574740376998 -51.571341997837 -51.5645603069998 -51.561456050697 -51.558128366616 -51.554016730383	-51,548599303213	-51.545492343312 -51.542419127966 -51.539379104284 -51.536331731394 -51.531396480116 -51.53042832657 -51.527540282295 -51.524658333094 -51.521806499613	-51,518984306632	-51.516191288883 -51.513426990793 -51.510690966232 -51.50530199834 -51.50264820896 -51.497419962726	-51,492294852031
ker(x)	51,201744157255	\$1,201C85543638 \$1,200431353910 \$1,199781530536 \$1,199136017107 \$1,198494758108 \$1,1978549890 \$1,19785488646 \$1,196595972381 \$1,1965957199886	51,195350420917	51.194733586169 51.194120647250 51.193511556665 51.192906267786 51.192904734318 51.19170612874 51.191705222226142 51.189935275451	51,189351863864	51.188771950294 51.188195494372 51.187052797495 51.185052797495 51.1850537464037 51.1855637464037 51.1855637464037 51.1855637464037 51.185563766037	51.183703702501	51.183156659513 51.182612706566 51.182071810194 51.181533937474 51.180999056024 51.179938133990 51.179938133990 51.179412043215	51.178368420336	51.177850834950 51.177336028185 51.176823971544 51.176314636973 51.175807996851 51.17580796851 51.174802691574 51.174802691574	51.173314275374
bei'(x)	49.749998022461	49.75499785858 49.75999788762 49.764997816637 49.764997744347 49.774997590154 49.7849975915912 49.78499735787	49,799997269334	49.804997182927 49.809997094347 49.814997003552 49.814996810502 49.82499681554 49.839996717467 49.83999651467 49.839996514901	49,849996302457	49.854996192418 49.854996079776 49.864995864483 49.876995866493 49.8799950235 49.884995478811 49.894995346618	49.899995079251	49.904994941037 49.909994799734 49.914994655291 49.92999436776 49.92999436776 49.934994045070 49.939993884135	49.949993551828	49.954993380343 49.959993205230 49.964993826430 49.9799284388 49.979922657339 49.984992273199 49.989992075086	49.999991666670
bei(x)	48.562499505615	43.570024485507 48.577599464722 48.58522444342 48.50062439123 48.60062439123 48.6083937446 48.616224349997 48.62409324757	48.639999271822	48.648024244085 48.656099215474 48.664224185965 48.67239915537 48.680624124167 48.68899031832 48.705599024172 48.705599024172	48.722498952363	48.731023914839 48.739598876202 48.748223836426 48.75623753347 48.77439870989 48.78323665382 48.792098619497	48.809998523775	48.819023473879 48.828098422586 48.837223369863 48.845398315681 48.85562326006 48.87423144047 48.883598083695 48.883598083695	48.902497958078	48.912022892742 48.921597825673 48.931222756835 48.950622613700 48.950922613700 48.960397539327 48.96097384778 46.990022304521	48.99999722223
ber'(χ)	-47.210937407303	-47.215164340389 -47.219487898297 -47.223848456020 -47.22245638552 -47.232742070886 -47.237275878017 -47.241665184937 -47.241665184937 -47.251229798118	-47.255999854364	-47.26082910372 -47.265720341133 -47.270671521640 -47.283757831859 -47.2837531859 -47.291091240564 -47.291091240564 -47.30167534884	-47.307062277377	-47,312512955547 -47,318627758383 -47,323607060877 -47,324561238018 -47,34960664795 -47,3467967116 -47,35467467216 -47,3544192838 -47,35845836853	-47.364499667849	-47.370608467215 -47.376785141138 -47.383030064606 -47.3839343612607 -47.402178082152 -47.408899753671 -47.421953845131	-47.428687015044	-47.435491434392 -47.442367478162 -47.449315521336 -47.456335938901 -47.463429105839 -47.470395397135 -47.47835187772 -47.485448852732 -47.492536767000	-47.499999305556
ber(x)	50,99992089845	50.59991876789 50.99991659458 50.99991437795 50.999991211742 50.99999051242 50.999990506671 50.999990506671 50.999990013613	50,99989760003	50.99989501593 50.99988238323 50.999888970132 50.999988488744 50.9999878448744 50.9999878448745 50.99998753320 50.999987553220	50.999986949848	50.9998640066 50.99986334801 50.99986003989 50.99985345455 50.999985345465 50.999984466397 50.999984314447 50.99988318981	50,999883597507	50.999983229959 50.999982476366 50.999982496365 50.999981697656 50.999981299765 50.999980893276 50.99980823277 50.9998082827	50.999979637355	50.99979205272 50.99978766349 50.99978120513 50.99977407817 50.999776940811 50.99976940811 50.99976946602 50.999976946602	50.999975000017
×	.150		.160	. 161 . 163 . 164 . 165 . 166 . 168	.170	177 177 177 178 178 178	.180	181 182 183 184 185 186 187 188	.190	.191 .192 .194 .195 .196 .197 .198	.200

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	Ker(K) 50.222926814650	50.223544104294 50.22415896423 50.224771407328 50.225391445782 50.225989092035 50.227197256161 50.227787759934 50.22778779934	50,223991864452	50.229585410405 50.23017664411 50.230765586926 50.231352240297 50.231936618772 50.233518733494 50.233676215749 50.233676215749	50.234824774221	50.235395733851 50.235964494522 50.23631066698 50.237095460756 50.23765786981 50.238775576623 50.239331460169 50.239331460169	50.240436654390	50.240986084682 50.24153341.703 50.242078660056 50.242621824265 50.243701952951 50.244238936082 50.2447738773877387	50,245837670957	50,246366541283 50,246893405894 50,247418273620 50,247941153243 50,24636085466 50,246346959191 50,250012963762 50,250526032072	50,251037163490
kei(x)	-50.758124933029	-50.757901697366 -50.75767845830 -50.75745380243 -50.75702618149 -50.75702618149 -50.75654943022 -50.75631932498 -50.756093835405	-50,755865141279	-50.755635852449 -50.755405911228 -50.75517549920 -50.754712796198 -50.754480568334 -50.754247599482 -50.754014371690	-50,753545869420	-50.753310758982 -50.753075078685 -50.752838830722 -50.752164646524 -50.752126702613 -50.75188820728 -50.75140954357	-50.751169332456	-50.750928670911 -50.750687410986 -50.750445604774 -50.75960361815 -50.749716929208 -50.749472938593 -50.749472938593	-50.748737839119	-50.748491736845 -50.748245106705 -50.747907950699 -50.74750270820 -50.747532060053 -50.74755347370 -50.74604107741 -50.74654352121	-50,746253300703
$\ker'(\chi)$	-51,492294852031	-51,489770011145 -51,487269815728 -51,484793901844 -51,482341912691 -51,4751816023 -51,475126029050 -51,475126020560 -51,475126020560	-51.468113272718	-51.465819330444 -51.461295068216 -51.459064154208 -51.45683564779 -51.45463316528 -51.452492831275 -51.450341935949 -51.448210362465	-51.446097847613	-51,444004132948 -51,441928964681 -51,439872093578 -51,437833274854 -51,433808337073 -51,423808337073 -51,429853778395	-51.425966291032	-51.424047333760 -51.422144629539 -51.420257954432 -51.418337108154 -51.41653188897 -51.4169209269 -51.41058026025 -51.41058026025	-51.407483370031	-51,405717861505 -51,40396659200 -51,40526476683 -51,398797140358 -51,397101427911 -51,393750198726 -51,39209435635	-51.390451493607
ke r(x)	51,173314275374	51,172523245012 51,172134727137 51,171348697286 51,171365131358 51,170884065603 51,170884065613 51,17088406517 51,169927821339 51,16993741310	51,168514172074	51.168047207560 51.167582526310 51.167120107167 51.16659529269 51.1665201972050 51.165746263878 51.164341223034 51.164391948484	51,163944795956	51.163499746522 51.163056781508 51.162177031303 51.161740210006 51.161740210006 51.160872586551 51.160872586551 51.16041749704	51.159585940749	51.15816093297 51.158737840657 51.158316640693 51.157897319472 51.157064250552 51.156650471993 51.156238510455 51.156238510455	51,155419978838	51.155013379421 51.154608538346 51.154205441393 51.1538040742285 51.153806475700 51.157610215516 51.152215532934 51.151822711742	51,151431439890
bei'(x)	49,99991666670	50,100499145624 50,100999124159 50,101499102264 50,102499657160 50,102499657160 50,102996033339 50,102996033339 50,102998033333 50,102998033333	50,104998936433	50.105498910867 50.105998884413 50.106498383208 50.106498331208 50.105498377551 50.108498746553 50.10849874551	50,109998657909	50.110495627129 50.110998595786 50.11149853138 50.11249849307 50.112998464638 50.113998395489 50.113998395489	50.114998323870	50.115496287114 50.115993249717 50.116498211668 50.117498133537 50.117498133537 50.117938095337 50.118458052802 50.118458052802	50.119997926402	50.120497882840 50.120997838550 50.121497797531 50.121497747745 50.122497703133 50.123997659314 50.12399756582 50.123997556982	50.124997456871
bci(x)	48.99999722223	49.101002213784 49.102009705133 49.103022196266 49.105062177653 49.105062177653 49.107122138540 49.108159643523 49.109202138261	49,110249627751	49.111302116968 49.112359665967 49.1134220946833131 49.114489583131 49.116639595202 49.117722046815 49.118809534139	49.120999507900	49.122101994326 49.123209480441 49.124321966239 49.125439451716 49.12561936855 49.127689421680 49.128621906156 49.12995990286	49.132249357483	49.13401840539 49.13455932324 49.136889287455 49.138061768988 49.13923950124 49.140421710856 49.140421710856	49.143999170560	49.145201649607 49.146409128215 49.147621606376 49.148839084083 49.150061561328 49.151289033104 49.152521514404 49.153758990219	49.156248940363
ber'(x)	-47.499999305556	-47.507536843382 -47.512838416774 -47.52838416774 -47.5306320202 -47.538444487024 -47.56536053973 -47.554358053973 -47.562431086159	-47,578811522847	-47,597119677306 -47,5955695813 -47,603973733343 -47,61252036434 -47,6147285383 -47,62985489845 -47,63643333235 -47,647513230529	-47.665498646725	-47.674614915574 -47.603814058222 -47.593096449642 -47.702462464805 -47.711912478683 -47.71146866247 -47.731066002467 -47.750560020557	-47,760435652764	-47.770397533306 -47.780446037348 -47.790881539859 -47.800804415805 -47.811115040152 -47.811513787866 -47.832001033912 -47.832200133912	-47.863997511681	-47.874842500691 -47.885777862849 -47.896803973116 -47.91921206452 -47.919129937817 -47.930430542171 -47.941823394471 -47.953308869677 -47.964887342745	-47,976559188632
ber(x)	50.99975000017	50.99974496256 50.99973465931 50.999734593216 50.999972464699 50.999971862302 50.999971311948 50.99997053560 50.999970187060	50.999969612369	50.99969029410 50.99966438104 50.99967338370 50.9999661330 50.99996531366 50.999965351566 50.999964364 50.999964353	50,999963397537	50.999962727487 50.999961359831 50.999961359831 50.99996662059 50.999958238206 50.999958211957 50.999957776046	50.999956274897	50.999955599487 50.999954734073 50.999953182881 50.999951328881 50.999951330621 50.999958733621 50.99994886590 50.99949865703871	50.999948160075	50.999947290662 50.99994610360 50.99994616721 50.999943703203 50.99994278431 50.999941842311 50.999940894753	50,999938964947
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kci'(x)	50.251037163490	50.25154636319 50.252053648799 50.25259019105 50.253062455351 50.25364055589 50.2540637808 50.255057469849 50.255057469849	50,256043744201	50.25634104089 50.25702262654 50.257509307477 50.257994166085 50.28477205948 50.28457482 50.25943763969 50.25943763969 50.25915486957 50.260391325464	50,260865381773	50.261337663036 50.261808176355 50.262276928781 50.262743927316 50.2637209178912 50.26413446849 50.264594520853 50.265052853243	50.265509472733	50.265964385989 50.266417599633 50.266869120242 50.26771894346 50.2677108433 50.268513588947 50.268658402287 50.269101554811 50.269101554811	50.269982902625	50,270421110420 50,270557682406 50,271292624733 50,271157644806 50,27157644806 50,273016219031 50,273443103903 50,273463103903	50.274292098732
kei(x)	-50.746253300703	-50.74602008777 -50.74575020810 -50.745497902117 -50.74499177779 -50.7449177779 -50.7448350930 -50.74428841269	-50,743717738817	-50.743461449739 -50.743204671223 -50.742947405105 -50.742689653217 -50.7421417379 -50.742172699409 -50.741653824289 -50.741653824289	-50.741133042232	-50.740871940562 -50.740610367496 -50.740348324797 -50.74008581423 -50.73955395445 -50.7395539545 -50.73953543754 -50.73953543751 -50.73954553	-50.738501022947	-50.738235285876 -50.737969094742 -50.737702451241 -50.737455357063 -50.737167813892 -50.736899823404 -50.736831387270 -50.736362507153	-50.735823421597	-50.73553219454 -50.73528257992 -50.735011504632 -50.73448053285 -50.734195680461 -50.73392878350 -50.73375992674	-50.733101912295
ker'(x)	-51.390451493607	-51,368621445837 -51,387204063976 -51,384506700587 -51,38460426220 -51,3005821991 -51,379301977134 -51,37757523067	-51,374703473628	-51.37193611630 -51.371695017088 -51.370207561719 -51.368731119185 -51.368731702076772 -51.36581076772 -51.36486633609 -51.361509808737	-51,360096894459	-51.358694160102 -91.357301493627 -51.355918784635 -51.35458292442 -51.35182805544 -51.35182932592 -51.350485371364 -51.349150849234 -51.347825655050	-51.346509689104	-51.345202853109 -51.343905050170 -51.342616184765 -51.340364891166 -51.337548234616 -51.33692234616 -51.33632670308	-51.333836630637	-51,332615983300 -51,331403471668 -51,330199012491 -51,32002524006 -51,32633134706 -51,32653134706 -51,324294676699	-51.321986525355
ker(x)	51,151431439890	51.151641804482 51.1565379277 51.150267392183 51.149499374707 51.14919773379 51.14873762463 51.14873762461 51.14882135329	51,147606672181	51.147232724582 51.146860281202 51.146489330835 51.146119864407 51.145751864607 51.145385327689 51.14565590913 51.144294370363	51,143933567864	51,143574173180 51,142859566874 51,14250435336 51,142150471778 51,1446809956 51,1446809956 51,1446809956 51,1446809956	51,140401338241	51.140055482727 51.139710929524 51.139367669647 51.139025694206 51.13864994405 51.138007386992 51.138007386992 51.13734778854	51,137000328478	51.136657102852 51.136335093799 51.13604293225 51.135746283118 51.135746285547 51.135746285547 51.134693016697 51.13464424819	51.133721863753
bei'(x)	50.124997456871	50.125497405600 50.125997353505 50.126497300577 50.127497192180 50.127497192180 50.127497680229 50.128497080229 50.1294969401	50,129996905895	50.130496845934 50.130996785046 50.131496663451 50.131996663451 50.132946342022 50.133996396721 50.1349463342	50,134996263311	50.135496193598 50.135996122849 50.135496051052 50.13749598196 50.13749582256 50.138495753150 50.138495753150 50.138495753150	50.139995518141	50,140495437534 50,140995355771 50,141495272841 50,141995188730 50,142495103426 50,142495016916 50,1434944929188 50,144494750025	50.144994658565	50.145494565834 50.145994471820 50.146494376508 50.146494279887 50.147494181943 50.14799408261 50.14899388031 50.14899388031	50,149993671888
bei(x)	49,156248940363	49.157501414676 49.138758888472 49.160021361743 49.161288834480 49.1612888346676 49.163838778321 49.165408719925 49.166408719925	49,168998659221	49.170301127981 49.171608596136 49.172921063678 49.17556096884 49.175588462529 49.178520925521 49.179558391852	49,182248318489	49.183600780774 49.184958242357 49.186320703228 49.18568816375 49.190438081457 49.191207996516 49.19420452885	49.195997908464	49.197400363244 49.198807817211 49.200220270355 49.201637722664 49.20360174126 49.203487654728 49.205920674460 49.207357523308	49.210247418304	49.211699864427 49.213157309617 49.216697197143 49.2156087197143 49.21903708077 49.220519521102 49.222059696413	49.224996835941
ber'(x)	-47.976559188632	-47.988324782294 -48.100018449669 -48.101213871276 -48.102418779947 -48.103633213378 -48.106690804497 -43.107334037176	-48.109849564247	-48.111121933629 -48.112404090235 -48.113696071559 -48.116399658343 -48.117631338791 -48.117631338791 -48.120304661270	-48.123018182490	-48.125772202401 -48.125772202401 -48.127164493101 -48.128567020956 -48.139979823458 -48.13783640238102 -48.132836402381 -48.13734589	-48.137199267962	-48.138674505714 -48.140160280567 -48.141656630013 -48.144681202657 -48.144681202657 -48.147748523587 -48.147748523587	-48.152430314135	-48.154012610060 -48.155605818010 -48.155209975476 -48.158825119950 -48.16045128923 -48.16208519887 -48.163736850334 -48.165396313753	-48,168748813478
ber(x)	50.999938964947	50,99937982513 50,9993698266 50,999935982113 50,999934964957 50,999931831205 50,999931836529 50,999931836529 50,999931836529	50.999928597642	50,99927492792 50,99926375170 50,99925244678 50,999922944686 50,999921974990 50,999910395696 50,999918185900	50.999916962535	50.99915725502 50.99914474699 50.999913210024 50.999911031375 50.9999910638650 50.99990831744 50.999906674982 50.99906674982	50.999903960256	50.999902580896 50.999901186731 50.999898777655 50.999898333563 50.9998954348 50.999895459903 50.999893590122 50.9998936122	50.999889487683	50.99988795547 50.999886407394 50.999883263158 50.999881666785 50.99988065785 50.999878424978 50.999875117014	50.999873437945
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kei'(x)	50,274292093732	50, 274714220401 50, 275134765988 50, 27553741254 50, 275971151929 50, 2768013003702 50, 277214053130 50, 277625261990	50,278443075752	50.278849691653 50.279254787508 50.279658368731 50.280660440706 50.28066078269 50.281257654457 50.281653742598	50.282441475586	50. 282833130781 50. 283223318625 50. 283612044214 50. 28438512867 50. 284769497976 50. 285152424922 50. 285152424922	50,286292602130	50.28669809630 50.287045599430 50.287419976336 50.286164710565 50.288534677369 50.288903450240 50.289903450240 50.289963450240	50,290001451859	50. 290364695471 50. 290726568255 50. 291087074756 50. 2913604006559 50. 292160401624 50. 29215527933 50. 29221627933	50,293572740801
kei(χ)	.50.733101912295	-50,732827409004 -50,73252486379 -50,732277139995 -50,7312219311418 -50,731446603929 -50,731446603929 -50,7318171596122 -50,730894176337	-50.730338106978	-50.729059460468 -50.729780408102 -50.729500951397 -50.728940331017 -50.728660170349 -50.72866170349 -50.7280379111359	-50,727533559333	-50.727250921907 -50.726967893560 -50.726684475757 -50.726116477616 -50.725831900182 -50.725546939101 -50.725546939101	-50,724689763343	-50.724403387019 -50.724116429197 -50.723829196291 -50.723541589713 -50.7225965861158 -50.722676541579 -50.72287654721 -50.72287654721	-50.721808181515	-50.721517998327 -50.721227452531 -50.72093645646 -50.7203645278836 -50.7203635360 -50.720961671323 -50.719769333226 -50.719476640715	-50.718390197814
ker'(χ)	-51.321986525355	-51,320843627973 -51,319708084169 -51,318579821361 -51,3184853151 -51,318238007296 -51,31818161496 -51,318045247790	-51,310879949200	-51.30807432736 -51.308741585180 -51.307682342831 -51.3056829642801 -51.3058342299 -51.30310179636 -51.30248303578	-51,300447408600	-51.299438809445 -51.298436277223 -51.297439755798 -51.295464524254 -51.29546454288 -51.294485702288 -51.2924485702288 -51.29244539333	-51.290627835722	-51.289677460806 -51.288732621427 -51.28793267925 -51.28689351235 -51.2859076349878 -51.285007634952 -51.283177091623	-51.281367349257	-51.280470164577 -51.27869094572 -51.278690945150 -51.27691634246 -51.27699337114 -51.275191889253 -51.27429249061	-51.272618227663
ker(x)	51,133721863753	51.133400449292 51.132080174046 51.132761030697 51.132443012000 51.132126110781 51.131810319937 51.13182041304 51.130869539650	51,130558120639	51.13024777507 51.129938503551 51.129630292135 51.129323136684 51.129017030689 51.128711967699 51.128407941125 51.128104945240	51.127502018916	51,127202076315 51,126903139275 51,12660520775 51,126408257779 51,126012301411 51,125717326781 51,125717326781 51,125423328069 51,125423328069	51.124547130037	51.124256977852 51.12396773271 51.1239218492 51.1231392184922 51.123105790262 51.122520321497 51.122535773248 51.122535773248	51,121687599263	51.121406680930 51.121126657246 51.120847523167 51.120591903877 51.120291903877 51.119739783586 51.119739783586 51.119455023416 51,119191123500	51.115918079090
bei'(x)	50,149993671888	50,150493565715 50,150993458120 50,151993238614 50,151993238614 50,1529931255 50,15299313255 50,15399286344 50,153992781926 50,153992781926	50,154992544510	50.155492423482 50.155992300887 50.156492176711 50.15649192351 50.157991794538 50.158991531564 50.158991531564	50,159991261891	50,160491124502 50,160990985390 50,161490844540 50,161990701935 50,16299041193 50,163490263424 50,163490113634 50,164489962006	50,164989808523	50.165489653168 50.165989495925 50.166489336776 50.1678889175703 50.167888847719 50.168483680772 50.169988511831	50.169988167900	50.170487992872 50.170987815779 50.171487455325 50.172487271927 50.172487086391 50.173486898696 50.173486898696	50.174986322482
$bei(\chi)$	49.224996835941	49.22800670251 49.228006707251 49.229519141288 49.23559006056 49.23559006056 49.23561886316 49.23156294718 49.237156294718	49,240246147993	49.241798572834 49.243355996457 49.244918418847 49.2464855139986 49.249613676452 49.251218095745 49.252805511724 49.252805511724	49,255995339670	49.257597751603 49.259205162154 49.260817571305 49.264034979039 49.264057185338 49.26584790184 49.267317193560 49.268954595447	49.272244394681	49.273896791991 49.275554187738 49.277216581903 49.28055367410 49.28223375474 49.283916142358 49.285603528323	49.288993295134	49.290695675939 49.292403054984 49.293832847709 49.297555181347 49.301014923068 49.302494657237	49,306242021435
ber'(x)	-48,163748913478	-48.170441916764 -48.173362021637 -48.175589096206 -48.17327574183 -48.179077457059 -48.180838874323 -48.182611773466	-48,186192257347	-48.187999917064 -48.189819238618 -48.191650259498 -48.19347549194 -48.19721349549194 -48.19721349288663 -48.200982165909	-48.204798135867	-48.206724100955 -48.208662102766 -48.210612178788 -48.21257436510 -48.2165352699 -48.2165352699 -48.220544984132 -48.222568292657	-48,224603937804	-48.22651957060 -48.22812387910 -48.230785267840 -48.23287063438 -48.231078976978 -48.231078976978 -48.241337715715	-48.245647150430	-48.247820972492 -48.250007581004 -48.25207013449 -48.254419307313 -48.25882262931 -48.261133732253 -48.26397846629	-48,267965259375
$ber(\chi)$	50.999873437945	50, 999871742001 50, 999863299037 50, 999864299037 50, 999864787217 50, 999861205201 50, 999861205201 50, 999859388385	50.999855700422	50.999853829471 50.999851940385 50.999848107341 50.999846163148 50.999842218830 50.999842218830 50.999840218469	50.999336160746	50.999834103144 50.99983922862 50.999823913940 50.9998235678334 50.99982352925 50.999821347589 50.999816936648	50.999814700798	50.999812444528 50.999810167717 50.999807870239 50.999803212785 50.999800822558 50.999800822558 50.99979647163 50.999796088475	50.999791198711	50.999788731381 50.999783731187 50.999783731187 50.99978198666 50.999778642758 50.99977605133 50.999770642415	50.999765528871
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Хеті	50.293572740801	50.293922477645 50.294270387996 50.294617976154 50.294961746398 50.29563202985 50.295631350046 50.296333733015 50.296672977077	50.297010928423	50.297347591175 50.29768269435 50.298017067281 50.29834988772 50.299011718814 50.299011718814 50.29968491607	50,300320238868	50, 300644237746 50, 30096691989 50, 301288505470 50, 301608182045 50, 30192782548 50, 302245639795 50, 302377595692 50, 303191744879	50.303504679884	50.303516404430 50.304126922220 50.304436236938 50.3044352253 50.30572121813 50.305661338175 50.305661338175 50.305664892187	50,306568059765	50, 306867880436 50, 307166530404 50, 307464013179 50, 30776033254 50, 308345491105 50, 30834549194 50, 308542341963 50, 308542341963 50, 308934040840	50.309514002548
kei(x)	-50.718890197814	-50.718596450094 -50.718302353300 -50.718007908758 -50.717713117787 -50.717417981703 -50.71712501618 -50.716530515867 -50.716530515867	-50,715937170344	-50.715639999977 -50.715342475590 -50.715044625465 -50.714746441881 -50.714149079428 -50.71350398377 -50.71350398377	-50,712950408812	-50,712649926470 -50,712349120751 -50,712047992899 -50,711746544152 -50,711444775746 -50,71142688911 -50,710840284875 -50,710840284875	-50.709931181776	-50,709627571133 -50,709323549370 -50,708714677295 -50,7087094677295 -50,708104575149 -50,707799065782 -50,7077935468 -50,7077935468	-50,706880718731	-50.706574000663 -50.706266983360 -50.705959667991 -50.7053652058722 -50.70534414713 -50.705035941125 -50.704727449111 -50.704418660324 -50.704418660324	-50.703800212019
ker'(x)	-51,272618227663	-51.27769765612 -51.27092949531 -51.270386740136 -51.2692269858 -51.2692269858 -51.2675963658 -51.266775199173 -51.266775199173	-51.264338053522	-51.262354335432 -51.262734889736 -51.261939681321 -51.261148675463 -51.260361837815 -51.259579134409 -51.25860531646 -51.258025996291	-51,256488996668	-51.255726467715 -51.254967876792 -51.254213192418 -51.253462383450 -51.25197246820 -51.251972268820 -51.25197268820 -51.251272902516 -51.25497290325	-51.249037210500	-51.248312684743 -51.246874518517 -51.246874518517 -51.246160821732 -51.24474062219 -51.244744062219 -51.244744100-51.2434129982	-51.241952321175	-51.241262934927 -51.240576916316 -51.29894239820 -51.298214680178 -51.29853812384 -51.237196453575 -51.235560968342	-51.235206993430
ker(χ)	51,118918079090	51.11%545865482 51.11%374538010 51.11%374538010 51.11%34365007 51.11%365526341 51.11%30333124 51.11%763965667 51.11%763965667	51.116233672060	51,115969736223 51,115706601965 51,115132721201 51,111132721201 51,114611966120 51,114612966147 51,114402806655 51,11444393728	51,113629881403	51.113373774001 51.11318427155 1.112863836945 51.112609998479 51.112165910897 51.112104567369 51.111652965098 51.111652965098	51.111102568276	51.110653893632 51.110605941693 51.1101538708335 51.110112191463 51.10986586007 51.109376895716 51.109376895716 51.108890212968	51,103647914541	51,108406307194 51,105165387548 51,107925152246 51,107685597953 51,10746721390 51,106970958286 51,106970412269 51,106497926994	51.106262390276
bei'(x)	50.174986322482	50.175486125970 50.175985927205 50.176485726169 50.17688532842 50.17798510937 50.1794898920 50.1784894686232 50.178484898920	50,179984253669	50,180484033752 50,180983811385 50,181483586548 50,181983359219 50,182483129378 50,182483685005 50,183432682078 50,183932424577 50,184480	50.184981941766	50.185481696414 50,185981443403 50.186481197710 50.187480688195 50.187480688195 50.188780459328 50.188460167693 50.18979903267	50,189979365954	50.190479093022 50.191478538493 50.191478538493 50.19247797260 50.19247797260 50.193477394138 50.19377100561 50.194476803942	50,194976504256	50.195476201462 50.195975895594 50.196475586568 50.197474959010 50.197474959010 50.198474318612 50.19897399337 50.19897399337	50.199973333511
$bei(\chi)$	49,306242021435	49.307914383679 49.309751743947 49.311514102216 49.315053812655 49.316831164799 49.318613518479 49.320400862770	49.323990552185	49.325792893624 49.327600232852 49.3294125996437 49.333052237019 49.334879557153 49.336731894051 49.35549220386	49,342238864067	49.344091182260 49.345948497987 49.3496781220 49.35155043097 49.3515043097 49.35310338674 49.353197339031 49.359089636730	49.360986931742	49.362889224040 49.364796513593 49.366708800374 49.368626084353 49.370548365501 49.374407919185 49.374407919185	49.380234727730	49.382186991261 49.384144251749 49.386106509163 49.38073763470 49.390046014640 49.392023265639 49.394005507437 49.395992749001	49.39998222293
ber'(x)	-48.267965259375	-48.270268632712 -43.272585167334 -43.274914900725 -48.277257870367 -48.27814113742 -48.281983668332 -48.286762861086 -48.289172574212	-48,291595748479	-48.294032421369 -48.296482630361 -48.298946412937 -48.30142386577 -46.301914848161 -48.306419556968 -43.306938028679 -48.31470241373 -48.314016252529	-48.316576099626	-48.319149820143 -48.321737451558 -48.324339031351 -48.32554586997 -48.33222783767 -48.3322783767 -48.337557467686 -48.340243524769	-48.342943792571	-46, 345658308567 -45, 346387110234 -48, 351130235047 -48, 35659604015 -48, 35659604015 -48, 36246715273 -48, 365062017947 -48, 365062017947	-48.370736304759	-48.373595363845 -48.376469083348 -48.37957500741 -48.38178579093 -48.38111314996 -48.391058698680 -48.394021367617 -48.396998759278	-48.399991111135
ber(x)	50.999765528871	50.999762837712 50.999760123454 50.999754625112 50.999754625112 50.99974602763 50.999746201046 50.999743345410 50.999740465744	50,999737561913	50.99973463784 50.999731681220 50.999728704086 50.999728702246 50.99972855564 50.999719623903 50.999713623903 50.999713623903	50.999707164726	50.999703086108 50.999700781683 50.999697551312 50.999694264856 50.999687703125 50.999684367570 50.999681005366 50.99968105365	50,999674200449	50.999670757450 50.999667287235 50.999660789660 50.999650264582 50.999653131342 50.99964582891 50.999645886359 50.999642221602	50.999638528473	50.999634806827 50.999631056517 50.99962727396 50.999623463318 50.99961346313 50.999611865858 50.999611865858 50.99603989381	50.999600004444
×	.350	. 351 . 352 . 353 . 355 . 355 . 355 . 358 . 358 . 358	.360	361 362 364 364 365 366 367 368	.370	371 372 373 374 375 376 377 378	.380	. 38.1 . 38.2 . 38.5 . 38.5 . 38.6 . 38.6	.390	. 391 . 393 . 394 . 396 . 396 . 397 . 398	, 400

×	007	. 402 . 403 . 403 . 404 . 405 . 406 . 407 . 403	410	. 411 . 412 . 413 . 415 . 415 . 416 . 417 . 418	420	. 421 . 423 . 424 . 424 . 425 . 426 . 427 . 427	.430	. 431 . 434 . 434 . 435 . 435 . 435 . 435 . 436 . 439	. 440	. 441 . 443 . 444 . 444 . 445 . 446 . 447 . 448	.450
$\ker^{-1}(\chi)$	50.309514002548	50,309802272;53 50,31089405416 50,310860376290 50,310860376290 50,3102264420551 50,31126644767 50,311738523198 50,312067799939	50.312345961691	50.312623016680 50.312898955118 50.313173019202 50.313447573114 50.313720233023 50.314262283434 50.314531680201 50.314799995496	50.315067232417	50, 31533394048 50, 315598483460 50, 315862503709 50, 316125457840 50, 31646481950 50, 3166481950 50, 317166673570 50, 317166673570	50.317680962873	50,317936538271 50,318191071422 50,318444565252 50,318697022675 50,31919883983 50,319696546134 50,319696546134	50,320190164262	50,3204,3544,376 50,3205,29716947 50,3205,22975781 50,32116,522669 50,32116,6712392 50,3216466715719 50,321885959407 50,32236,145839	50,322597720041
kei(x)	-50.703300212019	-50,703499553787 -50,703180497853 -50,702559857419 -50,702559857419 -50,70193765955 -50,70193765955 -50,701914953837 -50,701914953837	-50,700690818603	-50.700378334026 -50.700055522942 -50.699752536457 -50.69943225669 -50.699125641675 -50.69981178557 -50.69487658434 -50.698183261362 -50.698183261362	-50.697553661730	-50.697238461327 -50.69602295299 -50.69607264717 -50.695291270647 -50.695658496303 -50.695024680747 -50.695024680747	-50.694389832412	-50.694072023574 -50.6937399683 -50.693435641778 -50.693117070898 -50.6979174345 -50.692459890743 -50.691840279279	-50.691200390889	-50.690880077998 -50.690559520332 -50.690238718901 -50.69959538733 -50.688953095685 -50.6889109519 -50.6883108347604	-50,687986367932
ker'(x)	-51.235206993430	-51.23986461328 -51.233986461328 -51.23245857766 -51.23558331996 -51.23152583861401 -51.2313242389 -51.230673956386 -51.230673956386	+51,228776559935	-51.228149957739 -51.227526258493 -51.226905441289 -51.226287485421 -51.2250600383 -51.22506005867 -51.2236450581758 -51.223239915268	-51,222638703612	-51.22040213811 -51.221444426691 -51.220851323259 -51.220260884702 -51.219673092384 -51.219687927845 -51.218505372797 -51.217925409122	-51.216773134268	-51.21620087693 -51.215631111694 -51.215063838980 -51.215499052420 -51.2133473638 -51.213316870017 -51.212819440692 -51.21264490551	-51.211161602523	-51.21061375237 -51.210068256915 -51.209525100120 -51.208445746873 -51.207509557473 -51.207305551221 -51.206314409967	-51,205787195204
ker(x)	51.106262390276	51.150/027511958 51.105793283904 51.10559718001 51.10594520319 51.10465288728 51.104631834468 51.104401538439	51,103942725285	51,103714262269 51,103486424402 51,103159208792 51,103032612566 51,102036632874 51,102581266688 51,102356511789 51,102356511789	51.101685884050	51.101463544817 51.101241802722 51.101020655069 51.100500099187 51.100500132418 51.100360753126 51.100141955692 50.999237405167	50.994890436271	50,992725568019 50,99566410115 50,98641293441 50,987265125948 50,984122943162 50,9772998173 50,97729985295 50,97729982525	50.973495740875	50.971386866070 50.969287457980 50.967185493137 50.965092948233 50.96505800112 50.960924025775 50.956776507210 50.956776507210	50.952650211570
bei'(x)	50,19997333511	50.2001720-0111 50.200972660115 50.2014/2018/11 50.2018/1973/20 50.202471624675 50.2034727639 50.203470917099 50.20347011858053 50.20397058058	50, 204969829336	50.205459459604 50.205969086257 50.206468709267 50.207467944253 50.207467944253 50.208467164349 50.208966768746	50.209965966101	50.210465559004 50.210365148021 50.211464733124 50.211464314285 50.212462891476 50.21296346469 50.213463033876 50.213463033876 50.213462159976	50.214961716392	50.215461269668 50.215960818273 50.21646035280 50.217459438780 50.217958976414 50.218458497731 50.218458497735	50.219957053486	50.22045656328 50.220956068523 50.221455569311 50.222454557272 50.2224544534 50.223453526875 50.223453526875 50.223453526875	50.224951946314
$bei(\chi)$	49.399982222293	49,4019545252 49,403991632252 49,408021128609 49,410343346602 49,412070561091 49,416139979421 49,41613979421	49.420229383319	49,42281579767 49,424338772499 49,426400961480 49,420540326040 49,4261755545 49,43669679151 49,436786848819	49.440976176193	49.443078333822 49.445185487360 49.447297636769 49.461536923042 49.451536923042 49.45579619332 49.457933320489 49.467933320489	49.462222563675	49.464374678611 49.466531789054 49.46869384563 49.47308696294 49.473033093006 49.475210185055 49.477392272400 49.479579374995	49,483968505767	49 486170573854 49.488377637017 49.490589695209 49.495266748388 49.495028796506 49.497£5589518 49.499467877378 49.501724910040	49,506213959581
ber'(x)	-48.399991111135	-48.4029543765 -48.40502345719 -48.402110869734 -48.412110869734 -48.41517858429 -48.41359607137 -48.42359607137 -48.427601669267	-48,430745683936	-48.43305070964 -48.437079867821 -48.440270111972 -48.4449475840886 -48.44697092028 -48.449933902865 -43.453186310864 -48.459738068206	-48,463037492481	-48.466352663778 -48.469683619561 -48.473030397295 -48.47971568468 -48.4816603834 -48.480002926437	-48.496904002949	-48.500378704950 -48.503869566062 -48.510899915460 -48.510899915460 -48.517995350825 -48.521567569392 -48.525156171828	-48.532382678138	-48.536020656927 -45.539675169416 -48.543346253061 -48.560733245318 -48.55073328343 -43.55445935492 -48.558197048320 -48.551951549581	-48.569510977221
ber(x)	50.999600004444	50.999595989509 50.999587869042 50.99958785209 50.999579626774 50.999571261494 50.99957264358 50.999571261494 50.999567032343	50.999558480259	50.999554157018 50.999549802106 50.999540956652 50.999534096652 50.99953265259 50.999527547070 50.999528880 50.999518417931	50.999513804066	50.999509157129 50.999504476961 50.999499763404 50.999495016300 50.9994850450815 50.999480572116 50.999485720815 50.999475639232	50.999465820270	50.999460833870 50.999455812642 50.99945665056 50.9994353372 50.99943537612 50.999430178411 50.9994294306	50.999414369527	50.999409027524 50.999403649059 50.999392732078 50.999381782078 50.999381767257 50.9993762320 50.9993765320	50,999359288747
×	0047	. 401 . 403 . 404 . 405 . 405 . 407 . 408	.410	411 412 414 414 415 415 417 418	.420	. 422 . 423 . 424 . 425 . 426 . 427 . 428	. 430	431 432 434 434 435 435 437 6437	055.	. 441 . 443 . 444 . 444 . 445 . 445 . 445	.450

×	.450	. 451 . 453 . 453 . 454 . 455 . 456 . 456 . 457 . 458	. 460	. 461 . 462 . 463 . 464 . 465 . 465 . 467 . 468	0.470	. 471 . 473 . 473 . 475 . 476 . 477 . 477	. 480	.481 .483 .483 .485 .485 .485 .487 .487	067.	. 491 . 492 . 494 . 495 . 496 . 497 . 498	.500
kei'(x)	50.322597720041	50, 322832992520 50, 323067278977 50, 323306582101 50, 32352904571 50, 32334618312 50, 324224014685 50, 32445241111 50, 324679900114	50, 324906394309	50.325131926298 50.32536498677 50.32580114026 50.325802774918 50.326024483917 50.3264524353 50.32646505428 50.326683925015	50, 327118839459	50.327334890330 50.327550006960 50.327764191831 50.3271847417 50.328189776179 50.328401180572 50.32861263038 50.32882125014 50.329029871924	50, 329237603183	50. 329444422198 50. 32955331366 50. 329853333076 50. 330659429706 50. 330666312769 50. 330666312769 50. 33106641928	50,331265134891	50.31462961817 50.31855958856 50.31855958856 50.312245428765 50.312438846746 50.312631389768 50.32823060087	50.333203791604
kei(χ)	-50.687986367932	-50.687663652493 -50.687340702275 -50.687017518263 -50.686694101438 -50.686045073255 -50.685722463868 -50.68539812559 -50.685073559308	-50.684748766030	-50.684423746840 -50.684098502548 -50.683773034162 -50.683121428929 -50.683121428929 -50.68279529386 -50.68268938758 -50.682142364188	-50,681488560798	-50.68116133385 -50.680833891328 -50.680506234152 -50.6805082355 -50.679850279566 -50.67952193477512 -50.6798338333866	-50,678206701552	-50.677877360463 -50.677547813011 -50.677218060103 -50.67888102646 -50.678587941545 -50.6785897012009 -50.678566245372 -50.6785682	-50.674904112831	-50.674572748708 -50.6734241187202 -50.673909421198 -50.673577475579 -50.673574357724 -50.672912985014 -50.672580449822 -50.672580449822 -50.672580449822	-50.671581695094
(x),rey	-51.205737195204	-51.205262198064 -51.204739403995 -51.204218798571 -51.203700367497 -51.202184096601 -51.20269971836 -51.202157979280 -51.201648105129	-51,200634657438	-51.200131056890 -51.199629520730 -51.199130035745 -51.19862258835 -51.19764375707 -51.19764375707 -51.1976662923885	-51.195689987452	-51.195206449609 -51.194724849005 -51.194245173513 -51.193767411109 -51.192817577971 -51.19284583691 -51.191875255401 -51.191875255401	-51,190940350779	-51.190475651677 -51.190012773023 -51.189551703670 -51.189692432558 -51.188179241282 -51.187725299454 -51.187725299454	-51.186373961090	-51.18592697593 -51.185481703081 -51.185038133283 -51.184596256015 -51.184156061171 -51.183717538728 -51.183717538728 -51.183717538728 -51.18374746	-51.181979975332
ker(x)	50.952650211570	50.950594966446 50.948544960265 50.946500171071 50.944460577046 50.94246155520 50.938372749976 50.938372749976 50.936333721314 50.936333721314	50,932330907630	50.930327080784 50.92832479611 50.926334483532 50.924545672102 50.9203252055 50.918408943192 50.916439868487 50.916475678134	50,912516352452	50.910561871887 50.908612217003 50.90667368439 50.904727307155 50.900861469860 50.89393565111 50.893935656111 50.895098144822	50,893186410191	50.891279331700 50.889376891089 50.867479070208 50.83583851021 50.83164125 50.87993362488 50.876058634286 50.876188156822	50.874322175108	50.872460671856 50.870603629886 50.868751032119 50.866902861579 50.965059101391 50.863219734781 50.861384745075 50.881384787830173	50.855905872120
bei'(x)	50.224951946314	\$0.225451410010 \$0.225950863929 \$0.226450323039 \$0.225949772307 \$0.227948656189 \$0.228448090739 \$0.228446944894	50,229946364434	50.230445778904 50.230945188271 50.231444592503 50.23149391565 50.232443335425 50.23242774040 50.233441535449 50.23344153549	50.234940275497	50.235439637427 50.23593893915 50.236438344927 50.235937690427 50.237437030381 50.237435693508 50.238435693508 50.238435693508	50.239933645717	50.240432951649 50.240431546038 50.241431546038 50.241930834523 50.242430117053 50.2434229393642 50.2434229393642 50.243927928847 50.244427187389	50,244926439842	50.245425686167 50.245924926327 50.246422188803 50.24692388803 50.24792182609442 50.247921824564 50.24842103331 50.24842103331 50.249419431646	50.249918621116
bci(x)	49.506213959581	49.502465976367 49.510722887766 49.512984993730 49.515251994210 49.5152388160 49.52208262267 49.524369940326 49.52436912657	49.528958879207	49.531260839926 49.533567794468 49.535879743677 49.540515622491 49.542845554292 49.545177478954 49.547514397423	49.552203215568	49.556912008298 49.556912008298 49.5569273894997 49.564012643787 49.566389515768 49.56877137C064 49.571158229619	49.575946916280	49 .578348749272 49 .580755578294 49 .583167394288 49 .58584206196 49 .590432808517 49 .592864598812 49 .592301381782 49 .597743157369	49.600189925510	49.602641686145 49.605098439213 49.607560184651 49.61002692298 49.61498652390 49.61767088860 49.61767088860 49.61767088860	49.624932133822
ber'(x)	-43,569510977221	-48.577137886042 -46.577137888042 -48.580976743278 -48.58470544061 -48.58870544061 -48.6004265137711 -49.600426515784	-43.608326355489	-43,612302124577 -43,616295176421 -48,620305548469 -45,624333278168 -48,632440360312 -48,6352093764296 -48,64618522426 -48,644733602086	-48.648866264076	-48,653016545840 -48,65778448423 -48,661370118467 -48,665573484217 -48,66794619514 -48,674033561801 -48,678290345521 -48,688857605020	-43.691168149680	-48.695496688535 -48.69984325902 -48.704207893582 -48.70590644652 -48.717410606073 -43.717410606073 -45.72847696298 -48.726303442782	-48.735269454264	-48.73979994134 -48.744303940000 -48.748856329298 -48.753422199460 -48.758006587913 -48.762609532104 -48.771871237387 -48.771871237387	-48.781207614751
ber(x)	50.999359288747	50,999353574626 50,999347822371 50,999342031812 50,9993032428104 50,999324428614 50,999318493140 50,999318498510 50,999306474552	50.999300411096	50.993294307968 50.99928144996 50.999281982006 50.999275753827 50.999264931201 50.9992564460725 50.999250460722	50.999237565992	50.999231056592 50.999224505602 50.999217912844 50.999210781/1 50.999204601315 50.999191120584 50.999184316322 50.999184316322	50.999170579110	50.999163645301 50.999156669116 50.999142564894 50.99913427602 50.99912312507 50.999121128729 50.999113687988 50.999113687988	50,999099272331	50. 999001897149 50. 999084476720 50. 999077010909 50. 999061942403 50. 999065332 50. 999046590151 50. 999038994655 50. 999031252664	50.999023463991
×	.450	.451 .452 .454 .455 .456 .457 .458	097.		0.470	471 472 473 475 476 477 478	.480	.481 .483 .484 .485 .486 .486 .486	067.	. 491 . 493 . 495 . 496 . 496 . 498 . 498	, 50ġ